

The X-ray spectral properties of the AGN population in the *XMM-Newton* bright serendipitous survey

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Received , ; accepted ,

ABSTRACT

Context. X-ray surveys are a key instrument in the study of active galactic nuclei (AGN). Thanks to their penetrating ability, X-rays are able to map the innermost regions close to the central super massive black hole (SMBH) as well as to detect and characterize its emission up to high redshift.

Aims. We present here a detailed X-ray spectral analysis of the AGN belonging to the *XMM-Newton* bright survey (XBS). The XBS is composed of two flux-limited samples selected in the complementary 0.5-4.5 and 4.5-7.5 keV energy bands and comprising more than 300 AGN up to redshift ~ 2.4 .

Methods. We performed an X-ray analysis following two different approaches: by analyzing individually each AGN X-ray spectrum and by constructing average spectra for different AGN types.

Results. From the individual analysis, we find that there seems to be an anti correlation between the spectral index and the sources' hard X-ray luminosity, such that the average photon index for the higher luminosity sources ($> 10^{44}$ erg s⁻¹) is significantly ($> 2\sigma$) flatter than the average for the lower luminosity sources. We also find that the intrinsic column density distribution agrees with AGN unified schemes, although a number of exceptions are found (3% of the whole sample), which are much more common among optically classified type 2 AGN. We also find that the so-called “soft-excess”, apart from the intrinsic absorption, constitutes the principal deviation from a power-law shape in AGN X-ray spectra and it clearly displays different characteristics, and likely a different origin, for unabsorbed and absorbed AGN. Regarding the shape of the average spectra, we find that it is best reproduced by a combination of an unabsorbed (absorbed) power law, a narrow Fe K α emission line and a small (large) amount of reflection for unabsorbed (absorbed) sources. We do not significantly detect any relativistic contribution to the line emission and we compute an upper limit for its equivalent width (EW) of 230 eV at the 3σ confidence level. Finally, by dividing the type 1 AGN sample into high- and low-luminosity sources, we marginally detect a decrease in the narrow Fe K α line EW and in the amount of reflection as the luminosity increases, the “so-called” Iwasawa-Taniguchi effect.

Key words. X-rays: general- X-rays: diffuse background - surveys - galaxies:active

1. Introduction

Recent deep X-rays surveys carried out by *XMM-Newton* and *Chandra* have resolved most of the cosmic X-ray background (CXB) into discrete sources up to energies ~ 10 keV (although the resolved fraction decreases with energy; Worsley et al. 2005; Hickox & Markevitch 2006). The large majority of the sources that compose the CXB are active galactic nuclei (AGN), and CXB synthesis models make use of template AGN spectra to reproduce its shape following the AGN unified model (Antonucci 1993). The unified model, in its simplest version, states that the differences between the different observed AGN types are due to an orientation effect, i.e., as the inclination angle to the observer increases, the torus surrounding the central super massive black hole (SMBH) intercepts more nuclear emission. The CXB is then reproduced by a mixture of AGN spectra with different amounts of absorption.

However, there are still many unresolved questions regarding our knowledge about AGN. For example, the predicted fraction of heavily absorbed AGN (Compton-thick AGN) obtained from CXB synthesis models can vary from 30% to

9% between different works (Gilli et al. 2007; Treister et al. 2009). Besides, a small number of AGN that seem not to follow the unified scheme are usually found in X-rays surveys, i. e., their optical characteristics do not match their observed X-ray properties (Panessa & Bassani 2002, Akylas et al. 2004, Caccianiga et al. 2004, Cappi et al. 2006, Mateos et al. 2005a, Mateos et al. 2005b, Mateos et al. 2010). The evolution of these properties through cosmic time and the possible correlation between X-ray emission and source properties, like the bolometric luminosity or SMBH mass, are also a matter of debate.

The frequency and properties of some individual characteristics are also still unknown, such as the Fe K α emission line. This emission line is the most commonly observed line in AGN X-ray spectra, but its detailed study is strongly limited by the data quality and therefore, to sources in the local Universe (Nandra et al. 2007). To study its characteristics up to high redshifts, X-ray spectra have to be stacked to improve the signal-to-noise ratio (SNR) (Corral et al. 2008; Streblyanska et al. 2005; Brusa et al. 2005). Another intriguing AGN feature is the soft-excess emission in type 1 AGN, whose origin is still unclear.

Possible suggested explanations go from continuum emission (Ross et al. 1992; Shimura & Takahara 1993; Kawaguchi et al. 2001) to atomic processes (Crummy et al. 2006; Middleton et al. 2007).

Deep and medium surveys often lack good quality X-ray and multi-wavelength data, which limits the results, while samples composed by high-quality data are usually not well-defined flux-limited samples, which limits the applicability of the results. Given all that, well-defined X-ray samples, that contain both a significant number of reliable identifications and good enough X-ray data quality are the key to test the current hypotheses and to link the nearby and distant universe. We present here a detailed X-ray spectral analysis of the AGN within the XBS sample, which is composed of two flux-limited samples that are almost completely identified (identification rate $\sim 95\%$) and containing more than 300 AGN. Given the availability of both reliable optical spectroscopic identifications and good quality X-ray spectral data, this sample is the perfect laboratory to test AGN models and to better constrain the AGN properties and their evolution.

We assume the cosmological model $H_0=65 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_\Lambda=0.7$ and $\Omega_M=0.3$ throughout this paper. Reported errors are at 90% confidence level unless stated otherwise.

2. The XBS AGN sample

The sample of 305 AGN discussed here (XBS AGN sample hereafter) has been extracted from the *XMM-Newton* bright serendipitous survey¹.

The XBS consists of two flux-limited serendipitous (i.e. the targets of the *XMM-Newton* pointings were excluded) samples of X-ray selected sources at high galactic latitude ($|b| > 20^\circ$): the *XMM* bright serendipitous survey sample (BSS, 389 sources) and the *XMM* hard bright serendipitous survey sample (HBSS, 67 sources, with 56 sources in common with the BSS sample) having an EPIC MOS2 count rate limit, corrected for vignetting, of 10^{-2} cts/s and $2 \times 10^{-3} \text{ cts/s}$ in the 0.5-4.5 keV and 4.5-7.5 keV energy bands, respectively; the flux limit is $\sim 7 \times 10^{-14} \text{ erg cm}^{-2} \text{ s}^{-1}$ in both energy selection bands.

The details on the *XMM-Newton* fields selection strategy and the source selection criteria of the *XMM* BSS and HBSS samples are discussed in Della Ceca et al. (2004), while a description of the optical data and analysis, of the optical classification scheme and the optical properties of the extragalactic sources identified so far is presented in Caccianiga et al. (2007, 2008). The optical and X-ray properties of the Galactic population are discussed in López-Santiago et al. (2007). Previous X-rays spectral analyses of parts of the XBS sample have already been reported in previous works. In Caccianiga et al. (2004), the X-ray spectral analysis of a subsample extracted from the HBSS sample is reported. Galbiati et al. (2005) performed an analysis of the radio-loud AGN within the XBS. Severgnini et al. (2003) unveiled the AGN-nature of three sources previously considered as normal galaxies. Finally, Della Ceca et al. (2008) presented the cosmological properties of the HBSS AGN sample.

¹ The *XMM-Newton* Bright Serendipitous Survey is one of the research programs conducted by the *XMM-Newton* Survey Science Center (SSC, see <http://xmmssc-www.star.le.ac.uk>) a consortium of 10 international institutions, appointed by ESA to help the SOC in developing the software analysis system, to pipeline process all the *XMM-Newton* data, and to exploit the *XMM* serendipitous detections. The *Osservatorio Astronomico di Brera* is one of the Consortium Institutes.

2.1. AGN classification

The current classification breakdown of the XBS sample, which relies on dedicated optical spectroscopy, is as follows: 305 AGN (including 5 BL Lacs), 8 clusters of galaxies², 2 normal galaxies and 58 X-ray emitting stars, see Table 1 for a detailed summary. For 25 out of the 305 AGN that composed our sample, redshift and classification are reported here for the first time. These new identifications are marked in boldface in Table 3 (columns 2 and 3). The XBS AGN sample contains 35 sources that are optically classified as elusive AGN, i.e., sources for which a classification cannot be derived solely from our optical spectroscopy, although the redshift can be measured. These are sources that are characterized by a significant/dominant contamination of star-light from the host galaxy in the optical spectrum (Caccianiga et al. 2007). Even if the presence of an AGN in these sources is somehow suggested by the detection of a broad or strong emission line, the “dilution” caused by the host galaxy is critical because it avoids the quantification of the optical absorption that is necessary to classify a source as a type 1 or type 2 AGN. For these sources, the type 1/2 classification is assigned as a function of the absence/presence of a significant amount of intrinsic absorption in their X-ray spectra. There is one case however, XBSJ012654.3+191246, in which a type 1/2 classification cannot be inferred from either the optical or the X-ray data, and accordingly this source is classified as an AGN of uncertain type. At the time of writing, 27 X-ray sources belonging to the BSS sample are still unidentified. Out of these 27, two also belong to the HBSS sample, which results in a level of identification of 93% and 97% for the BSS and the HBSS samples, respectively.

2.2. X-ray data

The XBS source sample was defined using only the data from the MOS2 detector. However, to increase the statistics, the data from the MOS1 and the pn detectors were considered when available and were used for our spectral analysis.

In Table 2 we report the data used for the X-ray spectral analysis of each source: Source name; *XMM-Newton* observation ID; *XMM-Newton* filter for each detector; the values of Galactic column densities toward the used *XMM-Newton* pointings; resulting exposure time for each detector after removing high-background intervals; total counts for all available detectors in the 0.3-10 keV band, and corresponding sample. To increase the number of counts for the lowest quality spectra, we searched the *XMM-Newton* archive for additional observations and selected those with the longest MOS2 exposure times; we preferred not to combine different observation data sets to minimize possible problems related to source variability. As a result, some of the data sets used in this analysis are different from the ones used in Della Ceca et al. (2004) for the definition of the sample.

The *XMM-Newton* data were cleaned and processed with the *XMM-Newton* Science Analysis Software (SAS) and were analyzed with standard software packages (Ftools; Xspec, Arnaud 1996). Event files produced from the pipeline were filtered from high-background time intervals and only events corresponding to pattern 0-12 for MOS and 0-4 for pn were used. All spectra were accumulated from a circular extraction region with a radius of $\sim 20''$ - $30''$, depending on the source off-axis distance.

² The sample of cluster of galaxies is neither statistically complete nor representative of the cluster population because the source detection algorithm used in the construction of the sample is optimized for point-like sources.

Table 1. XBS classification summary.

Sample	Type 1 AGN	Type 2 AGN	BL Lacs	Stars	Clusters	Galaxies	Unidentified	AGN uncertain type
BSS	269(41)	19(10)	5	58(2)	8(1)	2	27(2)	1
HBSS	42	20	...	2	1	...	2	...
Total	270	29	5	58	8	2	27	1

The numbers between parenthesis for the BSS sample correspond to the number of sources in common with the HBSS sample.

Background counts were accumulated in nearby circular source free regions which an area usually about a factor ~ 4 larger than the one used to extract the source counts. To improve the statistics, the MOS1 and MOS2 spectra obtained with the same filters were combined *a posteriori* by using the FTOOLS task *mathpha*. The X-ray spectra usually cover the 0.3–10 keV energy range; the total (MOS1+MOS2+pn) counts range from ~ 100 to $\sim 10^4$, as can be seen in Fig. 1.

The ancillary response matrix and the detector response matrix were created by the XMM-SAS task *arfgen* and *rmfgen* at each source position in the EPIC detectors. For the MOS1 and MOS2 detectors, and provided that the observations were carried out by using the same filter, ancillary and detector response matrices for each source were combined by using *addrmf* and *addarf*.

3. X-ray spectral analysis

The availability of good *XMM-Newton* data for the sources in the XBS sample, which spans the energy range between ~ 0.3 and ~ 10 keV, allow us to perform a reliable X-ray spectral analysis for almost every AGN studied here. For 111 AGN an X-ray spectral analysis was already reported and discussed in Severgnini et al. (2003), Caccianiga et al. (2004, 2007), Galbiati et al. (2005) or in Della Ceca et al. (2008); for the remaining AGN the main X-ray spectral properties and parameters are discussed in detail here for the first time. Note, however, that there could be small differences in the best-fit model and parameters already published and the ones presented here owing to the different *XMM-Newton* observations used and/or our different way of defining the best-fit model for each source.

We grouped the spectra in bins containing at least 10 to 30 (depending on the spectral quality) source+background counts to use the χ^2 minimization technique. We fitted pn and MOS spectra simultaneously in the 0.3–10 keV band with Xspec version 12.5.0. We tied together all pn and MOS parameters except for a relative normalization, which accounts for the differences between pn and MOS flux calibrations. In the following, derived fluxes and luminosities refer to the MOS2 calibration.

To ensure a spectral analysis as uniform as possible, we defined a threshold of 10% for the null hypothesis probability to distinguish between an acceptable and an unacceptable fit, i.e., we consider as our best-fit model the simplest model for which the probability is $>10\%$.

As a starting point for the spectral modeling we first considered a simple absorbed power-law model that takes into account both the Galactic hydrogen column density along the line of sight (from Dickey & Lockman 1990) and a possible intrinsic absorption at the source redshift (Xspec model: *wabs*zwabs*zpo*). In the X-ray spectral modeling we made use of the redshifts obtained from the optical spectroscopy.

The results for this simple fit are shown in Table 3 along with the corresponding Galactic de-absorbed flux and intrinsic

luminosity in the standard hard (2–10 keV) energy band. In some cases, the spectral quality does not allow us to constrain the power-law photon index (Γ) and the intrinsic absorption at the same time. In other cases, the resulting photon index is ~ 1 , much lower than the typical values for unabsorbed AGN³. In those cases, we fixed Γ to 1.9, a common value for unabsorbed AGN (Caccianiga et al. 2004; Mateos et al. 2005a,b; Galbiati et al. 2005; Tozzi et al. 2006; Mateos et al. 2010). If there was no intrinsic absorption detected, an upper limit, at 90% confidence level, is given. The simple absorbed power-law model gives a good fit for 263 sources, but seems to fail in reproducing the spectral shape for 41 sources, marked with a ^p in Table 3. All X-ray spectra corresponding to sources classified as BL Lacs are well fitted by the simple power-law model.

For the 41 sources that are not well fitted by an absorbed power law, we tried several additional components to the absorbed power law model. We accept any of these additional components if the improvement of the fit was larger than 95% as measured by an F-test. These additional components are

- **Leaky absorbed power-law:** An additional unabsorbed power-law component, with the same photon index as the direct one, representing scattered emission into our line of sight (Xspec model: *wabs(zwabs*zpo+zpo)*). This model can also account for partially covered emission.
- **Ionized absorption:** (Xspec model: *wabs*zwabs*absori*zpo*, Magdziarz & Zdziarski 1995), since signatures of absorption from partially ionized gas have been found to be a quite common characteristic in the spectra of Seyfert galaxies.
- **Reflected component:** to account for a spectral hardening or change of curvature at high energies because of Compton reflection from neutral material (Xspec model: *wabs(zwabs*zpo+pexrav)*, Magdziarz & Zdziarski 1995). We fixed the inclination angle to $\sim 60^\circ$, an average value for Seyfert galaxies, because the spectral quality does not allow us to constrain it and the reflection factor R ($R=\Omega/2\pi$) at the same time.
- **Thermal component** to account for soft emission lines that could arise from ionized material far from the central source, like the narrow-line region (NLR) (Xspec model: *wabs(zwabs*zpo+mekal)*, Mewe et al. 1986; Liedahl et al. 1995). Although the NLR is likely photoionized, we can ap-

³ An alternative possibility is that these sources with an observed flat spectrum are Compton-thick AGN (i.e. sources with $N_H > 10^{24} \text{ cm}^{-2}$) in which all the direct emission is suppressed and only reflected emission is observed (in the 2–10 keV band). This effect, combined with the low statistics, may mimic a flat spectrum. However, this hypothesis does not seem to be valid in our sources because in all cases we find a significant amount of absorption (but not in the Compton-thick regime) even when leaving the photon index as a free parameter (N_H from 3.5×10^{21} to $2 \times 10^{23} \text{ cm}^{-2}$). We conclude that the best explanation for the sources with a very flat spectral index is the combination of (mild) absorption and the low statistics

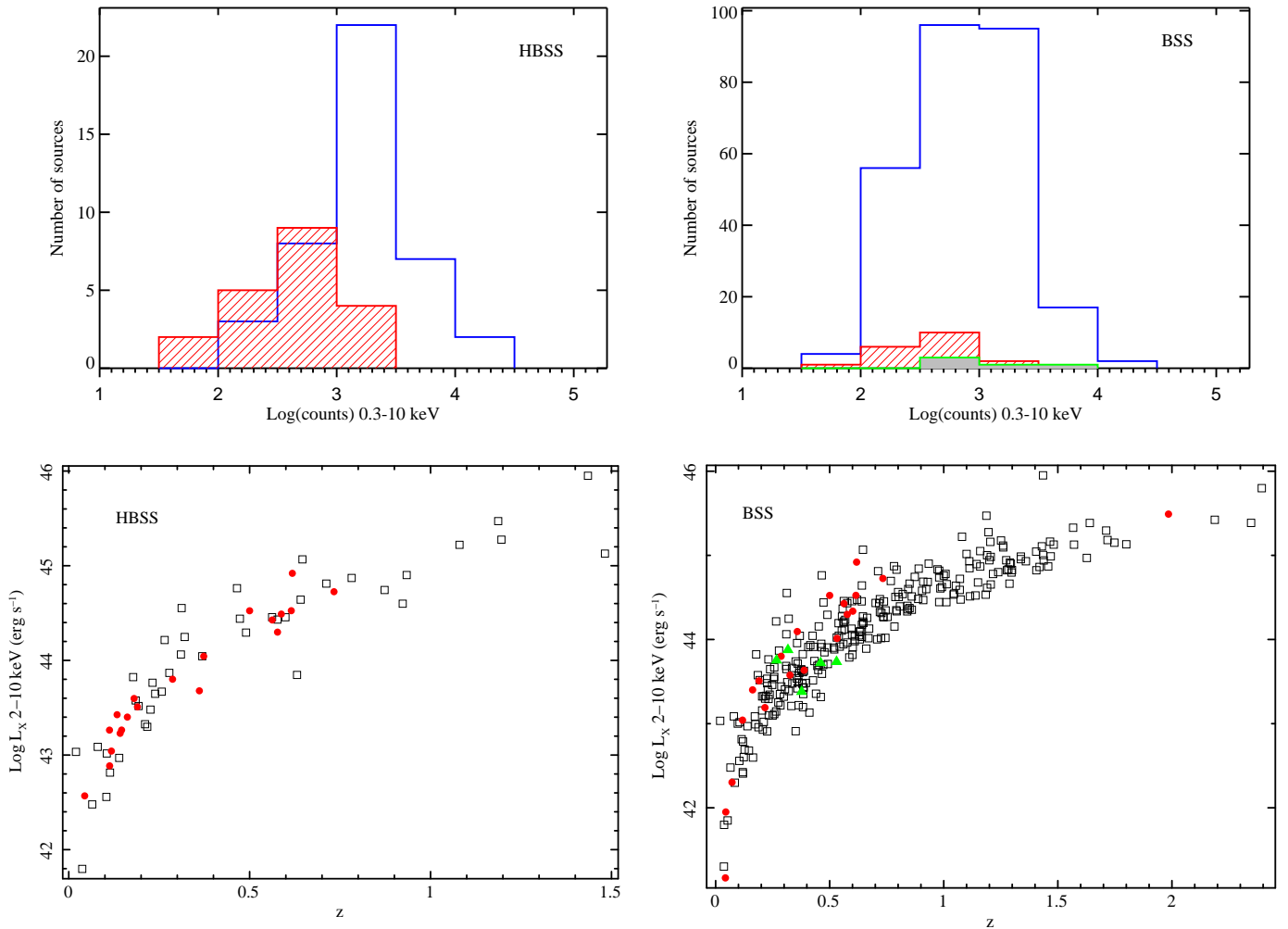


Fig. 1. Top panels: Counts distribution for the HBSS (left) and the BSS (right) samples. Empty histograms correspond to type 1 AGN and dashed histograms to type 2 AGN. Filled histogram for the BSS corresponds to BL Lacs. Bottom panels: X-ray intrinsic luminosity versus redshift for the HBSS (left) and the BSS (right). Squares and circles correspond to sources classified as type 1 and type 2 AGN, respectively. Triangles on the BSS sample represent BL Lacs.

proximate the resulting spectral shape by using this collisional model given the spectral resolution for the EPIC cameras.

- **Lines and edges** (`zgauss`, `zegde`) to model emission lines (such as the Fe K α emission line, the most commonly observed one in AGN X-ray spectra) and absorption edges. Energies were left free to vary.
- **A phenomenological black body model** (`Xspec` model: `wabs(zwabs*zpo+zbb)` to account for featureless soft-excess emission.

If different additional components significantly improved the fit, we selected the model that was more physically plausible and/or gave better residuals. These cases, 28 AGN, are discussed in more detail in the appendix. There are also two cases in which more than one additional component are required to obtain an acceptable fit: in both cases, one of the required additional components turned out to be an emission line. A summary of the models required during the spectral fit is shown in Table 4, while the results for the additional component fits are presented in Table 6 to Table 11. All models that significantly improve the simple

absorbed power-law fit are shown for each source in Tables 6 to 11. The model we considered as the “best-fit” for each source is marked in boldface in those tables and its parameters are the ones we consider in the interpretation of the results.

In 11 cases still no acceptable fit was found. This can be simply due to a statistical effect given the 10% probability limit imposed to consider a fit as acceptable. Roughly, 10% of the sources that actually display a power-law shape are expected to be not acceptably fitted by this model. We did not find these 11 sources to share any common spectral characteristic. For 9 of them, we found that no additional component significantly improved the fit, and we accordingly assume the simple absorbed power-law fit as our best fit and used the data in Table 3 in the subsequent analyses. For the remaining two sources, XBSJ021822.2–050615 and XBSJ153456.1+013033, a leaky absorber plus an emission line and a black body, respectively, did improve the fit significantly, consequently we consider these models as our best-fit model although the probability is still <10%.

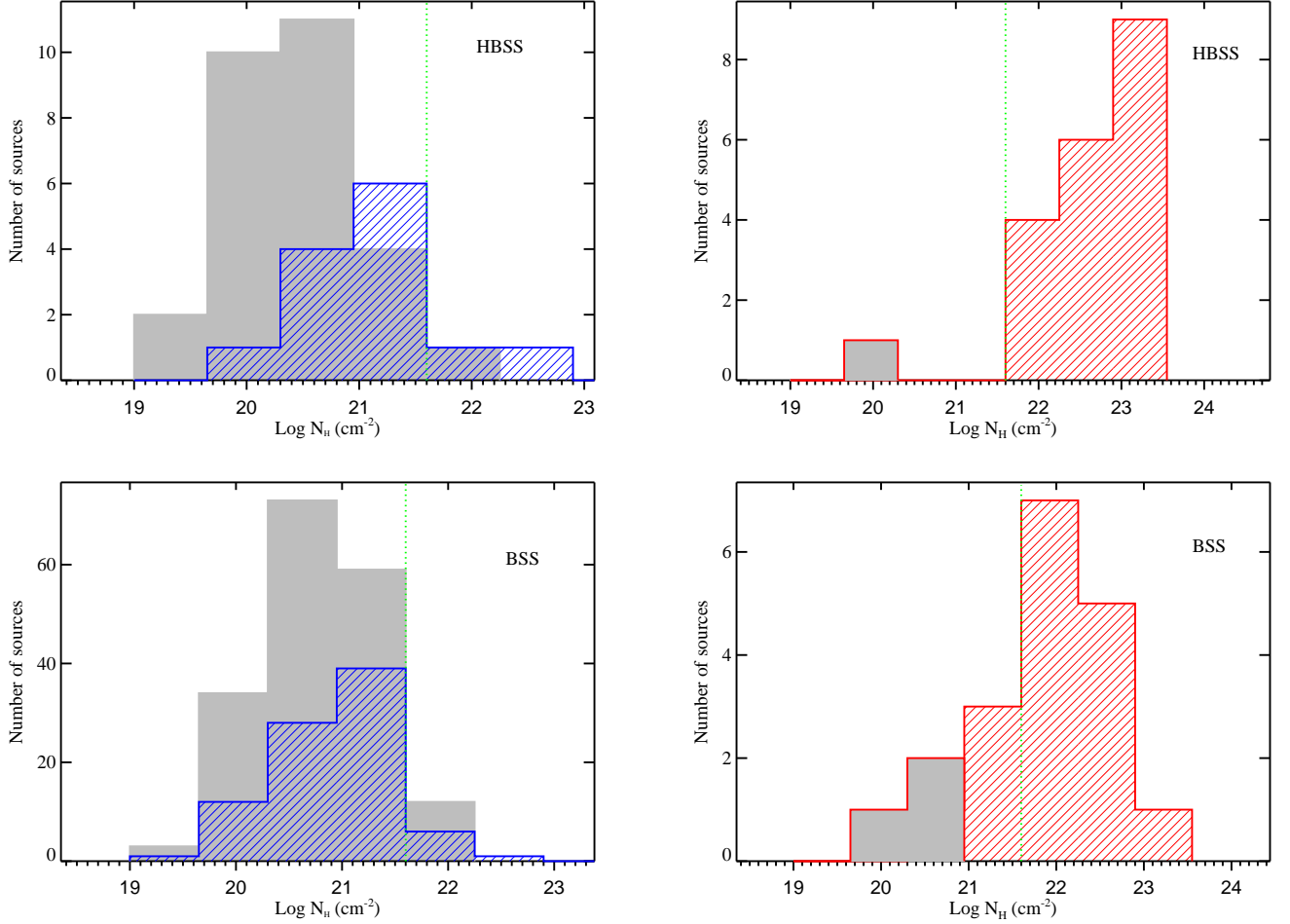


Fig. 2. Intrinsic absorption distribution corresponding to the HBSS (top) and BSS (bottom). Left and right panels correspond to type 1 and type 2 AGN, respectively. Dashed histograms correspond to detected intrinsic absorption, whereas filled histograms correspond to upper limits. Vertical dotted lines mark the threshold used to distinguish between absorbed and unabsorbed sources following the prescription of Caccianiga et al. (2007).

Table 4. X-ray spectral fit

Model	Type 1 AGN	Type 2 AGN	BL Lacs
SPL	243	24	5
Leaky	1	3	...
PL+T	...	2	...
PL+BB	13
WAPL	5
PL+R	4
PL+E	3
Leaky+L	1	1	...

SPL: absorbed power law; Leaky: absorbed plus unabsorbed power law. PL+T: absorbed power law plus thermal component; PL+BB: absorbed power law plus black body component; WAPL: ionized absorbed power law; PL+R: absorbed power law plus neutral reflection component; PL+E: absorbed power law plus absorption edge; Leaky+L: leaky model plus emission line.

4. Intrinsic absorption

We measured intrinsic absorption in excess of the Galactic one for 119 sources, 88 type 1 and 27 type 2 AGN. This absorption is significant (F-test > 95%) in 56 cases (30 type 1 and 26 type 2 AGN). For the remaining 65 sources for which the significance is below 95%, the measured amount of absorption is very low, except for one case, XBSJ161820.7+124116. The low significance in this case is likely owing to the extremely low number of counts in the available spectrum. The intrinsic absorption distribution for both studied samples is shown in Fig. 2.

In 17 cases (3 type 1 and 14 type 2 AGN) and because of the poor statistics, we fixed the value of Γ to 1.9 (which corresponds to the average value found for unabsorbed AGN) to better constrain the intrinsic absorption. For 7 out of these 17 AGN, Γ could be determined (although with large errors), but turned out to be flatter than the flattest Γ found for unabsorbed AGN ($\Gamma \sim 1.5$). This is probably owing to the low statistics available for the spectral analysis, which do not allow us to adequately constrain at the same time both the spectral index and the intrinsic absorption in these cases. We note here that the best fit N_H obtained with free Γ is usually within the reported errors

obtained when fixing $\Gamma = 1.9$; furthermore the variation of N_H is such that this problem does not have any effect on the X-ray source classification used here (absorbed vs. unabsorbed) or on the N_H distribution.

4.1. X-rays versus optical absorption

Following the criteria described in Caccianiga et al. (2008)⁴, we defined a sourced as absorbed if the measured intrinsic column density is higher than $4 \times 10^{21} \text{ cm}^{-2}$. According to this criterion, we find that 31 sources are absorbed, 8 type 1 (3% of type 1 AGN) and 23 type 2 AGN (80% of type 2 AGN). We find that the fraction of absorbed sources and the amount of absorption is much higher for type 2 AGN than for type 1 AGN, in agreement with the AGN unified picture. However, there is a number of sources that do not match within this scenario, 8 type 1 AGN are absorbed (3% total) and 6 type 2 AGN are unabsorbed (20% total). Nonetheless, for half of the absorbed type 1 AGN, the errors in the intrinsic column density are consistent with these sources being unabsorbed. It is important to note that the expected intrinsic column density derived from optical reddening depends on the assumed gas-to-dust ratio. If the intrinsic gas-to-dust ratio differs significantly from the Galactic value (the standard gas-to-dust ratio that is usually used), small differences between the expected value from optical observations and the measured value from X-rays are expected. Three of the unabsorbed type 2 AGN are consistent with being absorbed within errors, but only by low amounts of intrinsic absorption ($\leq 10^{22} \text{ cm}^{-2}$). This implies that the optical/X-ray type mismatch could be more common for sources that are optically classified as type 2 AGN. However, it is difficult to quantify how frequent the “mismatches” between optical and X-ray classification are.

In an X-ray selected sample, the least biased way to estimate the exceptions to the unified models is to compute the fraction of optically classified type 1/2 AGN among the absorbed/unabsorbed sources. Elusive AGN have to be removed because their classification is derived from X-rays. We also excluded those absorbed/unabsorbed sources that could be unabsorbed/absorbed within errors. In this way, we found that there are only four type 1 AGN among “truly” absorbed sources and two type 2 AGN among “truly” unabsorbed sources. Therefore, the resulting fraction of exceptions to the unified models for the XBS AGN sample turns out to be of $\sim 3\%$. For the HBSS only, which is almost completely identified, we find that the fraction of type 2 AGN among unabsorbed sources is only 3%. A similarly low value of 1% is obtained for the BSS. For the fraction of type 1 AGN among absorbed sources, we find 17% for the HBSS, whereas for the BSS it turns out to be $\sim 30\%$. However, if we take into account the larger number of unidentified sources on the BSS, seven of which are probably absorbed AGN (see discussion at the end of this section), this number could decrease to match the one obtained for the HBSS.

To compare our findings with previously reported results, we selected a 10^{22} cm^{-2} limit, which is the one that is usually used by other authors, to separate between absorbed and unabsorbed sources. Using this limit and considering the total XBS sample, we computed a fraction of unabsorbed type 2 AGN among the total number of type 2 AGN of $\sim 36\%$. This number agrees with reported values of unabsorbed type 2 AGN shown

in Panessa & Bassani (2002) and Akylas & Georgantopoulos (2009) ($\sim 10\text{-}30\%$ and $\sim 20\%$, respectively). However, it has to be pointed out that none of the reported fractions in Panessa & Bassani (2002) and Akylas & Georgantopoulos (2009) has been derived from complete samples. For example, if we consider only the HBSS, which is almost completely identified and that it is less biased against absorbed sources, the value decreases from 36% to 20%; this is expected to be a more reliable fraction than that computed using the total XBS sample. Moreover, errors in the resulting N_H values are not usually considered either. If we remove unabsorbed type 2 AGN that could be absorbed within 90% confidence errors, the fraction decreases to $\sim 5\%$. Therefore, caution must be exercised when computing the fraction of exceptions to unified models; these fractions have probably been overestimated in the past.

The existence of unabsorbed type 2 AGNs has no clear explanation so far. Some recent models (Elitzur & Shlosman 2006; Nicastro 2000) show that the BLR could not form under particular condition. For instance, it has been proposed that the BLR may disappear below bolometric luminosities of $\sim 10^{42} \text{ erg s}^{-1}$ (Elitzur & Shlosman (2006)) or below a critical accretion rate ($L_{\text{bol}}/L_{\text{Edd}} \sim 1\text{-}4 \times 10^{-3}$ for SMBH masses ranging from 10^6 to 10^9 solar masses; where L_{bol} and L_{Edd} are the bolometric and Eddington luminosities, respectively, Nicastro 2000). Nevertheless, the range of luminosity and accretion rates covered by the unabsorbed type 2 AGN in our sample (L_X from 10^{42} to $\sim 2 \times 10^{44} \text{ erg s}^{-1}$ and accretion rates from $\sim 10^{-3}$ up to ~ 1 , Caccianiga et al. in preparation) make these interpretations not applicable to the sources of the XBS sample.

A possible alternative explanation is that unabsorbed Type 2 AGN are indeed Compton-thick (CT) i.e. sources where the amount of intrinsic absorption is so high (above 10^{24} cm^{-2}) that the absorption cut-off falls outside the observed spectral range. Using X-ray data limited in the 2-10 keV energy band, it would not be possible to compute the actual column density and we would end up with an optically type 2 AGN with no sign of absorption in the X-rays. The unabsorbed type 2 AGN in our sample are XBSJ012057.4–110444, XBSJ031146.1–550702, XBSJ100032.5+553626, XBSJ141235.8–030909, XBSJ230522.1+122121 and XBSJ221951.6+120123. XBSJ100032.5+553626 is an elusive AGN, whose Compton-thick nature was studied and discarded as a possible explanation in Caccianiga et al. (2007). To test the Compton-thick hypothesis for the remaining five unabsorbed type 2 AGN, we used the diagnostic diagram by Bassani et al. (1999), which make use of the thickness parameter (T) and the Fe K α line EW to separate Compton-thick from Compton-thin sources. The thickness parameter represents the ratio between the 2-10 keV observed flux (corrected for Galactic absorption) and the reddening-corrected flux of the [OIII] $\lambda 5007\text{\AA}$ emission line. Compton-thick sources usually locate at $T < 1$ and large Fe K α equivalent widths. The values for the [OIII] fluxes used here were computed following the prescription and assumptions discussed in Caccianiga et al. (2007). Regarding the Fe K α EWs, only upper limits (at the 90% confidence limit) could be derived; we assumed in all cases an unresolved neutral emission line centered at 6.4 keV.

Our results for the five sources in consideration here are plotted in Fig. 3; obviously all our sources are well above the $F_X/F_{[\text{OIII}]}$ = 1 limit. The computed [OIII] fluxes are not corrected for extinction from the host-galaxy. Correcting for the host-galaxy extinction would increase the [OIII] fluxes thus decreasing the resulting T values. Nonetheless, the maximum expected extinction ($A_V \sim 1$ for galaxies with ongoing intense star-formation, Calzetti & Heckman (1999)) is not high enough to

⁴ The optimum dividing line between optical type 1/2 classification is found to correspond to an optical extinction of $A_V \sim 2 \text{ mag}$, which, assuming a Galactic A_V/N_H ratio, implies a column density of $N_H \sim 4 \times 10^{21} \text{ cm}^{-2}$ in X-rays.

place these sources below the $T = 1$ limit. We conclude that the CT hypothesis is not a valid explanation for the unabsorbed type 2 AGN in the XBS sample. Further investigations are thus required to understand the nature of these intriguing sources.

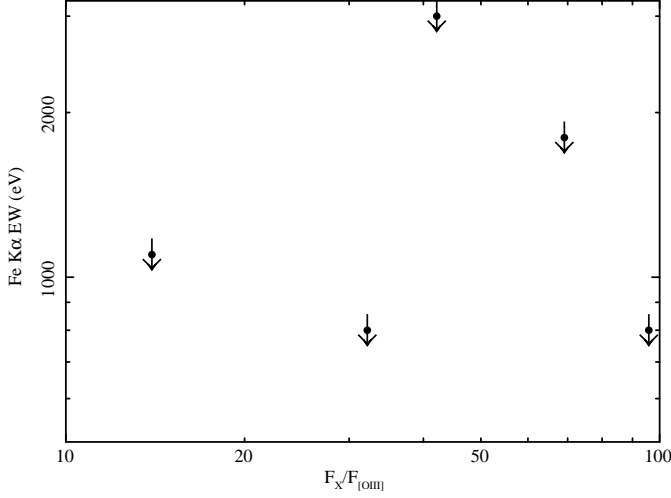


Fig. 3. Fe $K\alpha$ EW versus the thickness parameter $T = F_X/F_{[OIII]}$ for the 5 unabsorbed type 2 AGN. Vertical arrows correspond to upper limits on the Fe $K\alpha$ EW.

Finally the observed fractions of absorbed AGN (number of absorbed AGN to total number of AGN) are $\sim 35\%$ for the HBSS and $\sim 7\%$ for the BSS (34% and 6%, respectively, if we do not take into account those sources that are consistent with being unabsorbed within errors). This difference is expected because the selection at higher energies is less affected by obscuration, therefore the HBSS is much more efficient in selecting absorbed sources. To investigate how the larger number of unidentified sources in the BSS could affect this, we performed a simple power law X-ray spectral fit over the unidentified AGN in the XBS sample. The resulting photon index distribution is shown in Fig. 4. We can see that seven (including the two unidentified sources in the HBSS) out of the 27 unidentified sources display a photon index below 1.5. If these seven sources turn out to be absorbed AGN, the fraction of absorbed AGN in the BSS sample would increase to 10%.

4.2. Hardness ratios

Another way to estimate the source type by using X-rays is to compute the hardness ratios (X-ray colors). We extracted the hardness ratios HR2, HR3, and HR4 from the 2XMM-Newton catalog⁵ (Watson et al. 2009), which are defined as

$$HR_n = \frac{CR_{n+1} - CR_n}{CR_{n+1} + CR_n}, \quad (1)$$

where CR_n is the “vignetting”-corrected count rate in the energy band n . The 2, 3, 4, and 5 energy bands correspond in our case to the count rates in the 0.5–1.0, 1.0–2.0, 2.0–4.5, and 4.5–12.0 keV energy bands, respectively. To compare them with the unidentified sources, the hardness ratios for the identified AGN are plotted along with the ones for the unidentified sources in Fig. 5. For clarity errors are not plotted in this figure given the large number of sources. To see how the hardness ratios relate

to the fitted absorption, we made the symbol sizes proportional to the measured intrinsic column density. Filled squares refer to sources for which the best-fit model is a leaky model. If we do not take into account these latter sources, we can see how the most absorbed sources concentrate in the upper right in the HR3 vs. HR2 figures. For the HBSS, one of the unidentified sources is clearly within that region, whereas the remaining one is not, although it could be moderately absorbed given its photon index from the power-law fit. For the BSS, only about two lie within that region. Evidently also the intrinsic absorption seems to increase as the hardness ratio HR3 increases.

In Fig. 6 the measured column densities are plotted against HR3. Obviously there seems to be a correlation between the amount of intrinsic absorption and HR3 for absorbed sources, again not taking into account sources with a leaky shape. To derive an estimate of the intrinsic column density for unidentified sources, we fitted a linear model to this correlation. In this way, an intrinsic column density can be estimated even when the X-ray data quality is too poor to carry out a reliable spectral analysis. We performed the fit in two different ways. The first one was to fit a linear model by using χ^2 statistics. We considered only those sources with a detected value of the intrinsic column density higher than $4 \times 10^{21} \text{ cm}^{-2}$, i.e. absorbed sources, and we excluded sources with a leaky shape. This selection criterion resulted in a total of 25 sources. With a Spearman rank correlation analysis we confirmed a strong correlation ($\rho=0.82$, probability=0.0001). The resulting fitted relation corresponds to the dashed line in Fig. 6 and it is

$$\text{Log}(N_H) = 22.2(\pm 0.2) + 1.2(\pm 0.2)\text{HR3}. \quad (2)$$

Our second approach was to use all sources with a column density or upper limit above the $4 \times 10^{21} \text{ cm}^{-2}$ threshold and also all unabsorbed sources whose intrinsic column densities values were consistent with this limit within 90% errors, which were 85 sources in total. To perform this analysis, we used the ASURV package (*Astronomy Survival Analysis*, Isobe & Feigelson 1990, which implements the methods presented in Isobe & Feigelson 1986). Applying a Spearman rank analysis, including the upper limits, we found again that there is a strong correlation between the intrinsic column density and HR3 ($\rho=0.84$, probability<0.0001). We performed linear regression with the parametric EM algorithm, solid line in Fig. 6, which assumes Gaussian residuals as in χ^2 statistics, obtaining the relation

$$\text{Log}(N_H) = 22.00(\pm 0.04) + 1.46(\pm 0.10)\text{HR3}. \quad (3)$$

Both relations in Eq. 2 and 3 give similar estimates for the intrinsic column density and can be applied up to redshift ~ 1 , given the energy bands considered in the computation of HR3, and $\text{HR3} > 0$. Making use of these relations, an estimate of the intrinsic column density can be obtained for unidentified sources with small number of collected counts in X-ray surveys. As an example, we found for our unidentified sources that about eight out of the 27 unidentified sources could be absorbed AGN, consistent with what we obtained from the spectral analysis.

5. Photon index

Unfortunately, for half of type 2 AGN (14 out of 29) the photon index was fixed to 1.9 during the spectral fit to better constrain the intrinsic absorption, so we restricted our analysis of the power law index to those type 1 AGN in the XBS AGN

⁵ <http://xmmssc-www.star.le.ac.uk/Catalogue/2XMM/>

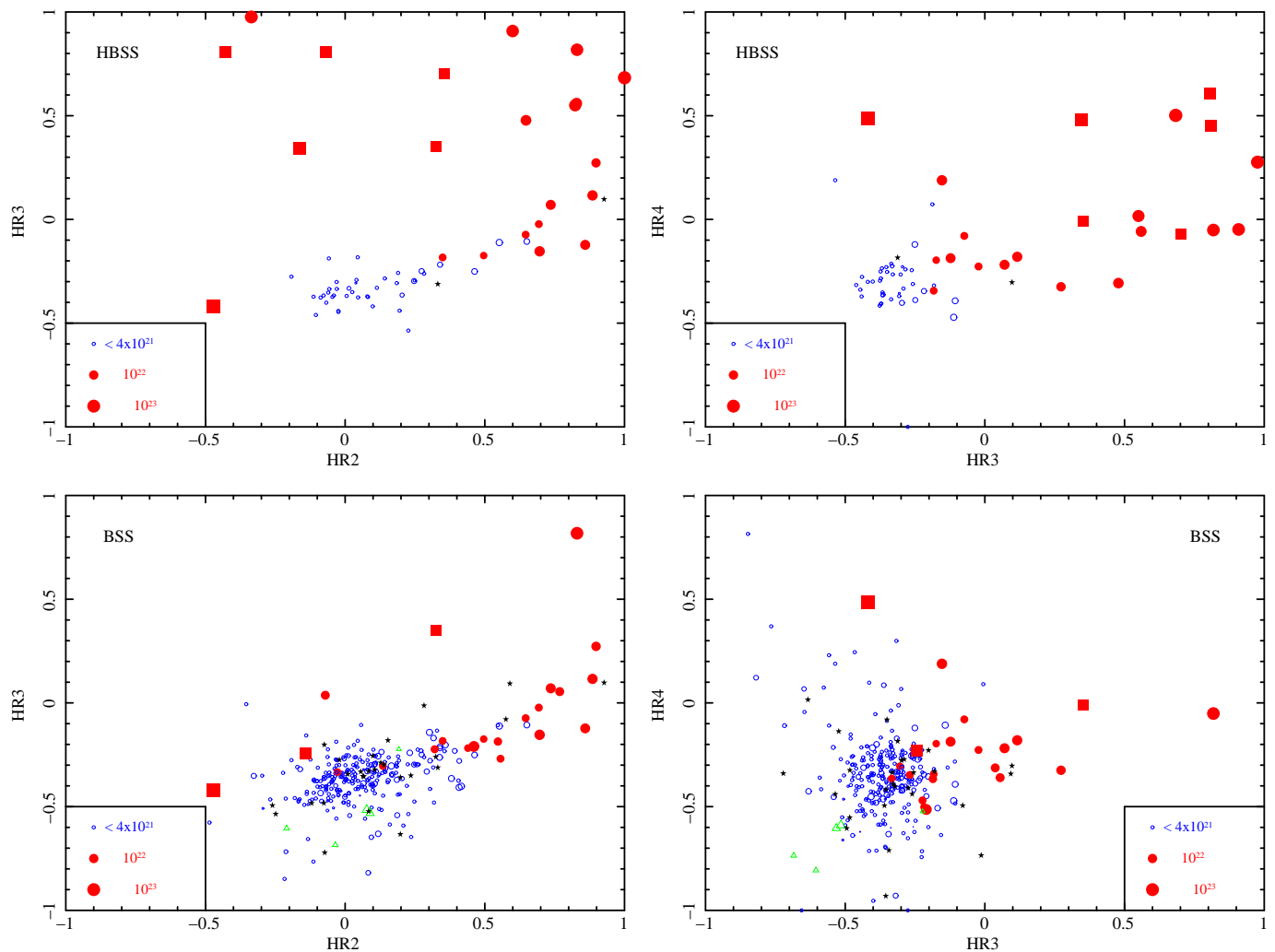


Fig. 5. Hardness ratios for the unabsorbed (open circles) and absorbed (filled circles) identified sources corresponding to the HBSS (top panels) and the BSS (bottom panels), the circle sizes are proportional to the measured intrinsic absorption. Stars correspond to unidentified sources. Filled squares correspond to sources whose best fit is a leaky model.

sample for which we were able to measure Γ , 267 type 1 AGN. We computed its mean and its intrinsic dispersion making use of Maccacaro et al. (1988) likelihood maximization technique. Because the errors on the photon index are not symmetric, we used the average value for the individual errors in each case. We obtain a mean value of $\langle\Gamma\rangle=2.05\pm0.03$ with an intrinsic dispersion of $\sigma=0.26\pm0.02$ for the BSS and $\langle\Gamma\rangle=1.98\pm0.08$ with an intrinsic dispersion of $\sigma=0.29\pm0.05$ for the HBSS, in agreement with previous works (Mateos et al. 2010, Young et al. 2009, Dadina 2008, Mainieri et al. 2007, Page et al. 2006, Mateos et al. 2005a, Mateos et al. 2005b). Errors were extracted from the 1σ confidence contours. The measured power-law index distribution for the BSS and the HBSS is shown in Fig. 7 along with the computed mean and intrinsic dispersion and their confidence contours. There are three type 1 AGN for which the resulting photon index is > 3 . Two of them are NLSy1s (narrow line Seyfert 1 galaxies), known to show these high values for the photon index, and the remaining one is a Seyfert galaxy with a low number of counts in its X-ray spectra.

To check for possible dependence of the photon index on redshift or luminosity, the values for the individual photon in-

dex measurements are presented along with the source redshift and luminosity in Fig. 9. We find an anti correlation between the power-law index and the X-ray luminosity ($\rho=-0.21$, probability= 8×10^{-4}) and marginally redshift ($\rho=-0.10$, probability=0.09) with the Spearman rank correlation analysis. We also find a correlation with the 0.5-2 keV flux ($\rho=0.15$, probability=0.01), which could be caused by undetected intrinsic absorption as pointed out in Mateos et al. (2010). In that work, the authors found a stronger correlation between the photon index and the source redshift, but in their case it mainly started above redshift 2 and our sample only contains 3 sources above that redshift. In our case, we find that the anti-correlation between photon index and luminosity seems to be the strongest one. Moreover, and given that this is a flux-limited sample, the dependence on redshift could be merely caused by the dependence on luminosity. To test this scenario, we selected two narrow luminosity ranges, 10^{43} to 10^{44} erg s $^{-1}$, and 10^{44} to 10^{45} erg s $^{-1}$ (below 10^{43} erg s $^{-1}$ there are not enough sources to perform a reliable analysis), and applied the same correlation analysis as for the whole type 1 AGN sample. For the lowest luminosity range, which reaches only $z \sim 0.8$, we find that the anti-correlations

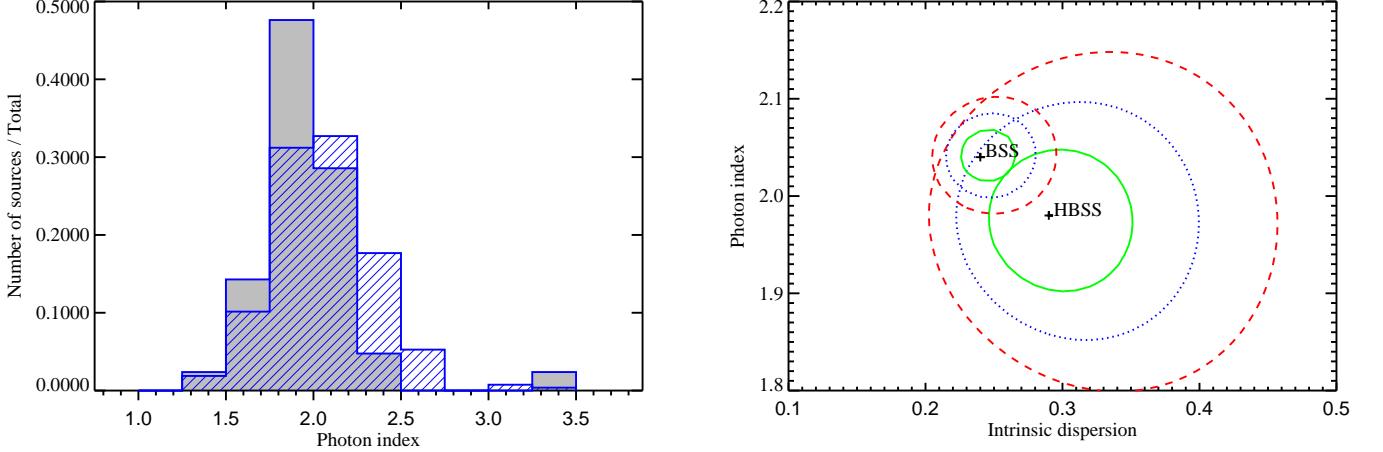


Fig. 7. Left panel: Photon index distribution for the BSS (dashed histogram) and the HBSS (filled histogram). Right panel: Computed mean photon index and intrinsic dispersion for the BSS and the HBSS along with the 1σ (solid line), 2σ (dotted line) and 3σ (dashed line) confidence contours.

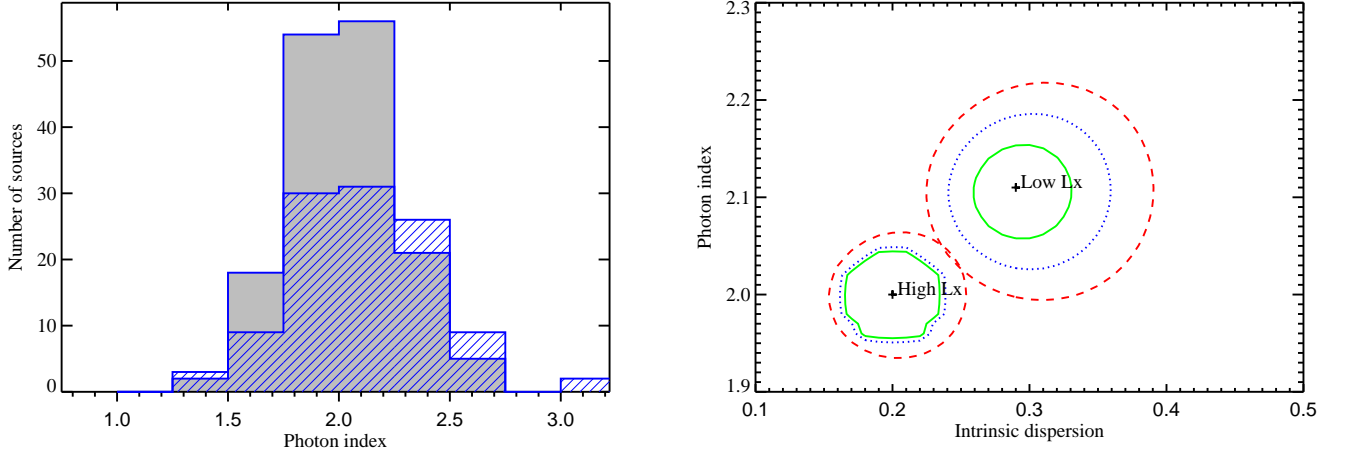


Fig. 8. Left panel: Photon index distribution corresponding to the low-luminosity (dashed histogram) and the high-luminosity (filled histogram) subsamples. Right panel: Computed mean and intrinsic dispersion for the low- and high-luminosity subsamples along with the 1σ (solid line), 2σ (dotted line) and 3σ (dashed line) confidence contours.

turns into a correlation between the photon index and redshift ($\rho = 0.29$, probability = 0.006), while for the high-luminosity bin, which reaches $z \sim 1.5$, we find that the correlation disappears ($\rho = 0.06$, probability = 0.69). This may imply that the observed correlation between the source photon indices and redshifts is mainly driven by an actual correlation between the photon index and the intrinsic luminosity.

To better explore these correlations, we constructed redshift and luminosity bins by dividing the sample into six bins with an equal number of sources (45 sources per bin, 42 in the last bin) and applied the likelihood-maximization technique to each bin. The results are presented in Fig. 9. An anticorrelation between the photon index and redshift and luminosity seems to be present, but it is within the intrinsic dispersion at each redshift or luminosity bin. A similar result is also found in Mateos et al. (2010). These authors pointed out that the hardening of the spectra at higher luminosities and redshifts can be caused by a decrease

in the detection efficiency for softer sources and an increment for harder sources, given that the sample is flux-limited.

To further constrain the correlation of the photon index with the intrinsic luminosity, which seems to be the strongest correlation in our case, we divided the sample into two subsamples, the criterion of which was if the intrinsic luminosity was higher or lower than 10^{44} erg s $^{-1}$. The results are displayed in Fig. 8. We compute a value of $\langle \Gamma \rangle = 2.11 \pm 0.04$ with an intrinsic dispersion of $\sigma = 0.29 \pm 0.04$ for the low-luminosity subsample and $\langle \Gamma \rangle = 2.00 \pm 0.05$ with an intrinsic dispersion of $\sigma = 0.20 \pm 0.04$ for the high-luminosity subsample. The mean and intrinsic dispersion results are different almost at the 3σ confidence level, although the main difference seems to be on the intrinsic dispersion. Applying a Kolmogorov-Smirnov (K-S) test, we find that the probability for both distributions to be drawn from the same parent distribution is only $\sim 1\%$. The opposite is found in Bianchi et al. (2009a), who analyzed high-quality X-ray spectral

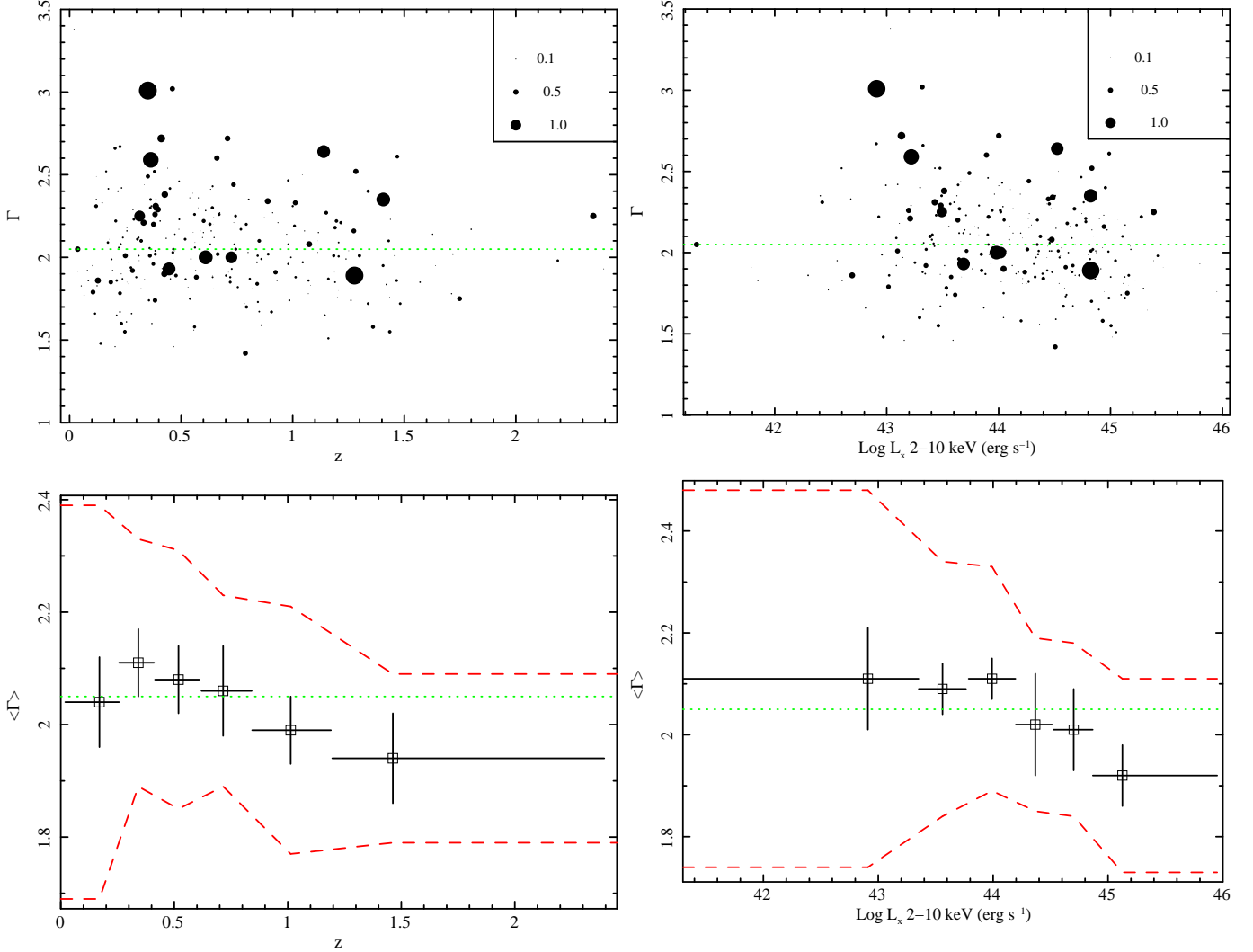


Fig. 9. Top panel: Type 1 AGN photon index versus redshift (left) and luminosity (right). The size of the circles indicates the size of the errors on the photon index. The horizontal dotted line corresponds to the average photon index for the whole sample. Bottom panel: Type 1 AGN average photon index versus redshift (left) and luminosity (right) for the luminosity and redshift bins. The error bars correspond to the mean error at 1σ confidence level, whereas the dashed lines mark the values for the intrinsic dispersion at each bin.

data and found that Seyfert galaxies show a flatter photon index than quasars.

In principle, the observed correlation of Γ with the luminosity could be owing to the presence of an undetected reflection-emission component that becomes increasingly important with the luminosity. In Section 7 we show that this correlation is present also when we analyze the average spectra, where the reflection component is already accounted for. This result excludes that the flattening of Γ with the luminosity is due to the reflection component.

Because the type 1 AGN sample we used in this particular analysis could be contaminated by radio-loud (RL) sources, which are expected to have a flatter photon index on average (Reeves et al. 1997; Reeves & Turner 2000), we performed a safety test. Making use of the NVSS/XBS cross-correlation and analysis presented in Galbiati et al. (2005), we removed all RL sources within these 267 type 1 AGN, which were 14 sources, and applied the same likelihood analysis by dividing into high-

and low-luminosity AGN. We obtained the same result as for the whole sample, the only difference was a small decrement on the intrinsic dispersion for the high-luminosity subsample ($\sigma=0.19\pm0.04$), and the two samples were still different almost at the 3σ level.

6. Soft-excess emission

We say a source shows a soft-excess emission when the extrapolated 2–10 keV power-law fit displays systematic positive residuals at low energies. We find that 35 AGN out of the 41 sources that were not well-fitted by an absorbed power law display a soft-excess. For 29 out of these 35 AGN, we are able to find an additional component that significantly improves the simple power-law fit, as measured by F-test. Assuming a fraction of spurious detections of 0.05, given our F-test significance limit of 95%, this number corresponds to $5^{+2}_{-4}\%$ of the total XBS AGN sample ($14^{+9}_{-7}\%$ for the HBSS and $4^{+2}_{-2}\%$ for the BSS). If we only take

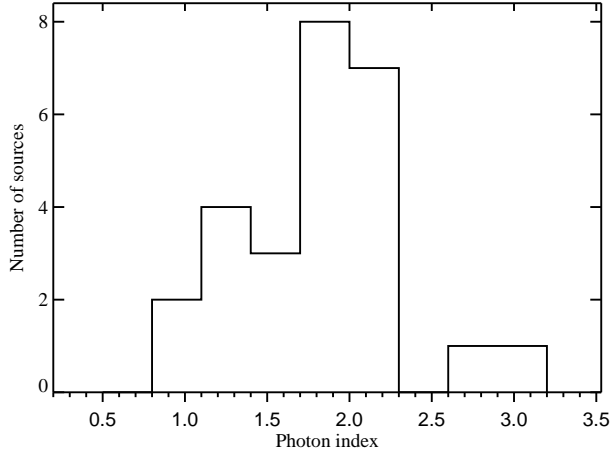


Fig. 4. Photon index distribution for the 27 unidentified sources in the XBS sample.

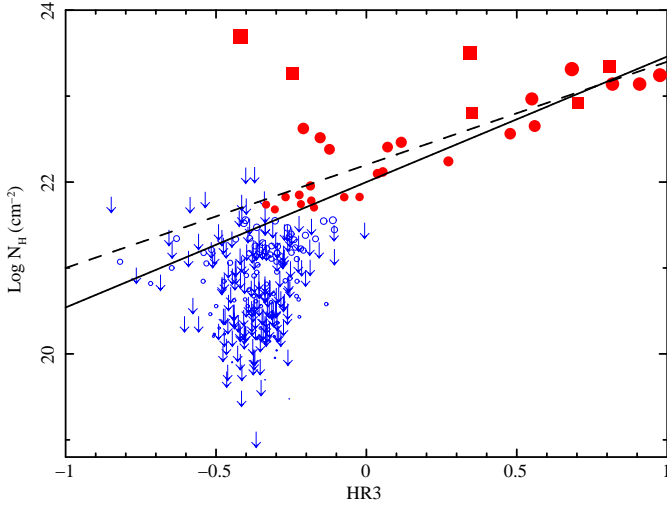


Fig. 6. Intrinsic column density against HR3 for unabsorbed (open circles) and absorbed (filled circles) sources. Filled squares correspond to sources whose best fit is a leaky model. Dashed and solid lines correspond to different fits to the observed correlation (see text for details).

into account sources below $z=0.5$ (beyond that value most of the soft-excess emission is redshifted outside the EPIC energy range) the fraction of sources increases to $11^{+6}_{-5}\%$ for the XBS sample ($20^{+12}_{-10}\%$ for the HBSS and $9^{+6}_{-4}\%$ for the BSS), a value closer to the reported values in recent works (Mateos et al. 2010; Bianchi et al. 2009a), although still lower. It should be noted that our computed value has to be considered as a lower limit because we did not search for soft-excess emission for all our sources, but only for the ones for which a simple power law gives a probability $< 10\%$. This could also explain why the fraction of sources showing soft-excess is larger for the HBSS than for the BSS. The difference in the fraction of detected soft-excess in both samples is likely caused by differences in the data quality. The collected number of counts for the HBSS spectra is larger on average than for the BSS. As the spectral quality increases, it becomes easier to detect and characterize additional components. Undetected soft-excess would also increase the value for the measured photon index, and this in turn could be contributing to increase the

computed average photon index for the BSS and resulting in a higher value than the one for the HBSS.

In the case of absorbed AGN and thanks to Chandra and *XMM-Newton* grating spectra, this soft-excess is known to be associated to scattered emission hundreds of pc far from the central source, likely by the NLR clouds (see for example Bianchi et al. 2006). Indeed, all absorbed AGN that display soft-excess, five type 2 and two type 1 AGN, are best-fitted either by a leaky model or by a power law plus a thermal component that could arise from the host galaxy given its low luminosity.

The case of unabsorbed AGN is more complex. Soft-excess emission has usually been attributed to the hard tail of the thermal emission from the accretion disk or to optically-thick comptonization of EUV disk photons (Ross et al. 1992; Shimura & Takahara 1993), but these models are unable to explain either the higher temperatures usually detected or the fact that these temperatures seem not to vary with AGN properties such as the intrinsic luminosity. A recent model, also invoking continuum emission, explains this soft-excess emission via optically-thin comptonization of the disk photons (Kawaguchi et al. 2001), which would explain the non-dependence on the source luminosity. Two alternative models, based on atomic processes within the accretion disk, have been proposed in recent works: the soft-excess emission could come from relativistically blurred reflection from a partially ionized accretion disk (Crummy et al. 2006) or from velocity-smeared absorption from partially ionized material coming from a disk wind (Middleton et al. 2007), although they are indistinguishable at the EPIC energies. Besides, the quality of our data prevents us from applying them in our spectral analysis. Using the additional components described in Sect.3, we find a great variety of best-fit models within the unabsorbed AGN that show soft-excess in our sample, 27 sources in total: 3 reflection components, 3 ionized absorbers, 13 black body models and two power laws plus an absorption edge. For six of them no acceptable fit was found, therefore the simple power law model was adopted as the best-fit model.

As mentioned before, when the soft-excess in unabsorbed AGN is modeled with a black body model, a value of $kT \sim 0.1$ keV is obtained that does not depend on the source flux, redshift or luminosity. To compare our results with previous works, we also attempted to fit a black body plus a power law to all unabsorbed sources with a soft-excess (see Table 6). This improved the simple power-law fit, F-test $> 95\%$, in all but 6 cases, 21 sources in total. The values obtained for the black body temperature are presented in Fig. 10 against the 2-10 keV luminosity. By using a Spearman rank correlation analysis, we found a significant correlation between the black body temperature and the source's luminosity ($\rho = 0.60$, probability=0.004). This can be because higher luminosity sources are at higher redshifts, given that the sample is flux-limited, which means that the soft-emission is shifted outside the observed energy range. Therefore, at higher luminosities only black body components of higher temperatures can be detected. Indeed, if we remove sources at $z > 0.5$, the correlation disappears ($\rho = 0.06$, probability = 0.78). Therefore, and given the present statistics, we cannot confirm if there is an actual correlation between the black body temperature and the hard X-ray luminosity, although our results suggest that high black body temperatures can only be reached by sources with high intrinsic luminosities.

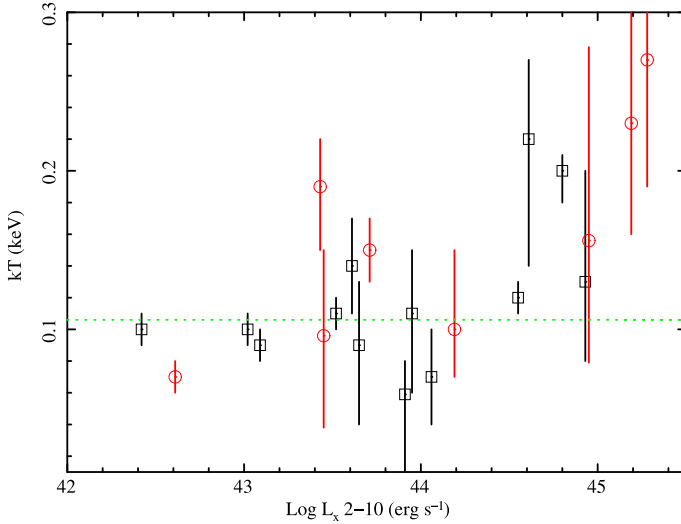


Fig. 10. Black body temperature versus 2-10 keV luminosity for unabsorbed AGN. Squares correspond to the temperature values if a power law plus a black body component is considered the best-fit model, whereas circles correspond to the temperature values in the remaining unabsorbed AGN with a soft-excess emission. The horizontal dotted line corresponds to the weighted mean.

7. Average spectrum

A narrow Fe $K\alpha$ emission line (6.4-6.9 keV, depending on ionization state of the material) is almost ubiquitously observed in good quality X-ray spectra of AGN (Nandra et al. 2007; Bianchi et al. 2009b). To measure a broad and/or a relativistically broadened component requires an even better quality (Guainazzi et al. 2006). We searched for this emission line in our sample, but its presence is only suggested in ~ 20 sources and its parameters are only well constrained in a couple of cases, all corresponding to a narrow Fe $K\alpha$ line. To improve the SNR and to detect spectral features that would remain hidden otherwise, such as the Fe $K\alpha$ line, we averaged the whole sample following the process described in Corral et al. (2008). The basic steps of this averaging method are

- Selecting those individual spectra, pn or MOS, with more than 80 counts in the total 0.2-12 keV EPIC band. For the XBS AGN sample, this means to exclude 12 sources, 8 type 1 and 4 type 2 AGN.
- Fitting pn and MOS spectra for each source individually by using a simple absorbed power-law model and leaving the power-law index, intrinsic column density, and the normalization as free parameters.
- Obtaining the incident spectra, i. e. before entering the detectors, in flux units ($\text{keV cm}^{-2} \text{s}^{-1} \text{keV}^{-1}$) by using the parameters from the previous spectral fit.
- Correcting for the absorption from our Galaxy and shifting to rest frame.
- Rescaling the individual spectra so that every spectrum has the same 2-5 keV rest frame flux.
- Binning every spectrum to a common energy grid so that the final averaged spectrum has at least 1000 counts per bin.
- Averaging by using a standard mean.

As a final step and to quantify the significance of any spectral feature, we used simulations: we simulated each source 100 times by using the fitted model and keeping the same spectral

quality as for the real data. By averaging all the simulations we obtained a simulated “continuum” that should account for the average of absorbed power laws. Taking one simulation for each real spectrum and averaging these, we constructed 100 simulated continua from which we can compute 1σ and 2σ limits by removing the 32 and 5 extreme values at each bin. In this way, we can say that any excursion over or below these limits is detected at 1σ or 2σ confidence level. By using the simulated “continuum” and the confidence limit and comparing them to our data, we can determine if there are any significant deviations from a power-law shape and estimate their significance. Only the energies in the 2 to 15 keV rest-frame energy band are used in this spectral analysis. For energies below 2 keV, the averaging method is highly dependent on the model used to unfold the spectra, and for energies above 15 keV noise becomes too important. The resulting best-fit models and parameters for each case are shown in Table 5.

First, and to compare the two samples under study here, we constructed the average spectra and confidence limits for the type 1 and type 2 AGN within the BSS and HBSS separately. The resulting averaged spectrum, simulated continuum, and confidence limits for the BSS and HBSS samples are shown in Fig. 11.

BSS: For the type 1 AGN we find that the best-fit model consists of a power law of $\Gamma=2.02^{+0.04}_{-0.03}$ plus a narrow Fe $K\alpha$ line centered at $E=6.40^{+0.04}_{-0.06}$ keV and equivalent width $EW=110^{+30}_{-40}$ eV and a reflection component with a reflection factor $R\sim 0.6$. The inclination angle of the reflection component is always fixed to its default value, $i\sim 60$ deg, because it cannot be determined at the same time as the reflection factor. For the type 2 AGN the best-fit model turns out to be an absorbed power law (Γ fixed to 1.9) with an intrinsic column density $N_H < 1.1 \times 10^{22} \text{cm}^{-2}$ plus a narrow emission line ($E=6.53^{+0.14}_{-0.04}$ keV, $EW=200^{+150}_{-150}$ eV) and a reflection component ($R\sim 1$). The fitted intrinsic column density does not represent the actual average absorption of the sources, but it is related to the fraction of absorbed sources among the type 2 AGN as well as their column densities, and as such does not have a meaningful physical interpretation. In any case, and given that not only the statistics are much lower for type 2 AGN but also that we are fitting above 2 keV, the column density cannot be very well constrained.

HBSS: For the type 1 AGN we find that the best-fit model consists of a power law of $\Gamma=2.00^{+0.05}_{-0.07}$ plus a narrow Fe $K\alpha$ line centered at $E=6.44^{+0.05}_{-0.04}$ keV and equivalent width $EW=80^{+60}_{-40}$ eV and a reflection component with a reflection factor $R\sim 0.9$. For the type 2 AGN the best-fit model turns out to be an absorbed power law (Γ fixed to 1.9) with an intrinsic column density $N_H \sim 2 \times 10^{22} \text{cm}^{-2}$ plus a narrow emission line ($E=6.42^{+0.09}_{-0.13}$ keV, $EW=90^{+50}_{-60}$ eV) and a reflection component ($R\sim 1$).

The small differences between type 2 AGN average spectra for the BSS and HBSS are caused by the larger number of absorbed sources for the HBSS. This is expected because selecting at harder energies makes the sample less biased against absorbed sources. The relatively low value for the average column density for the type 2 AGN in the BSS is consistent with it being due to contribution of unabsorbed type 2 AGN, which are more numerous in the BSS (which contains six unabsorbed type 2 AGN out of 19 type 2 AGN) than in the HBSS (which only contains one unabsorbed type 2 AGN out of 20). For both samples and AGN classes, the detected Fe $K\alpha$ line turns out to be narrow and likely comes from neutral material, i.e., far from the central source.

To better characterize the differences in the spectral shape from absorbed to unabsorbed sources, we divided the whole

sample into absorbed (31 AGN) and unabsorbed (274 AGN) sources and constructed the average spectra. The results are shown in Fig. 12.

For the **unabsorbed sources** we find that the best-fit model consists on a power law of $\Gamma = 2.10^{+0.03}_{-0.03}$ plus a narrow Fe K α line centered at $E = 6.40^{+0.04}_{-0.06}$ keV and equivalent width $EW = 100^{+30}_{-40}$ eV and a reflection component with a reflection factor $R \sim 0.5$. A relativistic Fe K α line is not clearly present and when we attempted to fit one, we did not obtain a significant improvement, and the line parameters resulted in unphysical values (such as a very large inclination angle, $i \sim 80^\circ$, or a extremely high line energy ~ 8 keV). From the best-fit model, we estimate an upper limit for the EW of a relativistic Fe K α line contribution of 230 eV at the 3σ confidence level for a relativistic emission line (Laor 1991) centered at 6.4 keV and inclination angle of 30° . For the **absorbed sources** the best-fit model turns out to be an absorbed power law (Γ fixed to 1.9) with an intrinsic column density $N_H \sim 1 \times 10^{22} \text{ cm}^{-2}$ plus a narrow emission line ($E = 6.47^{+0.06}_{-0.07}$ keV, $EW = 100^{+70}_{-50}$ eV) and a reflection component ($R \sim 1.5$). The main difference between absorbed and unabsorbed sources seems to be a larger amount of reflection in the case of absorbed AGN besides the amount of absorption. This difference is not due to a hidden dependence of the reflection with luminosity because absorbed and unabsorbed AGN in the XBS sample display a very similar luminosity distribution with an average luminosity of $\langle L_X(2-10 \text{ keV}) \rangle \sim 4 \times 10^{44}$ in both cases and about the same dispersion. The averaging process used here is designed to study the 2 to 10 keV rest-frame range to study the Fe K α line properties that minimize the contribution of the underlying continuum and observational effects. However, for highly absorbed sources ($N_H > 10^{23} \text{ cm}^{-2}$), the way the rescaling is carried out can give larger weights during the averaging process to the more absorbed sources. In the final average spectra, this produces a feature that could mimic the shape of a reflection component (Corral et al. 2011, in preparation). However, the number of highly absorbed sources is too small to be responsible for the whole amount of observed reflection. As a safety test, we removed the eight most absorbed AGN ($N_H > 10^{23} \text{ cm}^{-2}$) from the average of absorbed AGN. Given that the remaining number of sources is small in this case, R cannot be well constrained, but we obtain a lower limit of $R > 1.1$ at the 90% confidence level. In summary, although the values obtained for the reflection components reported here have to be taken as tentative, the difference in the amount of reflection between absorbed and unabsorbed AGN seems to be real. For unabsorbed sources, our results excellently agree with those from studies of local AGN (Nandra et al. 2007), from the average of large samples of distant AGN (Chaudhary & Brusa 2010) and with the predictions of theoretical models (Ballantyne 2010).

7.1. Dependence on redshift and luminosity

As we showed in Sect. 5, there seems to be a difference between the spectral shape for low- ($L_X < 10^{44} \text{ erg s}^{-1}$) and high- ($L_X > 10^{44} \text{ erg s}^{-1}$) luminosity type 1 AGN. To explore this possible difference, we constructed the average spectrum for both luminosity subsamples. The resulting ratios of the average spectra to the simulated continua are shown in Fig. 13. We find that the best-fit model for both samples consists in a power law with a narrow Fe K α line centered on ~ 6.4 keV and a reflection component. Consistently with the results reported in Sect. 5 we find the photon index of the average spectrum for the low-luminosity subsample to be (marginally) larger than the one for

the high-luminosity subsample. The large error on Γ in the low-luminosity subsample is likely caused by the larger dispersion of the photon index distribution (see Fig. 8). For the low-luminosity subsample, the line EW seems to be higher, $EW = 110 \pm 30$ eV, than for the high-luminosity subsample, $EW = 80 \pm 30$ eV (the so-called Iwasawa-Taniguchi effect, Iwasawa & Taniguchi 1993), but both values are consistent within errors. The resulting reflection fraction also turns out to be marginally larger for the low-luminosity subsample, $R = 0.8^{+0.8}_{-0.5}$, than for the high-luminosity sample, $R = 0.3 \pm 0.1$, although it is not well constrained for the former sample. These results agree with models that predicts a decrease of the torus covering fraction as the luminosity increases (Lawrence 1991), thus decreasing the reprocessing of the radiation within the torus and also explaining the Iwasawa-Taniguchi effect. Evidence of this decrement of the covering fraction as a function of the luminosity have been reported in recent works (Maiolino et al. 2007, Della Ceca et al. 2008), which point out the need for the simplest unified schemes to be revised.

We also explored the possible dependence of the spectral shape on redshift. To this end, we again constructed average spectra by dividing the sample in different redshift bins. In this case we did not detect any significant trend of the resulting averaged spectral shape with redshift. Nevertheless, we point out again that our sample only reaches redshift ~ 2 , and it is above this value where, for example, Mateos et al. (2010) found this dependence to become stronger.

8. Discussion

In the previous sections we discussed the possible existence of several statistical correlations. First, we found an anti-correlation between the photon index and the X-ray luminosity. This correlation is significant in the analysis of the single spectra, but it is also marginally present in the analysis of the average spectra. The lower significance in the latter case is likely caused by the high intrinsic dispersion in the photon index distribution. The second correlation, found in the analysis of the average spectra, is the dependence of the intensity of the reflection component with both the AGN “type” and luminosity. In particular, the reflection component seems to be stronger in absorbed AGN and in low-luminosity AGN.

An anti-correlation between the photon index and the X-ray luminosity has been recently reported by Green et al. (2009) and Young et al. (2009), whereas Saez et al. (2008) found a “positive” correlation. It has to be noted that Saez et al. (2008) took into account type 1 and type 2 AGN at the same time, the latter being more numerous than the former, whereas we here only considered type 1 AGN. The physical explanation for this anti-correlation is still a matter of debate in the recent literature. Several authors (Shemmer et al. 2008; Risaliti et al. 2009; Grupe et al. 2010) have reported a correlation between the photon index and the Eddington ratio. This dependence could explain the correlations we find if the Eddington ratio was the actual driver of the luminosity/photon index anti-correlation. In that case, the low-luminosity subsample could be sampling a different AGN population or a mixture of very different accretion states, which could explain the higher dispersion found for the photon indices. We are currently studying this hypothesis in deeper detail.

The observed anti-correlation between the reflection component intensity and the luminosity found in the analysis of the average spectra confirms similar trends that were already observed in other samples (e.g. Nandra et al. 1997). If the observed reflection component is related to the molecular torus, the trend

Table 5. Average spectra fit results.

Sample	N_H 10^{22} (cm^{-2})	Γ	R	E (keV)	σ (eV)	EW (eV)	$\chi^2/\text{d.o.f}$
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)
BSS type 1 AGN	...	$2.02^{+0.04}_{-0.03}$	$0.6^{+0.2}_{-0.2}$	$6.40^{+0.04}_{-0.05}$	< 160	110^{+30}_{-40}	21/16
HBSS type 1 AGN	...	$2.00^{+0.03}_{-0.07}$	$0.9^{+0.2}_{-0.2}$	$6.44^{+0.05}_{-0.04}$	< 140	80^{+60}_{-40}	15/16
BSS type 2 AGN	< 1.1	1.9^f	$1.0^{+0.7}_{-0.8}$	$6.53^{+0.14}_{-0.06}$	< 200	200^{+150}_{-150}	13/16
HBSS type 2 AGN	$2.0^{+0.5}_{-0.4}$	1.9^f	$1.0^{+0.3}_{-0.3}$	$6.42^{+0.09}_{-0.13}$	< 250	90^{+50}_{-60}	6/16
Unabsorbed AGN	...	$2.10^{+0.03}_{-0.03}$	$0.5^{+0.1}_{-0.1}$	$6.40^{+0.04}_{-0.06}$	< 160	100^{+30}_{-40}	18/16
Absorbed AGN	$1.2^{+0.4}_{-0.3}$	1.9^f	$1.5^{+0.2}_{-0.3}$	$6.47^{+0.08}_{-0.07}$	< 160	100^{+40}_{-50}	9/16
Low-luminosity type 1 AGN	...	$2.11^{+0.10}_{-0.20}$	$0.8^{+0.8}_{-0.5}$	$6.43^{+0.06}_{-0.13}$	< 180	110^{+30}_{-30}	12/16
High-luminosity type 1 AGN	...	$2.00^{+0.03}_{-0.03}$	$0.3^{+0.1}_{-0.1}$	$6.39^{+0.04}_{-0.05}$	< 130	80^{+30}_{-30}	25/16

Columns: (1) Sample used to construct the average spectrum; (2) Intrinsic column density; (3) Photon index; (4) Reflection scaling factor; (5) Fe $K\alpha$ central energy; (6) Fe $K\alpha$ width upper limit; (7) Fe $K\alpha$ equivalent width; (8) χ^2 to number of degrees of freedom.

f : Fixed parameter.

can be well explained in the context of the receding torus model according to which the molecular torus-covering fraction (and thus the intensity of the reflected component) decreases with the luminosity.

On the contrary, the difference seen in the observed reflection component between absorbed and unabsorbed AGN is, at the moment, troublesome because it is not easily reconciled with the unified model (which predicts a larger reflection component for unobscured sources, see e.g. Krolik et al. 1994; Murphy & Yaqoob 2009). It is worth noting that similar results as those reported here (i.e. more reflection in absorbed objects) were obtained by Malizia et al. (2003) and by Deluit & Courvoisier (2003) from the analysis of the average spectra of local type 1 and type 2 AGN observed with BeppoSAX; yet Burlon et al. (2011) have recently obtained the opposite trend from the analysis of the type 1 and type 2 AGN observed in the SWIFTBAT, although with large uncertainties. Further detailed studies on this particular aspect are clearly needed.

9. Conclusions

We have analyzed the X-ray spectra corresponding to all identified AGN within the XBS sample.

From the individual analysis and according to our fitting criteria, we find that

- Most AGN are well fitted by a simple unabsorbed power law model. The most common deviation from this shape are neutral intrinsic absorption and soft-excess emission.
- In agreement with the AGN unified model, most type 2 AGN are absorbed ($N_H > 4 \times 10^{21} \text{ cm}^{-2}$) and by larger amounts of intrinsic material than type 1 AGN, which are most unabsorbed. Nonetheless, deviations from this simple version of the unified model are found, and are more frequent in type 2 AGN.
- We find that the fraction of exceptions to the unified model is of 5% for the whole sample (only 3% if we take into account the errors on the measured intrinsic absorption). The fraction of type 1 AGN among absorbed sources is 17% and 31% for the HBSS and the BSS, respectively. The different values for the two samples are likely due to the larger number of unidentified sources within the BSS. The fraction of type 2

AGN among unabsorbed sources turns to be 3% and 1% for the HBSS and the BSS, respectively.

- We find that the X-ray spectral photon index for type 1 AGN is anti-correlated with the hard X-ray luminosity. When the type 1 AGN sample is split into high- and low-luminosity subsamples, we find that the intrinsic photon index for both samples is different almost at the 3σ level in the plane photon index vs. intrinsic dispersion. We compute an average photon index and intrinsic dispersion of $\langle \Gamma \rangle = 2.11 \pm 0.04$ (2.00 ± 0.05) and $\sigma = 0.29 \pm 0.04$ (0.20 ± 0.04) for the low- (high) luminosity sample.
- We find that the so-called “soft-excess” is a common characteristic of AGN and it clearly displays different properties and origin for unabsorbed and absorbed AGN.

From the constructed average spectra we find that

- The average spectrum for type 2 AGN is different in the HBSS and the BSS samples as a result of a larger amount of absorbed sources in the HBSS. We do not find any significant differences between the type 1 AGN average spectra for these two samples.
- Apart from the amount of absorption, the differences between average spectra corresponding to absorbed and unabsorbed AGN are caused by an increase in the amount of reflection.
- We do not detect a significant relativistic Fe $K\alpha$ emission line on the average spectrum for unabsorbed sources. We estimate an upper limit for a broad relativistic contribution to the line of 230 eV at the 3σ confidence level.
- When dividing the type 1 AGN sample into high- and low-luminosity sources, we find that the narrow Fe $K\alpha$ line EW seems to decrease as the luminosity increases, which is consistent with the so-called Iwasawa-Taniguchi effect, although the resulting values for the high- and low-luminosity subsamples are consistent within errors (EW = 110 ± 30 and 80 ± 30 eV for the low- and high-luminosity subsamples, respectively). We find moreover that the amount of reflection may also decrease with luminosity, which supports models in which the covering fraction of the putative torus decreases as the intrinsic luminosity increases.

Acknowledgements. We would like to thank the referee for providing us with constructive comments and suggestions. We acknowledge financial support from ASI (grant n.1/088/06/0 and COFIS contract. It is a pleasure to thank Tommaso

Maccacaro, Mike Watson and Valentina Braito for their initial efforts and contributions to the XBS project.

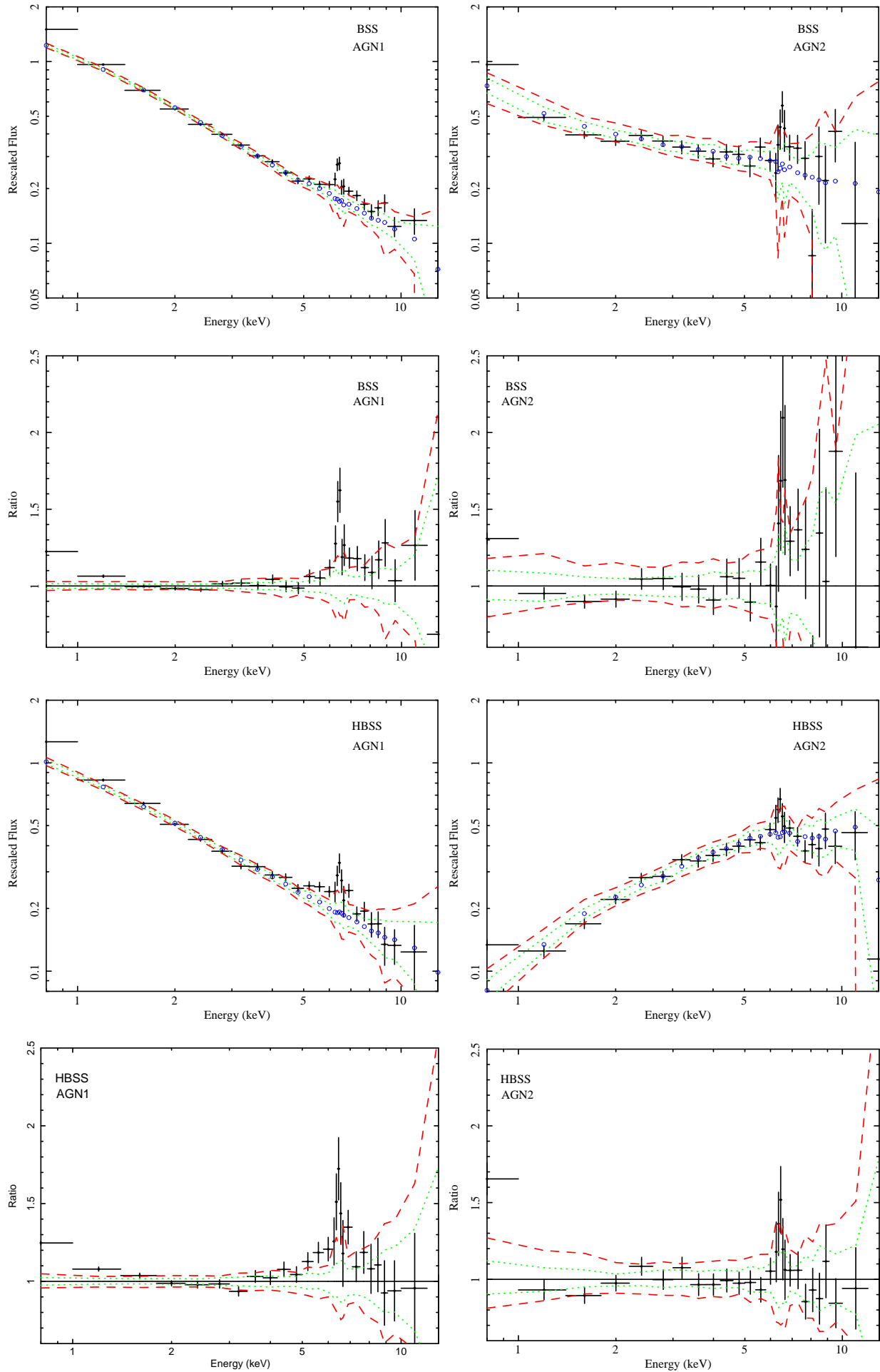


Fig. 11. Type 1 (left) and type 2 (right) AGN average spectrum and average spectrum to simulated continuum ratio (bottom) corresponding to the HBSS and BSS samples. Error bars: real average spectrum, circles: average continuum, dashed line: 2σ limit, dotted line: 1σ limit.

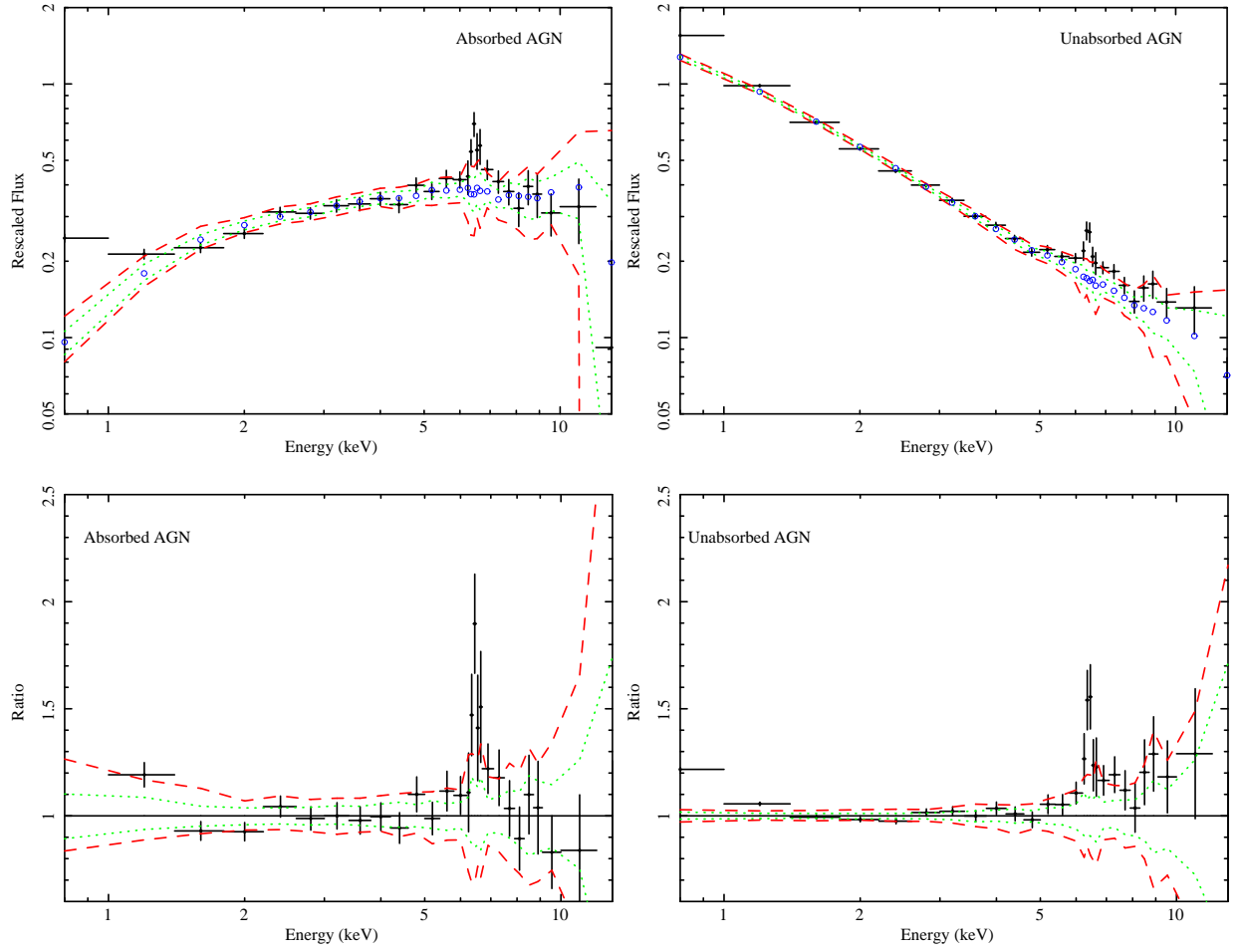


Fig. 12. Average spectrum (top) and average spectrum to simulated continuum ratio (bottom) corresponding to the absorbed (left) and unabsorbed (right) sources. Error bars: real average spectrum, circles: average continuum, dashed line: 2σ limit, dotted line: 1σ limit.

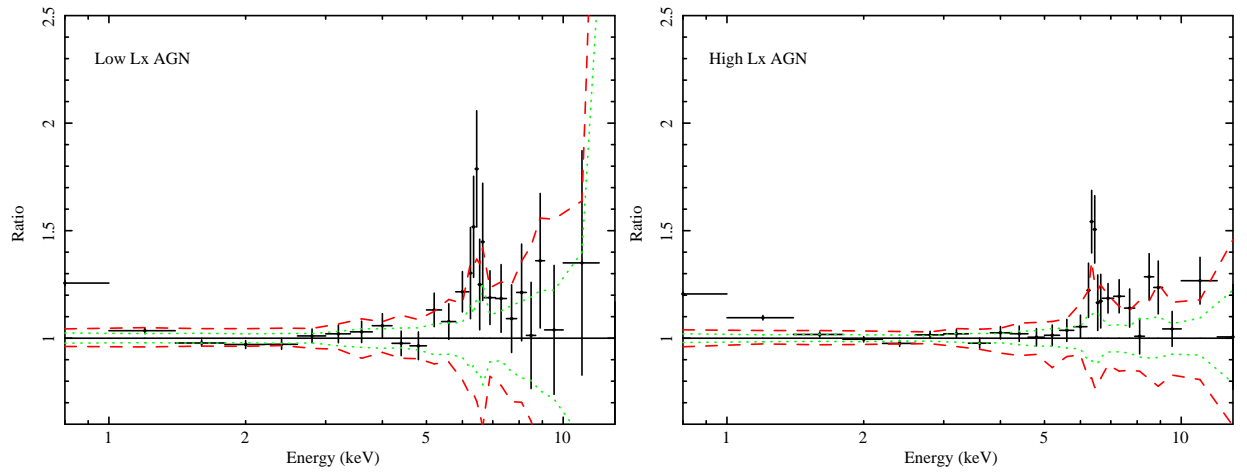


Fig. 13. Average spectrum to simulated continuum ratio corresponding to the luminosity subsamples. Error bars: real average spectrum, dashed lines: 2σ limit, dotted lines: 1σ limit.

Table 6. Black body model fit results.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	kT keV	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{ s}^{-1}$	Log $L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability (%)	Sample
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ000532.7+200716	AGN1	0.119	2.31 ^{+0.38 -0.33}	0.02 ^{+0.22 -0.02}	0.10 ^{+0.01 -0.01}	0.60	42.42	36/35	40.7	BSS
XBSJ005031.1-520012	AGN1	0.463	2.03 ^{+0.35 -0.19}	<0.26	0.11 ^{+0.04 -0.05}	0.94	43.95	88/78	19.7	BSS
XBSJ015957.5+003309	AGN1	0.310	2.01 ^{+0.21 -0.07}	<0.15	0.07 ^{+0.03 -0.01}	3.18	44.06	93/80	15.3	HBSS,BSS
XBSJ021808.3-045845	AGN1	0.712	1.91 ^{+0.07 -0.05}	<0.03	0.20 ^{+0.02 -0.02}	2.58	44.81	446/417	15.9	HBSS,BSS
XBSJ023530.2-523045	AGN1	0.429	1.9 ^{+0.12 -0.14}	<0.125	0.059 ^{+0.021 -0.053}	1.10	43.91	22/18	23.1	BSS
XBSJ023713.5-522734	AGN1	0.193	1.92 ^{+0.12 -0.14}	<0.06	0.11 ^{+0.01 -0.01}	2.71	43.52	71/87	88.9	HBSS,BSS
XBSJ030641.0-283559	AGN1	0.367	2.10 ^{+0.47 -0.39}	<0.14	0.096 ^{+0.054 -0.058}	0.51	43.45	33/24	11.0	BSS
XBSJ031311.7-765428	AGN1	1.274	1.78 ^{+0.17 -0.16}	<0.66	0.156 ^{+0.122 -0.077}	0.98	44.95	59/50	18.8	BSS
XBSJ031851.9-441815	AGN1	1.360	1.58 ^{+0.49 -0.26}	0.13 ^{+1.51 -0.13}	0.13 ^{+0.07 -0.05}	0.94	44.93	17/15	32.8	BSS
XBSJ052108.5-251913	AGN1	1.196	1.70 ^{+0.10 -0.20}	<0.28	0.27 ^{+0.05 -0.08}	2.51	45.28	78/64	11.8	HBSS,BSS
XBSJ065839.5-560813	AGN1	0.211	1.86 ^{+0.20 -0.24}	0.04 ^{+0.08 -0.04}	0.10 ^{+0.01 -0.01}	0.71	43.02	66/68	54.1	BSS
XBSJ074312.1+742937	AGN1	0.312	1.98 ^{+0.07 -0.07}	<0.01	0.12 ^{+0.01 -0.01}	9.79	44.55	190/192	51.7	HBSS,BSS
XBSJ095218.9-013643 ^p	AGN1	0.01993	1.90 ^f	4.2 ^{+4.4 -1.2}	0.15 ^{+0.01 -0.01}	5.73	41.92	133/57	5e-6	HBSS,BSS
XBSJ100926.5+533426	AGN1	1.718	1.72 ^{+0.20 -0.27}	0.17 ^{+0.60 -0.17}	0.23 ^{+0.10 -0.07}	0.87	45.19	41/37	30.7	BSS
XBSJ101922.6+412049	AGN1	0.239	1.51 ^{+0.15 -0.11}	<0.04	0.15 ^{+0.02 -0.02}	2.87	43.71	107/110	57.8	HBSS,BSS
XBSJ140100.0-110942	AGN1 ^e	0.164	2.31 ^{+0.35 -0.23}	0.12 ^{+0.17 -0.03}	0.07 ^{+0.01 -0.01}	0.45	42.61	62/61	46.1	BSS
XBSJ141736.3+523028	AGN1	0.985	1.74 ^{+0.21 -0.16}	<0.29	0.22 ^{+0.05 -0.08}	0.82	44.61	111/97	17.0	BSS
XBSJ153456.1+013033^p	AGN1	0.310	2.27 ^{+0.42 -0.24}	<0.30	0.09 ^{+0.04 -0.03}	1.13	43.65	77/62	9.7	BSS
XBSJ223547.9-255836	AGN1	0.304	1.43 ^{+0.23 -0.34}	<0.03	0.19 ^{+0.05 -0.04}	0.91	43.43	32/32	49.5	BSS
XBSJ225118.0-175951	AGN1	0.172	2.09 ^{+0.27 -0.20}	0.16 ^{+0.17 -0.09}	0.09 ^{+0.01 -0.01}	1.25	43.09	95/88	28.9	BSS
XBSJ231601.7-424038	AGN1	0.383	1.74 ^{+0.42 -0.43}	<0.10	0.14 ^{+0.03 -0.03}	0.75	43.61	56/51	29.0	BSS

Columns: (1) Source name; (2) Class; (3) Redshift; (4) Photon index; (5) Intrinsic column density; (6) Black body temperature; (7) Observed flux in the 2-10 keV band, de-absorbed by our Galaxy; (8) Intrinsic luminosity in the 2-10 keV band; (9) χ^2 to number of degrees of freedom; (10) Null hypothesis probability; (11) Sample the source belongs to.

^p: Null hypothesis probability < 10%. ^e: Elusive AGN. ^f: Fixed parameter. Source and parameters in bold face indicates that this model is considered as our best-fit model. Note: Errors and upper limits are at 90% confidence level. Fluxes and luminosities refer to the MOS2 calibration,

Table 7. Ionized absorbed power law fit results.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$N_{\text{H}i}$ 10^{22} cm^{-2}	ξ	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{ s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability (%)	Sample
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)
XBSJ005031.1–520012	AGN1	0.463	$2.30^{+0.23}_{-0.11}$	<0.049	$0.60^{+1.73}_{-0.45}$	300^{+2300}_{-250}	0.79	43.89	91/78	14.8	BSS
XBSJ015957.5+003309 ^p	AGN1	0.310	$2.30^{+0.14}_{-0.12}$	<0.025	$0.69^{+4.67}_{-0.29}$	160^{+2100}_{-110}	2.79	44.03	92/80	8.1	HBSS,BSS
XBSJ023713.5–522734	AGN1	0.193	$2.34^{+0.08}_{-0.09}$	<0.016	$2.34^{+0.83}_{-1.01}$	900^{+4100}_{-500}	2.35	43.52	80/87	69.3	HBSS,BSS
XBSJ030641.0–283559	AGN1	0.367	$2.31^{+0.23}_{-0.11}$	<0.06	$11.90^{+11.40}_{-7.14}$	3884^{+1116}_{-1744}	0.57	43.60	27/24	29.2	BSS
XBSJ052543.6–334856	AGN1	0.735	$2.44^{+0.43}_{-0.45}$	$0.14^{+0.17}_{-0.14}$	$6.46^{+5.14}_{-4.33}$	1561^{+660}_{-1270}	0.47	44.27	19/17	34.8	BSS
XBSJ065839.5–560813	AGN1	0.211	$2.85^{+0.18}_{-0.13}$	<0.016	$5.0^{+2.0}_{-2.0}$	1600^{+400}_{-300}	0.52	42.94	76/67	21.6	BSS
XBSJ074312.1+742937	AGN1	0.312	$2.23^{+0.03}_{-0.03}$	<0.006	$3.2^{+1.3}_{-1.3}$	1300^{+500}_{-300}	9.13	44.58	206/192	23.22	HBSS,BSS
XBSJ101922.6+412049	AGN1	0.239	$1.82^{+0.09}_{-0.04}$	<0.02	$8.3^{+9.3}_{-3.3}$	1700^{+1600}_{-800}	2.72	43.70	109/110	71.1	HBSS,BSS
XBSJ140100.0–110942	AGN1	0.164	$2.52^{+0.28}_{-0.11}$	<0.03	$0.43^{+0.35}_{-0.18}$	60^{+130}_{-50}	0.41	42.60	63/61	41.1	BSS
XBSJ140127.7+025605	AGN1	0.265	$1.84^{+0.10}_{-0.09}$	$0.17^{+0.04}_{-0.03}$	$0.60^{+0.38}_{-0.27}$	38^{+51}_{-30}	6.81	44.22	381/359	22.2	HBSS,BSS
XBSJ223547.9–255836	AGN1	0.304	$2.11^{+0.25}_{-0.11}$	$0.08^{+0.08}_{-0.02}$	$15.60^{+8.64}_{-1.1}$	2772^{+2061}_{-1321}	0.78	43.41	31/32	50.7	BSS
XBSJ225118.0–175951	AGN1	0.1718	$2.89^{+0.11}_{-0.24}$	$0.05^{+0.02}_{-0.01}$	$4.3^{+1.3}_{-1.1}$	4900^{+100}_{-2700}	1.07	43.14	104/88	12.4	BSS

Columns: (1) Source name; (2) Class; (3) Redshift; (4) Photon index; (5) Intrinsic column density; (6) Ionized absorber column density; (7) Ionization parameter; (8) Observed flux in the 2-10 keV band, de-absorbed by our Galaxy; (9) Intrinsic luminosity in the 2-10 keV band; (10) χ^2 to number of degrees of freedom; (11) Null hypothesis probability; (12) Sample the source belongs to.

^p: Null hypothesis probability < 10%. ^e: Elusive AGN. ^f: Fixed parameter. Source and parameters in bold face indicates that this model is considered as our best-fit model. Note: Errors and upper limits are at 90% confidence level. Fluxes and luminosities refer to the MOS2 calibration.

Table 8. Power law plus reflected component fit results.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	R	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{ s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability (%)	Sample
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ021808.3–045845 ^p	AGN1	0.712	$2.38^{+0.19}_{-0.04}$	<0.04	0.3(<0.5)	2.67	44.80	459/417	7.8	HBSS,BSS
XBSJ023713.5–522734	AGN1	0.193	$2.80^{+0.08}_{-0.15}$	<0.017	2(<7)	2.95	43.53	75/88	84.4	HBSS,BSS
XBSJ031311.7–765428	AGN1	1.274	$2.16^{+0.75}_{-0.16}$	<0.25	<0.10	1.11	44.94	60/50	17.2	BSS
XBSJ043448.3–775329	AGN1^e	0.097	1.9^f	$0.20^{+0.27}_{-0.20}$	0.4(<0.6)	4.00	43.00	9/11	62.1	BSS
XBSJ052108.5–251913	AGN1	1.196	$2.22^{+0.51}_{-0.15}$	<0.17	0.6(<1.20)	2.73	45.28	78/64	11.2	HBSS,BSS
XBSJ074312.1+742937	AGN1	0.312	$2.51^{+0.06}_{-0.06}$	<0.006	1.3(<1.6)	10.6	44.56	199/192	35.2	HBSS,BSS
XBSJ101922.6+412049	AGN1	0.239	$2.12^{+0.27}_{-0.08}$	<0.03	0.3(< 0.6)	2.97	43.65	102/110	69.0	HBSS,BSS
XBSJ141736.3+523028	AGN1	0.985	$2.16^{+0.60}_{-0.11}$	<0.14	0.3(<1.7)	0.89	44.60	112/97	15.0	BSS
XBSJ223547.9–255836	AGN1	0.304	$2.88^{+0.68}_{-0.69}$	$0.08^{+0.09}_{-0.08}$	1.7(<5)	0.86	43.42	32/32	47.7	BSS
XBSJ231601.7–424038	AGN1	0.383	$3.05^{+0.78}_{-0.25}$	<0.072	2(< 5.5)	0.73	43.59	56/52	33.9	BSS

Columns: (1) Source name; (2) Class; (3) Redshift; (4) Photon index; (5) Intrinsic column density; (6) Reflection scaling factor and upper limit at 90% confidence (between parenthesis); (7) Observed flux in the 2-10 keV band, de-absorbed by our Galaxy; (8) Intrinsic luminosity in the 2-10 keV band; (9) χ^2 to number of degrees of freedom; (10) Null hypothesis probability; (11) Sample the source belongs to.

^p: Null hypothesis probability < 10%. ^e: Elusive AGN. ^f: Fixed parameter. Source and parameters in bold face indicates that this model is considered as our best-fit model. Note: Errors and upper limits are at 90% confidence level. Fluxes and luminosities refer to the MOS2 calibration.

Table 9. Leaky/Leaky+line fit results.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	Ratio	E keV	σ eV	EW eV	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{s}^{-1}$	Log $L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)
XBSJ021822.2-050615^p	AGN2^e	0.044	1.97^{+0.13}_{-0.24}	22.20^{+2.63}_{-1.55}	0.004^{+0.003}_{-0.001}	6.28^{+0.15}_{-0.10}	210⁺¹²⁰₋₇₀	370⁺¹⁹⁰₋₁₉₀	2.81	42.57	112/79	0.8	HBSS
XBSJ033845.7-352253	AGN2	0.113	1.99^{+0.61}_{-0.40}	31.30^{+9.70}_{-8.20}	0.024^{+0.072}_{-0.022}	1.70	43.26	42/36	22.3	HBSS
XBSJ040758.9-712833	AGN2	0.134	1.90^f	21.90^{+14.10}_{-10.60}	0.032^{+0.057}_{-0.030}	2.23	43.43	11/11	42.7	HBSS
XBSJ091828.4+513931	AGN1	0.185	1.85^{+0.43}_{-0.41}	6.41^{+2.17}_{-2.05}	0.045^{+0.060}_{-0.029}	2.56	43.58	12/15	67.8	HBSS,BSS
XBSJ095218.9-013643	AGN1	0.020	3.38^{+0.12}_{-0.04}	49.00^{+6.60}_{-5.60}	0.007^{+0.002}_{-0.002}	0.92^{+0.02}_{-0.03}	50⁺³⁰₋₂₂₀	150⁺⁴⁰₋₄₀	9.96	43.03	60/54	27.0	HBSS,BSS
XBSJ112026.7+431520	AGN2^e	0.146	1.88^{+0.66}_{-0.90}	8.25^{+2.25}_{-3.56}	0.021^{+0.057}_{-0.029}	1.88	43.26	20/15	15.8	HBSS

Columns: (1) Source name; (2) Class; (3) Redshift; (4) Photon index; (5) Intrinsic column density; (6) Scattered component to direct component ratio; (7) Emission line central energy; (8) Emission line width; (9) Emission line equivalent width; (10) Observed flux in the 2-10 keV band, de-absorbed by our Galaxy; (11) Intrinsic luminosity in the 2-10 keV band; (12) χ^2 to number of degrees of freedom; (13) Null hypothesis probability; (14) Sample the source belongs to.

^p: Null hypothesis probability < 10%. ^e: Elusive AGN. ^f: Fixed parameter. Source and parameters in bold face indicates that this model is considered as our best-fit model. Note: Errors and upper limits are at 90% confidence level. Fluxes and luminosities refer to the MOS2 calibration.

Table 10. Power law plus thermal component fit results.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	kT keV	E keV	σ eV	EW eV	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{s}^{-1}$	Log $L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)	(13)	(14)
XBSJ012654.3+191246	AGN^e	0.043	1.90^f	0.066^{+0.968}_{-0.066}	1.04^{+0.29}_{-0.11}	0.29	41.17	7/11	83.0	BSS
XBSJ021822.2-050615^p	AGN2	0.044	1.67^{+0.13}_{-0.19}	19.6^{+1.8}_{-1.7}	1.8^{+1.6}_{-0.4}	6.28^{+0.15}_{-0.40}	210⁺¹²⁰₋₇₀	360⁺²⁰₋₂₀	2.93	42.51	109/78	1.2	HBSS
XBSJ033845.7-322253^p	AGN2	0.113	1.9^f	24.1^{+8.7}_{-8.2}	0.21^{+0.5}_{-0.5}	1.61	43.17	49/36	7.3	HBSS
XBSJ040758.9-712833	AGN2	0.134	1.9^f	11.0^{+8.2}_{-4.2}	0.24^{+0.10}_{-0.08}	1.80	43.22	14/10	16.3	HBSS
XBSJ112026.7+431520	AGN2	0.146	1.9^f	7.5^{+1.7}_{-1.3}	0.87^{+0.45}_{-0.25}	1.87	43.26	18/15	25.5	HBSS
XBSJ231546.5-590313	AGN2	0.0446	1.90^f	1.25^{+0.58}_{-0.51}	0.46^{+0.10}_{-0.18}	1.49	41.95	15/14	35.9	BSS

Columns: (1) Source name; (2) Class; (3) Redshift; (4) Photon index; (5) Intrinsic column density; (6) Plasma temperature; (7) Emission line central energy; (8) Emission line width; (9) Emission line equivalent width; (10) Observed flux in the 2-10 keV band, de-absorbed by our Galaxy; (11) Intrinsic luminosity in the 2-10 keV band; (12) χ^2 to number of degrees of freedom; (13) Null hypothesis probability; (14) Sample the source belongs to.

^p: Null hypothesis probability < 10%. ^e: Elusive AGN. ^f: Fixed parameter. Source and parameters in bold face indicates that this model is considered as our best-fit model. Note: Errors and upper limits are at 90% confidence level. Fluxes and luminosities refer to the MOS2 calibration.

Table 11. Power law and absorption edges fit results.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	E keV	τ	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{ s}^{-1}$	Log $L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)	(12)
XBSJ100926.5+533426	AGN1	1.718	2.02 ^{+0.08} _{-0.11}	<0.18	2.09 ^{+0.14} _{-0.21}	0.63 ^{+0.37} _{-0.33}	0.73	45.18	40/37	34.5	BSS
XBSJ102412.3+042023	AGN1	1.458	2.01 ^{+0.16} _{-0.10}	<0.17	3.25 ^{+0.13} _{-0.13}	0.9 ^{+0.5} _{-0.4}	0.53	44.87	37/33	28.9	BSS
XBSJ204159.2-321439	AGN1	0.738	2.08 ^{+0.18} _{-0.11}	<0.10	2.10 ^{+0.15} _{-0.11}	0.70 ^{+0.30} _{-0.30}	1.05	44.45	27/26	42.3	BSS

Columns: (1) Source name; (2) Class; (3) Redshift; (4) Photon index; (5) Intrinsic column density; (6) Threshold energy; (7) Absorption depth at threshold; (8) Observed flux in the 2-10 keV band, de-absorbed by our Galaxy; (9) Intrinsic luminosity in the 2-10 keV band; (10) χ^2 to number of degrees of freedom; (11) Null hypothesis probability; (12) Sample the source belongs to. ^p: Null hypothesis probability < 10%. ^e: Elusive AGN. ^f: Fixed parameter. Source and parameters in bold face indicates that this model is considered as our best-fit model. Note: Errors and upper limits are at 90% confidence level. Fluxes and luminosities refer to the MOS2 calibration.

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Appendix A: Notes on individual sources

During the spectral fit, those sources which are not well fitted using a simple power law model can be usually well fitted by using different additional components. Here, we describe how we decided between the different models that are an acceptable fit for each source:

- **Leaky model:** All sources for which a leaky model was selected as our best-fit model share a common spectral shape: a power law shape at high energies that drops around ~ 2 –3 keV and an additional soft component. These sources are four type 2 AGN (XBSJ021822.2–050615, XBSJ033845.7–322253, XBSJ040758.9–712833 and XBSJ112026.7+431520) and two type 1 AGN (XBSJ091828.4+513931 and XBSJ095218.9–013643). XBSJ095218.9–013643 is a NLSy1 (narrow line Seyfert 1) whose intriguing X-ray spectral shape (a very steep photon index and large amount of absorption that partially covers the central source) and variability (variability of a factor of 4 in the soft X-rays) have been already studied in detail and presented in Grupe et al. (2004). In Fig. A.1 is shown an example of a leaky model fit. We find that leaving the soft photon index free to vary for all these sources does not significantly improve the fit. However, in all cases, this soft photon index steepens if it is left free to vary, which suggests the contribution of an additional soft component, most likely a thermal component given the low luminosity observed for this sources. In no case adding a thermal component to the leaky model significantly improves the fit, and by fitting a simple absorbed power law plus a thermal component always gives worse residuals at low energies than the leaky model.
- **Warm absorption:** We find that an additional ionized absorber gives a best fit in five cases (XBSJ030641.0–283559, XBSJ052543.6–334856, XBSJ140100.0–110942, XBSJ140127.7+025605, XBSJ223547.9–255836, all type 1 AGN), although the ionized absorber parameters, mainly the ionization state of the absorber, are not well constrained in all cases. We selected this model as our best-fit model when the power-law residuals at low energies showed some evidence of an structured shape resembling absorption lines or edges. In two cases, the source also displays a soft-excess (XBSJ030641.0–283559 and XBSJ223547.9–255836). One example of this model is again shown in Fig. A.1. The ionized absorber was added to the neutral one because of the way the spectral fit is carried out, i.e., our base-line model is a simple power law including neutral intrinsic absorption. It is worth noting, however, that none of the sources for which the best-fit model includes warm absorption need significant additional cold absorption, as can be seen in Table 7.
- **Absorption edges:** In three cases (XBSJ100926.5+533426, XBSJ102412.3+042023 and XBSJ204159.2–321439), an

absorption edge has to be added to the simple power law model to obtain an acceptable fit. It is not clear whether these edges are real or an instrumental effect given the energies at which they are found, but they could be caused by a warm absorber that our simple *Xspec* *absori* model is not able to fit properly. See again Fig. A.1.

- **Reflection component:** We find that a simple power law plus a neutral reflection component is a good fit in four cases (XBSJ031311.7–765428, XBSJ043448.3–775329, XBSJ052108.5–251913 and XBSJ101922.6+412049; all type 1 AGN). We use neutral reflection in all cases (*pexrav* model in *Xspec*) since our intention is not to determine where this reflection component originates in. Given the data quality, we can only estimate the amount of reflection by the reflection fraction *R* in the *pexrav* model. Nevertheless, we find that in all cases but one (XBSJ031311.7–765428) the reflection component is most likely coming from Compton-thick material far away from the central source, the putative torus in unified models, given the spectral shape, a rather flat continuum at high energies, and the characteristics of a possible Fe *K α* line. Note also that all but XBSJ043448.3–775329 show a soft-excess.
- **Black body:** A phenomenological black body component is needed to obtain an acceptable fit in 13 cases (XBSJ000532.7+200716, XBSJ005031.1–520012, XBSJ015957.5+003309, XBSJ021808.3–045845, XBSJ023530.2–523045, XBSJ023713.5–522734, XBSJ031851.9–441815, XBSJ065839.5–560813, XBSJ074312.1+742937, XBSJ141736.3+523028, XBSJ153456.1+013033, XBSJ225118.0–175951 and XBSJ231601.7–424038; all type 1 AGN), all showing soft-excess. The physical origin for this soft component is not clear although a host galaxy thermal contribution is ruled out given its high luminosity in all cases but XBSJ000532.7+200716. In that case the low luminosity found for the black body component ($\sim 4 \times 10^{42}$ erg s^{−1} in the 0.5–2.0 keV energy range) could be caused by thermal emission, but adding an *Xspec* *mekal* component does not improve the simple power law fit. More complex models recently proposed in the literature (Crummy et al. 2006; Middleton et al. 2007) cannot be used in our case given the data quality, and in any case, they are indistinguishable in the EPIC-covered energy range. In some cases, the need for a black body component instead of a more physically motivated model, could be just due to the data quality. For example, in seven cases, XBSJ005031.1–520012, XBSJ015957.5+003309, XBSJ021808.3–045845, XBSJ065839.5–560813, XBSJ074312.1+742937, XBSJ153456.1+013033 and XBSJ225118.0–175951 (see Table 7), an ionized absorber is also a good fit, but gives worse residuals than the black body model. This could be because of both the data quality and the need of a better representation of the ionized absorber. And for XBSJ021808.3–045845, XBSJ023713.5–522734, XBSJ074312.1+742937, XBSJ141736.3+523028 and XBSJ231601.7–424038 (see Table 8, note that XBSJ021828.3–045845 and XBSJ074312.1+742937 can be also fitted by using ionized absorption), the addition of a reflection component instead of a black body also significantly improves the fit. For the first two cases, this reflection component could derive from ionized material.

- **Sources for which no best fit is found:** We are unable to find an acceptable fit in 11 cases. We do not find that these sources share any common characteristic, and that an acceptable fit is not found could be simply due to our selection criteria based on the resulting null hypothesis probability. The simple power law fit for these sources is shown in Fig. A.2. In the case of XBSJ021822.2–050615 and XBSJ153456.1+013033, the fit corresponds to a leaky model and a power law plus a black body, which significantly improve the fit, but not enough to obtain a probability > 10%.

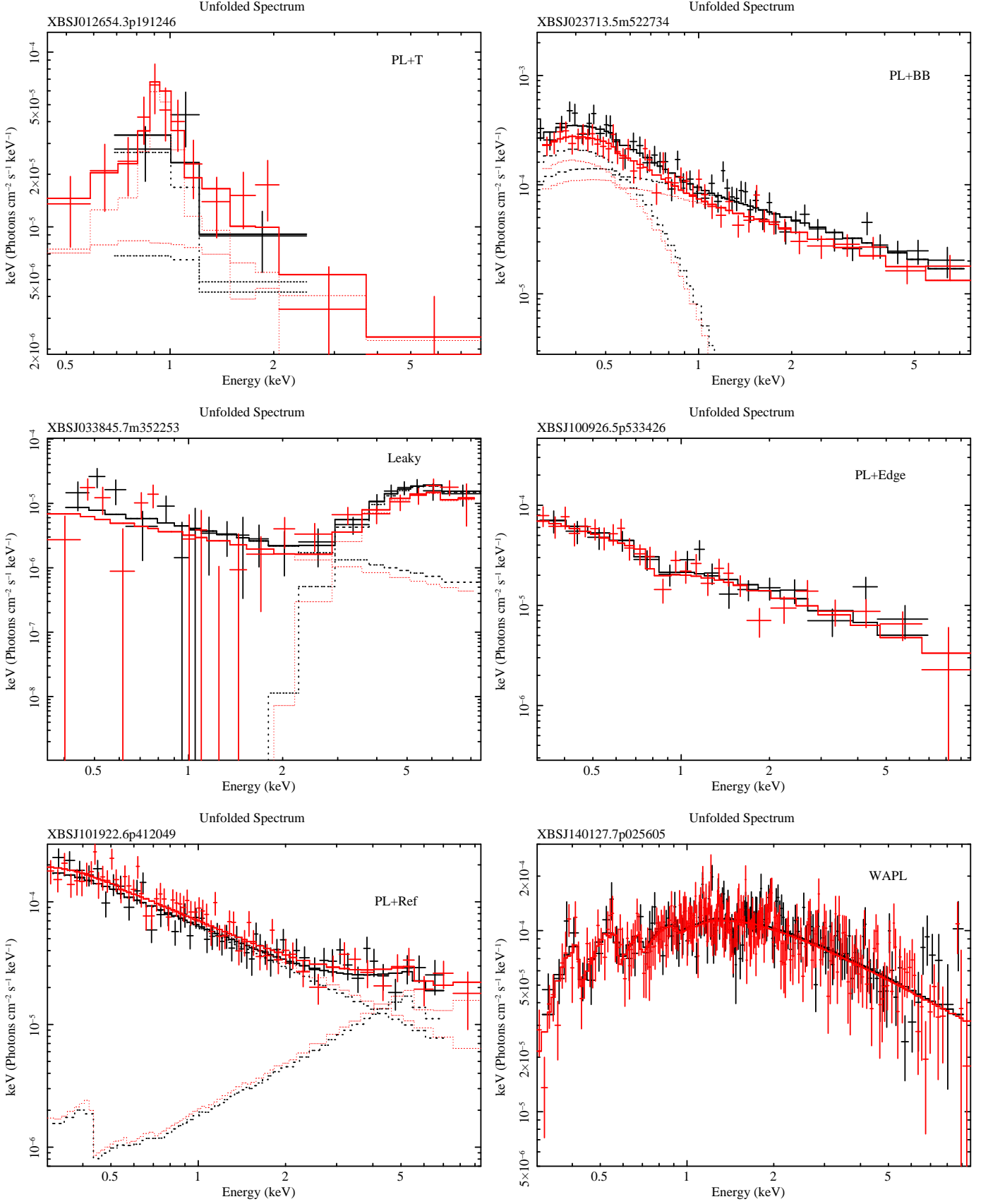


Fig. A.1. Unfolded spectra corresponding to the different additional components used during the spectral fit. From top to bottom and left to right: Power law plus thermal component, power law plus black body, leaky model, power law and absorption edge, power law plus reflection component and power law and ionized absorber.

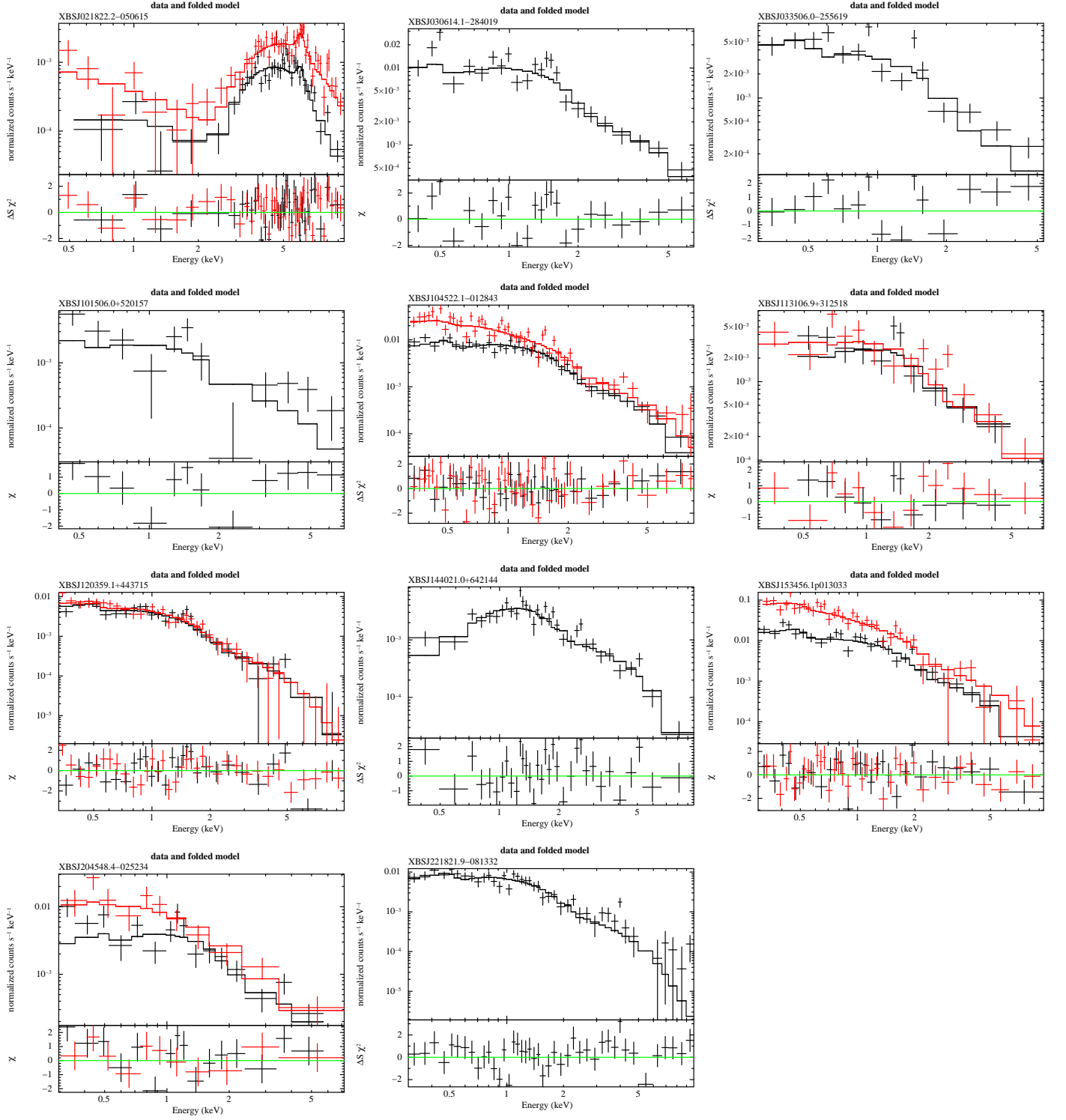


Fig. A.2. Data and residuals corresponding to the eleven sources for which no acceptable fit is found.

Table 2. Summary of *XMM-Newton* observations used.

Source name	OBSID	pn	FILTER MOS1	MOS2	N_H Gal (10^{20}cm^{-2})	pn (ks)	EXPOSURE TIME MOS1 (ks)	MOS2 (ks)	TOTAL COUNTS (0.3-10 keV)	SAMPLE
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ000027.7–250442	0125310101	Med	Med	Med	1.88	21.28	18.62	17.98	1270	BSS
XBSJ000031.7–245502	0125310101	Med	Med	Med	1.88	21.28	18.65	17.98	705	BSS
XBSJ000100.2–250501	0125310101	Med	Med	Med	1.88	21.27	18.64	17.98	830	BSS
XBSJ000102.4–245850	0125310101	Med	Med	Med	1.88	21.28	18.65	17.98	810	BSS
XBSJ000532.7+200716	0101040101	...	Med	Med	3.99	...	31.38	30.71	1000	BSS
XBSJ001831.6+162925	0111000101	Med	Med	Med	4.07	22.75	29.62	29.33	2950	BSS
XBSJ002618.5+105019	0001930101	Thin	Thin	Thin	5.07	14.18	17.91	17.51	2348	HBSS,BSS
XBSJ002637.4+165953	0050140201	Thin	Thin	Thin	4.23	37.41	46.02	45.91	3200	BSS
XBSJ002707.5+170748	0050140201	Thin	Thin	Thin	4.23	37.43	45.97	45.90	1740	BSS
XBSJ003255.9+394619	0065770101	Med	Med	Med	6.57	4.07	6.91	6.91	50	BSS
XBSJ003315.5–120700	0125920201	...	Thin	Thin	2.49	...	21.71	12.90	320	BSS
XBSJ003316.0–120456	0125920201	...	Thin	Thin	2.49	...	21.71	13.27	170	BSS
XBSJ003418.9–115940	0125920201	...	Thin	Thin	2.49	...	21.70	13.27	230	BSS
XBSJ005009.9–515934	0125320701	Med	Thin	Thin	3.21	30.23	30.22	29.53	1310	BSS
XBSJ005031.1–520012	0125320701	Med	Thin	Thin	3.21	30.23	30.22	29.53	2080	BSS
XBSJ005032.3–521543	0125320701	Med	Thin	Thin	3.21	28.30	29.64	30.08	920	BSS
XBSJ010421.4–061418	0112650401	Thin	Thin	Med	5.17	14.57	21.46	21.95	420	BSS
XBSJ010432.8–583712	0301890101	...	Med	Med	2.77	...	96.67	97.31	2270	BSS
XBSJ010701.5–172748	0025540101	Med	Med	Med	1.52	2.58	8.79	8.81	470	BSS
XBSJ010747.2–172044	0025540101	Med	Med	Med	1.52	2.58	8.82	8.80	490	BSS
XBSJ012000.0–110429	0113040801	Thin	Thin	Thin	3.26	5.51	7.81	7.92	204	BSS
XBSJ012025.2–105441	0113040801	Thin	Thin	Thin	3.26	5.51	7.82	7.92	540	BSS
XBSJ012057.4–110444	0113040801	...	Thin	Thin	3.26	...	7.82	7.92	220	BSS
XBSJ012119.9–110418	0113040801	Thin	Thin	Thin	3.26	5.51	7.82	7.92	640	BSS
XBSJ012505.4+014624	0109860101	Thin	Thin	Thin	3.10	27.81	36.92	36.03	2380	BSS
XBSJ012540.2+015752	0109860101	Thin	Thin	Thin	3.10	27.37	36.95	36.05	800	BSS
XBSJ012654.3+191246	0112600601	Med	Med	Med	4.80	5.98	3.93	3.75	163	BSS
XBSJ013204.9–400050	0112630201	Thin	Thin	Thin	1.89	24.10	28.32	28.22	920	BSS
XBSJ013240.1–133307	0084230301	Med	Med	Med	1.64	12.67	18.40	18.10	323	HBSS,BSS
XBSJ013811.7–175416	0111430101	...	Med	Med	1.44	...	32.61	32.43	700	BSS
XBSJ013944.0–674909	0032140401	Med	Med	Med	2.49	6.68	8.81	8.89	814	HBSS,BSS
XBSJ014109.9–675639	0032140401	Med	Med	Med	2.49	6.65	8.81	8.90	410	BSS
XBSJ014227.0+133453	0093641001	Med	Med	Med	4.86	6.11	10.03	10.13	230	BSS
XBSJ014251.5+133352	0093641001	Med	Med	Med	4.86	6.11	10.03	10.13	246	BSS
XBSJ015957.5+003309	0101640201	Med	Med	Med	2.59	8.00	8.15	7.95	1898	HBSS,BSS
XBSJ020029.0+002846	0101640201	Med	Med	Med	2.59	8.00	8.15	7.95	800	BSS
XBSJ020757.3+351828	0084140101	Thin	Med	Med	6.28	29.62	36.44	35.88	1040	BSS
XBSJ020845.1+351438	0084140101	Thin	Med	Med	6.28	29.63	36.47	35.88	2140	BSS
XBSJ021640.7–044404	0112371701	Thin	Thin	Thin	2.42	19.72	20.92	17.93	1675	HBSS,BSS
XBSJ021642.3–043553	0112372001	Thin	Thin	Thin	2.42	23.45	27.32	27.31	880	BSS
XBSJ021808.3–045845	0112371001	Thin	Thin	Thin	2.52	42.21	48.39	48.50	13833	HBSS,BSS

Table 2. continued.

Source name	OBSID	pn	FILTER MOS1	MOS2	N_H Gal (10^{20}cm^{-2})	pn (ks)	EXPOSURE TIME MOS1 (ks)	MOS2 (ks)	TOTAL COUNTS (0.3-10 keV)	SAMPLE
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ021817.4-045113	0112371001	Thin	Thin	Thin	2.52	42.21	48.45	48.54	4788	HBSS,BSS
XBSJ021820.6-050427	0112371001	Thin	Thin	Thin	2.52	42.21	48.57	48.59	2540	BSS
XBSJ021822.2-050615	0112371001	Thin	...	Thin	2.52	87.30	...	48.39	1800	HBSS
XBSJ021923.2-045148	0112370301	Thin	Thin	Thin	2.57	43.09	50.93	51.48	2570	BSS
XBSJ022253.0-044515	0109520501	Thin	Thin	Thin	2.64	16.47	22.57	22.77	560	BSS
XBSJ022707.7-050819	0111110201	Thin	Thin	Thin	2.63	6.58	10.02	10.00	505	BSS
XBSJ023459.7-294436	0103861101	Med	Med	Med	1.77	3.73	6.01	6.82	281	BSS
XBSJ023530.2-523045	0098810101	...	Thin	Thin	3.28	...	23.34	23.35	470	BSS
XBSJ023713.5-522734	0098810101	Thin	Thin	Thin	3.28	17.76	21.69	22.13	2136	HBSS,BSS
XBSJ023853.2-521911	0305370101	Med	3.13	111.96	1130	BSS
XBSJ024200.9+000020	0111200201	Med	Med	Med	3.53	29.90	36.36	33.62	1830	BSS
XBSJ024204.7+000814	0111200201	Med	3.53	33.60	90	BSS
XBSJ024207.3+000037	0111200201	...	Med	Med	3.53	...	36.36	33.62	600	BSS
XBSJ024325.6-000413	0111200201	Med	Med	Med	3.53	29.90	36.37	33.63	490	BSS
XBSJ025606.1+001635	0056020301	Thin	Thin	Thin	6.50	14.35	19.46	19.37	590	BSS
XBSJ025645.4+000031	0056020301	Thin	Thin	Thin	6.50	14.34	19.46	19.38	460	BSS
XBSJ030206.8-000121	0041170101	Thin	Thin	Thin	7.16	37.42	47.01	47.01	3044	HBSS,BSS
XBSJ030614.1-284019	0042340501	Thin	Thin	Thin	1.36	7.94	12.60	12.41	508	HBSS,BSS
XBSJ030641.0-283559	0042340501	Thin	Thin	Thin	1.36	7.94	12.60	12.41	620	BSS
XBSJ031015.5-765131	0122520201	Thick	Thick	Thick	8.21	23.76	28.98	28.33	4569	HBSS,BSS
XBSJ031146.1-550702	0110970101	Thin	Thin	Thin	2.55	5.48	9.79	9.38	803	HBSS,BSS
XBSJ031311.7-765428	0122520201	Thick	Thick	Thick	8.21	23.80	28.98	28.33	1170	BSS
XBSJ031401.3-545959	0110970101	Thin	Thin	Thin	2.55	5.48	9.79	9.36	180	BSS
XBSJ031549.4-551811	0129320801	Med	Med	Med	2.68	5.40	9.90	9.86	285	BSS
XBSJ031851.9-441815	0105660101	Med	Med	Med	2.60	15.60	22.12	21.91	430	BSS
XBSJ031859.2-441627	0105660101	Med	Med	Med	2.60	15.60	22.09	21.91	369	HBSS,BSS
XBSJ033208.7-274735	0108060501	Thin	Thin	Thin	0.90	38.31	48.16	48.17	1600	BSS
XBSJ033226.9-274107	0108060501	Thin	Thin	Thin	0.90	38.27	48.20	48.15	2420	BSS
XBSJ033435.5-254259	0111320101	...	Thin	Thin	1.03	...	21.23	21.37	165	BSS
XBSJ033453.9-254154	0111320101	...	Thin	Thin	1.03	...	21.23	21.37	380	BSS
XBSJ033506.0-255619	0111320101	...	Thin	Thin	1.03	...	21.28	21.38	310	BSS
XBSJ033845.7-352253	0055140101	Thin	Thin	Thin	1.31	37.28	42.23	42.78	843	HBSS
XBSJ033851.4-352646	0400620101	Med	Med	Med	1.31	69.24	121.19	120.55	3700	BSS
XBSJ033912.1-352813	0400620101	Med	Med	Med	1.31	69.23	121.18	120.59	1490	BSS
XBSJ033942.8-352411	0055140101	Thin	Thin	Thin	1.31	37.28	42.23	42.75	3250	BSS
XBSJ040658.8-712457	0111970301	Med	Med	Thin	7.57	13.36	18.10	17.08	131	HBSS
XBSJ040758.9-712833	0111970301	Med	...	Thin	7.57	13.36	...	17.08	151	HBSS
XBSJ041108.1-711341	0111970301	...	Med	Thin	7.57	...	18.06	17.09	152	HBSS,BSS
XBSJ043448.3-775329	0103861701	Med	Med	Med	7.57	4.75	7.31	7.49	240	BSS
XBSJ045942.4+015843	0112880401	Thick	Thick	Thick	7.80	13.65	18.02	17.30	225	BSS
XBSJ050011.7+013948	0112880401	Thick	Thick	Thick	7.80	13.64	18.00	17.31	340	BSS

Table 2. continued.

Source name	OBSID	pn	FILTER MOS1	MOS2	N_H Gal (10^{20}cm^{-2})	pn (ks)	EXPOSURE TIME MOS1 (ks)	MOS2 (ks)	TOTAL COUNTS (0.3-10 keV)	SAMPLE
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ050446.3–283821	0111160201	Thin	...	Thin	1.49	30.95	...	38.55	660	BSS
XBSJ050453.4–284532	0111160201	Thin	Thin	Thin	1.49	30.95	37.85	38.56	1250	BSS
XBSJ050501.8–284149	0111160201	Thin	Thin	Thin	1.49	30.95	37.96	38.55	1730	BSS
XBSJ050536.6–290050	0111160201	Thin	Thin	Thin	1.49	30.97	37.98	38.58	1000	HBSS,BSS
XBSJ051651.9+794314	0094400101	Med	Med	Med	8.00	24.51	30.72	30.05	2025	BSS
XBSJ051655.3–104104	0103860701	Med	Med	Med	9.29	4.31	3.99	4.10	200	BSS
XBSJ051822.6+793208	0094400101	Med	Med	Med	8.00	24.49	30.78	30.07	1000	BSS
XBSJ051955.5–455727	0206390101	Thin	Thin	Thin	4.16	39.15	51.26	52.01	2770	BSS
XBSJ052022.0–252309	0085640101	Thin	Thin	Thin	1.92	6.91	9.89	9.97	320	BSS
XBSJ052108.5–251913	0085640101	Thin	Thin	Thin	1.92	6.91	9.89	9.96	1512	HBSS,BSS
XBSJ052116.2–252957	0085640101	Thin	Thin	Thin	1.92	6.92	9.89	9.96	165	BSS
XBSJ052128.9–253032	0085640101	Thin	Thin	Thin	1.92	6.92	9.88	9.96	70	HBSS
XBSJ052144.1–251518	0085640101	Thin	Thin	Thin	1.92	6.91	9.89	9.97	180	BSS
XBSJ052543.6–334856	0050150301	Thin	Thin	Thin	2.22	14.41	23.77	23.79	480	BSS
XBSJ061342.7+710725	0009220601	Thin	Med	Med	8.38	7.64	10.76	10.97	1630	BSS
XBSJ062134.8–643150	0103860101	...	Med	Med	5.41	...	7.76	7.86	60	BSS
XBSJ065214.1+743230	0061540101	Med	Med	Med	6.39	21.08	17.09	17.30	2180	BSS
XBSJ065400.0+742045	0061540101	Med	Med	Med	6.39	23.02	17.08	17.27	1430	BSS
XBSJ065744.3–560817	0112980201	Thin	Thin	Thin	6.53	27.34	32.61	33.32	1050	BSS
XBSJ065839.5–560813	0112980201	Thin	Thin	Thin	6.53	27.32	32.61	33.33	1640	BSS
XBSJ074202.7+742625	0123100101	...	Thin	Thin	3.54	...	46.90	46.84	1963	HBSS,BSS
XBSJ074312.1+742937	0123100101	Thin	3.54	46.84	9740	HBSS,BSS
XBSJ074338.7+495431	0400070301	...	Med	Med	5.68	...	19.83	19.84	510	BSS
XBSJ074352.0+744258	0123100101	...	Thin	Thin	3.54	...	46.96	46.87	1430	BSS
XBSJ075117.9+180856	0111100301	...	Thin	Thin	4.12	...	28.53	28.81	860	BSS
XBSJ080504.6+245156	0094530401	Thin	Thin	Thin	3.85	14.26	20.13	20.40	640	BSS
XBSJ080608.1+244420	0094530401	Thin	Thin	Thin	3.85	14.26	20.17	20.37	3770	BSS
XBSJ083049.8+524908	0092800201	Med	Med	Med	3.85	63.62	76.57	77.11	3640	BSS
XBSJ083737.0+255151	0025540301	...	Med	Med	3.61	...	10.27	10.08	248	HBSS,BSS
XBSJ083737.1+254751	0025540301	Med	Med	Med	3.61	4.61	10.27	10.08	1593	HBSS,BSS
XBSJ083838.6+253616	0025540301	Med	Med	Med	3.61	4.61	10.27	10.07	180	BSS
XBSJ083905.9+255010	0025540301	Med	Med	Med	3.61	4.60	10.28	10.08	135	BSS
XBSJ084026.2+650638	0111400101	...	Thick	Thick	4.29	...	42.25	39.53	600	BSS
XBSJ084651.7+344634	0107860501	...	Thin	Thin	3.29	...	58.59	58.33	1010	BSS
XBSJ085530.7+585129	0085030101	Thin	Thin	Thin	4.14	7.02	11.30	11.30	590	BSS
XBSJ090729.1+620824	0110660201	Med	...	Med	4.43	5.44	...	9.13	90	BSS
XBSJ091828.4+513931	0084230601	Med	Med	Med	1.44	15.37	13.53	14.80	401	HBSS,BSS
XBSJ094526.2–085006	0017540101	Med	Med	Med	3.59	4.25	7.33	7.43	135	BSS
XBSJ094548.3–084824	0017540101	Med	Med	Med	3.59	4.25	7.32	7.43	125	BSS
XBSJ095054.5+393924	0111290101	...	Thin	Thin	1.57	...	20.63	18.49	190	BSS
XBSJ095218.9–013643	0065790101	Thin	Thin	Thin	4.00	3.57	8.37	8.27	1413	HBSS,BSS

Table 2. continued.

Source name	OBSID	pn	FILTER MOS1	MOS2	N_H Gal (10^{20}cm^{-2})	pn (ks)	EXPOSURE TIME MOS1 (ks)	MOS2 (ks)	TOTAL COUNTS (0.3-10 keV)	SAMPLE
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ095309.7+013558	0070940101	Thin	Thin	Thin	3.53	2.79	7.73	7.72	230	BSS
XBSJ095509.6+174124	0112850101	Thin	Thin	Thin	3.18	22.03	24.95	24.81	1890	BSS
XBSJ095606.4+411814	0111290201	Thin	Thin	Thin	0.79	9.10	12.30	12.37	360	BSS
XBSJ100032.5+553626	0147760101	Thin	...	Med	0.79	12.37	...	33.17	440	BSS
XBSJ100100.0+252103	0041170201	Thin	Med	Med	2.76	35.23	44.09	44.42	1930	BSS
XBSJ100309.4+554135	0147760101	Thin	Med	Med	0.79	12.38	31.94	33.15	1450	BSS
XBSJ100828.8+535408	0070340201	Thin	Thin	Thin	0.74	20.93	24.47	24.44	560	BSS
XBSJ100921.7+534926	0070340201	Thin	Thin	Thin	0.74	20.96	24.46	24.40	1390	BSS
XBSJ100926.5+533426	0070340201	Thin	Thin	Thin	0.74	20.96	24.44	24.37	900	BSS
XBSJ101506.0+520157	0086750101	...	Thin	Thin	0.76	...	8.32	8.24	126	BSS
XBSJ101511.8+520708	0086750101	Thin	Thin	Thin	0.76	6.34	8.32	8.24	480	BSS
XBSJ101706.5+520245	0086750101	Thin	Thin	Thin	0.76	6.43	8.31	8.23	700	BSS
XBSJ101838.0+411635	0028740301	Thin	Thin	Thin	1.12	19.02	23.60	23.50	1290	BSS
XBSJ101843.0+413515	0028740301	Thin	Thin	Thin	1.12	19.02	23.57	23.46	850	BSS
XBSJ101850.5+411506	0028740301	Thin	Thin	Thin	1.12	19.02	23.57	23.50	3001	HBSS,BSS
XBSJ101922.6+412049	0028740301	Thin	Thin	Thin	1.12	19.02	23.58	23.45	2642	HBSS,BSS
XBSJ102252.0+194837	0101040301	Med	Med	Med	2.15	29.33	34.37	34.78	1320	BSS
XBSJ102412.3+042023	0108670101	Thin	Thin	Thin	3.00	43.19	49.49	49.69	800	BSS
XBSJ102417.5+041656	0108670101	Thin	Thin	Thin	3.00	43.19	49.48	49.69	1380	BSS
XBSJ103120.0+311404	0102040301	Thin	Thin	Thick	1.96	20.33	25.29	24.75	755	BSS
XBSJ103154.1+310732	0102040301	Thin	Thin	Thick	1.96	20.32	25.26	24.75	370	BSS
XBSJ103745.7+532353	0112810301	Thin	Thin	Thin	1.13	13.25	17.36	17.14	890	BSS
XBSJ103909.4+205222	0059800101	Thin	Thin	Thin	2.02	4.77	11.82	12.63	425	BSS
XBSJ103932.7+205426	0059800101	Thin	Thin	Thin	2.02	4.78	11.79	12.63	490	BSS
XBSJ103935.8+533036	0112810301	Thin	Thin	Thin	1.13	11.83	17.36	17.15	1630	BSS
XBSJ104026.9+204542	0059800101	Thin	Thin	Thin	2.02	4.77	11.79	12.63	3428	HBSS,BSS
XBSJ104034.3+205110	0059800101	Thin	Thin	Thin	2.02	4.76	11.82	12.63	300	BSS
XBSJ104425.0-013521	0125300101	Thin	Thin	Thin	4.19	31.52	29.93	29.89	950	BSS
XBSJ104509.3-012442	0125300101	Thin	Thin	Thin	4.19	31.51	29.92	29.90	1040	BSS
XBSJ104522.1-012843	0125300101	Thin	Thin	Thin	4.19	31.51	29.92	29.90	1854	HBSS,BSS
XBSJ104912.8+330459	0055990201	Thin	Thin	Thin	1.98	26.13	31.18	31.55	2436	HBSS,BSS
XBSJ105014.9+331013	0055990201	...	Thin	Thin	1.98	...	31.23	31.53	310	BSS
XBSJ105239.7+572431	0147511001	Med	Med	Med	0.56	60.30	82.23	82.56	6860	BSS
XBSJ105316.9+573551	0147511801	Med	Med	Med	0.56	68.81	88.82	88.89	6380	BSS
XBSJ105335.0+572540	0147511801	Med	Med	Med	0.56	68.78	88.76	88.94	1960	BSS
XBSJ105339.7+573104	0147511801	Med	Med	Med	0.56	68.81	88.66	88.94	2824	BSS
XBSJ105624.2-033522	0094800101	Thin	Thin	Thin	3.58	26.11	34.14	34.33	1730	BSS
XBSJ110652.0-182738	0112630101	...	Thin	Thin	4.62	...	22.21	21.97	740	BSS
XBSJ111928.5+130250	0082140301	Med	Med	Med	2.43	25.74	32.62	32.63	1960	BSS
XBSJ111933.0+212756	0111290401	...	Thin	Thin	1.28	...	7.39	7.99	105	BSS
XBSJ111942.1+211516	0111290401	Thin	Thin	Thin	1.28	5.17	7.42	8.01	165	BSS

Table 2. continued.

Source name	OBSID	pn	FILTER MOS1	MOS2	N_H Gal (10^{20}cm^{-2})	pn (ks)	EXPOSURE TIME MOS1 (ks)	MOS2 (ks)	TOTAL COUNTS (0.3-10 keV)	SAMPLE
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ112022.3+125252	0093641101	Med	Med	Med	2.43	4.84	9.81	9.07	620	BSS
XBSJ112026.7+431520	0107860201	Thin	Thin	Thin	2.08	15.37	19.44	19.97	402	HBSS
XBSJ112046.7+125429	0093641101	Med	Med	Med	2.43	4.84	9.80	9.08	495	BSS
XBSJ113106.9+312518	0102040201	...	Thin	Thick	2.02	...	17.85	20.75	258	HBSS,BSS
XBSJ113121.8+310252	0102040201	Thin	Thin	Thick	2.02	14.30	17.86	20.76	706	HBSS,BSS
XBSJ113128.6-195903	0042341001	...	Thin	Thin	4.50	...	12.26	12.35	270	BSS
XBSJ113148.7+311358	0102040201	Thin	Thin	Thick	2.02	14.31	17.86	20.76	526	HBSS,BSS
XBSJ113837.9-373402	0112210501	...	Med	Med	8.50	...	123.64	123.75	4816	BSS
XBSJ115317.9+364712	0112551401	Med	Thin	Thin	1.90	6.97	12.19	12.10	180	BSS
XBSJ115846.9+551625	0090020101	Thin	Thin	Thin	1.22	7.85	11.04	10.45	290	BSS
XBSJ120359.1+443715	0109141401	...	Thin	Med	1.32	...	89.46	89.37	2030	BSS
XBSJ120413.7+443149	0109141401	...	Thin	Med	1.32	...	89.39	89.37	1000	BSS
XBSJ121501.7+140113	0112610101	Med	Med	Med	2.74	45.24	50.57	50.24	1320	BSS
XBSJ122350.4+752231	0124110101	Med	Med	Med	2.90	14.72	16.32	16.23	825	BSS
XBSJ122656.5+013126	0110990201	Thin	Thin	Thin	1.84	17.58	23.19	23.10	1454	HBSS,BSS
XBSJ123036.2+642531	0124900101	...	Thin	Thin	1.98	...	41.57	42.08	500	BSS
XBSJ123116.5+641115	0124900101	Thin	Thin	Thin	1.98	38.37	41.54	42.07	2840	BSS
XBSJ123218.5+640311	0124900101	Thin	Thin	Thin	1.98	38.42	41.55	42.08	1490	BSS
XBSJ123538.6+621644	0111550101	Thin	Thin	Thin	1.49	36.34	41.54	41.65	2690	BSS
XBSJ123759.6+621102	0111550101	Thin	Thin	Thin	1.49	36.31	41.55	41.64	3065	BSS
XBSJ123800.9+621338	0111550401	Thin	...	Thin	1.49	75.62	...	89.85	2520	BSS
XBSJ124214.1-112512	0136950201	Thin	Thin	Thin	3.64	23.49	28.09	28.25	1280	BSS
XBSJ124557.6+022659	0051760101	Thin	Thin	Thin	1.75	5.68	9.49	9.28	340	BSS
XBSJ124607.6+022153	0051760101	Thin	Thin	Thin	1.75	5.68	9.47	9.28	765	BSS
XBSJ124641.8+022412	0051760101	Thin	Thin	Thin	1.75	5.69	9.45	9.27	1582	HBSS,BSS
XBSJ124647.9+020955	0051760101	Thin	Thin	Thin	1.75	5.69	9.48	9.29	150	BSS
XBSJ124903.6-061049	0060370201	Thin	Thin	Thin	2.13	28.12	37.35	37.95	1760	BSS
XBSJ124914.6-060910	0060370201	Thin	Thin	Thin	2.13	28.12	37.35	37.95	1025	BSS
XBSJ124949.4-060722	0203270301	Med	Med	Med	2.13	34.87	44.48	46.03	1210	BSS
XBSJ125457.2+564940	0081340201	Med	Med	Med	1.27	16.57	20.19	19.98	820	BSS
XBSJ125648.4+570349	0081340201	Med	Med	Med	1.27	16.57	20.20	19.99	550	BSS
XBSJ130619.7-233857	0002940101	Med	Med	Med	9.16	4.23	6.82	7.10	320	BSS
XBSJ130658.1-234849	0002940101	Med	Med	Med	9.16	4.30	6.82	7.10	165	BSS
XBSJ132038.0+341124	0093640401	Med	Med	Med	1.00	13.39	18.31	18.65	2078	HBSS,BSS
XBSJ132052.5+341742	0093640401	Med	Med	Med	1.00	13.39	18.36	18.65	545	BSS
XBSJ132101.6+340656	0093640401	Med	Med	Med	1.00	13.38	18.35	18.65	3295	BSS
XBSJ132105.5+341459	0093640401	Med	Med	Med	1.00	13.38	18.36	18.65	1250	BSS
XBSJ133023.8+241707	0100240201	Med	Med	Med	1.16	29.51	34.84	34.80	1510	BSS
XBSJ133026.6+241520	0100240201	Med	Med	Med	1.16	29.51	34.87	34.78	3510	BSS
XBSJ133807.5+242411	0402080301	...	Thin	Thin	1.16	...	59.75	59.79	710	BSS
XBSJ133942.6-315004	0035940301	Thin	Med	Med	3.87	28.88	36.77	37.43	1081	HBSS,BSS

Table 2. continued.

Source name	OBSID	pn	FILTER MOS1	MOS2	N_H Gal (10^{20}cm^{-2})	pn (ks)	EXPOSURE TIME MOS1 (ks)	MOS2 (ks)	TOTAL COUNTS (0.3-10 keV)	SAMPLE
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ134656.7+580315	0112250201	Med	Thin	Thin	1.28	32.42	34.88	34.55	1236	HBSS
XBSJ134749.9+582111	0112250201	Med	Thin	Thin	1.28	32.07	34.82	34.57	19091	HBSS,BSS
XBSJ140100.0-110942	0109910101	Thin	Thin	Thin	4.29	36.79	44.70	45.93	1450	BSS
XBSJ140102.0-111224	0109910101	Thin	...	Thin	4.29	36.83	...	45.97	8037	HBSS,BSS
XBSJ140113.4+024016	0098010101	Thin	Thin	Thin	2.32	30.69	34.51	34.66	596	HBSS
XBSJ140127.7+025605	0098010101	Thin	Thin	Thin	2.32	30.68	34.56	34.67	8933	HBSS,BSS
XBSJ140921.1+261336	0092850101	Thin	Thin	Thin	1.49	12.14	55.80	53.18	2465	BSS
XBSJ141235.8-030909	0013140101	Med	Med	Med	3.81	11.74	15.06	15.78	490	BSS
XBSJ141531.5+113156	0112250301	Thin	Thin	Thin	1.82	18.21	24.13	22.75	3015	HBSS,BSS
XBSJ141643.8+521434	0127921001	Thin	Thin	Thin	1.30	39.90	51.17	50.40	1830	BSS
XBSJ141658.8+521202	0127921001	Thin	Thin	Thin	1.30	39.89	51.17	50.41	930	BSS
XBSJ141722.6+251335	0109960101	...	Thin	Med	1.69	...	20.39	21.63	280	BSS
XBSJ141736.3+523028	0127921001	Thin	Thin	Thin	1.30	44.27	54.47	54.51	2324	BSS
XBSJ141809.1+250040	0109960101	...	Thin	Med	1.69	...	20.39	21.63	420	BSS
XBSJ142741.8+423335	0111850201	...	Med	Med	1.38	...	48.65	48.83	487	HBSS
XBSJ143835.1+642928	0111530101	...	Thick	Thick	1.68	...	49.91	49.12	394	HBSS,BSS
XBSJ143911.2+640526	0111530101	...	Thick	Thick	1.68	...	49.93	49.16	125	HBSS
XBSJ144021.0+642144	0111530101	...	Thick	Thick	1.68	...	49.97	49.12	720	BSS
XBSJ144937.5+090826	0057560301	Thin	Thin	Thin	2.02	32.52	38.37	38.47	1840	BSS
XBSJ144945.8+085921	0057560301	Thin	Thin	Thin	2.02	32.53	38.32	38.45	1310	BSS
XBSJ145857.1-313535	0067750101	Thin	Thin	Thin	8.43	30.59	40.63	40.74	1190	BSS
XBSJ150428.3+101856	0205340401	...	Thin	Thin	2.34	...	20.62	20.71	500	BSS
XBSJ151815.0+060851	0018741001	Thin	Thin	Thin	3.20	4.25	6.96	7.07	75	BSS
XBSJ153205.7-082952	0100240801	Med	Med	Med	9.03	19.62	24.25	23.95	940	BSS
XBSJ153419.0+011808	0112190401	Med	Med	Med	4.89	8.87	13.41	13.73	480	BSS
XBSJ153452.3+013104	0112190401	Med	Med	Med	4.89	8.87	13.39	13.72	2637	HBSS,BSS
XBSJ153456.1+013033	0112190401	Med	Med	Med	4.89	8.87	13.39	13.72	1435	BSS
XBSJ160645.9+081525	0067340601	Med	Med	Med	4.01	7.93	13.68	13.73	211	HBSS,BSS
XBSJ160706.6+075709	0067340601	Med	Med	Med	4.01	8.00	13.74	13.74	685	BSS
XBSJ160731.5+081202	0067340601	Med	Med	Med	4.01	7.98	13.72	13.71	660	BSS
XBSJ161544.2+121708	0103460801	Med	Med	Thin	4.58	9.58	13.08	13.06	330	BSS
XBSJ161615.1+121353	0103460801	Med	Med	Thin	4.58	9.17	12.76	12.53	440	BSS
XBSJ161820.7+124116	0103461001	Med	Med	Thin	4.55	7.75	10.77	10.95	40	HBSS
XBSJ161825.4+124145	0103461001	Med	Med	Thin	4.55	7.75	10.77	10.95	120	BSS
XBSJ162813.9+780342	0061940301	...	Med	Med	4.04	...	4.63	4.74	170	BSS
XBSJ163141.1+781239	0400920201	Thin	Med	Med	4.04	7.44	13.59	13.50	560	BSS
XBSJ163223.6+052547	0112230301	Med	Med	Med	5.67	15.88	19.75	19.84	1070	BSS
XBSJ163309.8+571039	0049540101	Med	Med	Med	1.87	6.17	8.52	8.61	1070	BSS
XBSJ163332.3+570520	0049540101	Med	Med	Med	1.87	6.17	8.52	8.52	700	BSS
XBSJ163427.5+781002	0400920201	Thin	Med	Med	4.04	7.44	13.58	13.50	250	BSS
XBSJ164237.9+030014	0067340501	Med	Med	Med	5.47	6.99	12.92	13.14	255	BSS

Table 2. continued.

Source name	OBSID	pn	FILTER MOS1	MOS2	N_H Gal (10^{20}cm^{-2})	pn (ks)	EXPOSURE TIME MOS1 (ks)	MOS2 (ks)	TOTAL COUNTS (0.3-10 keV)	SAMPLE
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ165314.4+141943	0113070101	Thin	Thin	Med	5.54	13.91	20.09	20.70	535	BSS
XBSJ165406.6+142123	0113070101	Thin	Thin	Med	5.54	13.93	20.09	20.70	1040	BSS
XBSJ165425.3+142159	0113070101	Thin	Thin	Med	5.54	13.92	20.11	20.70	7683	HBSS,BSS
XBSJ165448.5+141311	0113070101	Thin	Thin	Med	5.54	13.90	20.08	20.70	2184	HBSS,BSS
XBSJ165800.7+352333	0111060101	Med	1.77	10.10	140	BSS
XBSJ172230.6+341344	0102040101	Thin	Thin	Thick	3.11	3.71	6.37	6.06	210	BSS
XBSJ185518.7-462504	0067340101	Med	Med	Med	5.29	4.99	7.59	9.69	840	BSS
XBSJ185613.7-462239	0067340101	Med	Med	Med	5.29	5.15	7.58	9.69	1015	BSS
XBSJ193138.9-725115	0081341001	Thin	Thin	Thin	5.95	14.74	18.40	18.60	860	BSS
XBSJ193248.8-723355	0081341001	Thin	Thin	Thin	5.95	14.75	18.38	18.59	782	HBSS,BSS
XBSJ204043.4-004548	0111180201	Med	Med	Med	6.72	9.59	13.53	13.34	302	HBSS,BSS
XBSJ204159.2-321439	0111510101	...	Med	Med	5.08	...	19.56	18.67	660	BSS
XBSJ204204.1-321601	0111510101	...	Med	Med	5.08	...	19.56	18.67	250	BSS
XBSJ204208.2-323523	0111510101	...	Med	Med	5.08	...	19.56	18.76	280	BSS
XBSJ204548.4-025234	0112600501	Med	Med	Med	4.98	6.10	9.64	10.35	260	BSS
XBSJ205411.9-160804	0083210101	Thin	Thin	Thin	4.34	5.90	8.73	9.23	215	BSS
XBSJ205429.9-154937	0083210101	Thin	Thin	Thin	4.34	5.90	8.72	9.24	200	BSS
XBSJ205635.7-044717	0112190601	Med	Med	Med	4.96	8.32	15.63	14.91	567	HBSS,BSS
XBSJ205829.9-423634	0081340401	Thin	Thin	Thin	3.89	12.76	16.24	16.05	2823	HBSS,BSS
XBSJ210325.4-112011	0041150101	Med	Med	Med	4.70	26.25	34.65	33.67	820	BSS
XBSJ210355.3-121858	0038540301	Thin	Thin	Thin	4.48	11.41	14.12	13.99	385	BSS
XBSJ212759.5-443924	0088020201	Thin	Thin	Thin	3.50	19.73	24.23	24.71	860	BSS
XBSJ213002.3-153414	0103060101	...	Med	Med	4.99	...	19.12	19.21	642	HBSS,BSS
XBSJ213719.6-433347	0109463501	Thin	Thin	Thin	2.91	4.53	7.33	7.33	370	BSS
XBSJ213729.7-423601	0061940201	...	Med	Med	2.68	...	10.40	10.41	375	BSS
XBSJ213733.2-434800	0109463501	Thin	Thin	Thin	2.91	4.53	7.32	7.32	120	BSS
XBSJ213757.6-422334	0061940201	...	Med	Med	2.68	...	10.41	10.41	220	BSS
XBSJ213820.2-142536	0092850201	Med	Med	Med	4.73	38.80	29.66	30.84	1996	HBSS,BSS
XBSJ213824.0-423019	0061940201	Thin	Med	Med	2.68	6.72	10.40	10.41	2120	BSS
XBSJ213829.8-423958	0061940201	...	Med	Med	2.68	...	10.40	10.41	300	BSS
XBSJ213852.2-434714	0109463501	Thin	Thin	Thin	2.91	4.53	7.31	7.32	175	BSS
XBSJ214041.4-234720	0008830101	Thin	Thin	Thin	3.55	14.79	16.86	17.67	2234	HBSS,BSS
XBSJ215244.2-302407	0103060401	Med	Med	Med	2.12	18.13	22.54	22.35	1640	BSS
XBSJ220320.8+184930	0130920101	Thin	Thin	Thick	6.02	12.16	13.97	15.69	850	BSS
XBSJ220446.8-014535	0012440301	Thin	Thin	Thin	6.13	23.03	29.77	29.56	575	BSS
XBSJ220601.5-015346	0012440301	...	Thin	Thin	6.13	...	29.75	29.57	532	HBSS,BSS
XBSJ221623.3-174317	0106660101	Thin	Thin	Thin	2.36	46.53	54.07	53.86	450	BSS
XBSJ221722.4-082018	0009650201	...	Med	Med	5.31	...	27.18	27.59	790	BSS
XBSJ221729.3-081154	0009650201	...	Med	Med	5.31	...	27.23	27.60	740	BSS
XBSJ221821.9-081332	0009650201	...	Med	Med	5.31	...	27.17	27.59	955	BSS
XBSJ221951.6+120123	0103861201	Med	Med	Med	5.41	7.33	10.08	10.19	270	BSS

Table 2. continued.

Source name	OBSID	pn	FILTER MOS1	MOS2	N_H Gal (10^{20}cm^{-2})	pn (ks)	EXPOSURE TIME MOS1 (ks)	MOS2 (ks)	TOTAL COUNTS (0.3-10 keV)	SAMPLE
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ223547.9–255836	0111790101	...	Med	Med	1.46	...	35.92	35.88	835	BSS
XBSJ223555.0–255833	0111790101	...	Med	Med	1.46	...	35.92	35.88	735	BSS
XBSJ223949.8+080926	0103860301	Med	Med	Med	6.68	3.62	6.23	6.03	105	BSS
XBSJ224756.6–642721	0112240101	...	Med	Med	2.95	...	27.93	28.31	320	BSS
XBSJ225025.1–643225	0112240101	Med	Med	Med	2.95	22.51	27.96	28.31	740	BSS
XBSJ225050.2–642900	0112240101	Med	Med	Med	2.95	22.51	27.96	28.30	2355	BSS
XBSJ225118.0–175951	0081340901	Med	Med	Med	2.69	16.90	21.48	21.23	2220	BSS
XBSJ230400.4–083755	0109130701	...	Thin	Thin	3.60	...	9.49	9.78	110	BSS
XBSJ230401.0+031519	0033541001	Thin	Thin	Thin	5.26	7.96	11.57	11.99	210	BSS
XBSJ230434.1+122728	0025541001	Med	Med	Med	4.96	6.87	10.19	9.98	215	BSS
XBSJ230443.8+121636	0025541001	Med	Med	Med	4.96	7.33	10.18	9.98	425	BSS
XBSJ230459.6+121205	0025541001	Med	Med	Med	4.96	7.33	10.19	9.99	220	BSS
XBSJ230522.1+122121	0025541001	Med	Med	Med	4.96	6.87	10.18	9.99	215	BSS
XBSJ231342.5–423210	0123900101	Thin	Thin	Thin	1.79	45.40	50.16	50.23	3620	BSS
XBSJ231546.5–590313	0079940101	...	Thin	Thin	2.90	...	17.14	16.63	270	BSS
XBSJ231601.7–424038	0093640701	Med	Med	Med	1.85	12.00	17.27	17.27	1170	BSS
XBSJ233421.9–151219	0093550401	Med	Med	Med	2.01	17.10	19.72	21.55	575	BSS
XBSJ235036.9+362204	0100241001	Thick	Thick	Thick	8.09	5.43	8.39	8.39	345	BSS

Columns: (1) Source name; (2) Observation identifier; (3), (4), (5) Filter in use during the observation for the pn, MOS1 and MOS2 cameras, respectively; (6) Galactic column density toward the used *XMM-Newton* pointing. (7), (8), (9) Resulting exposure time after removing high-background intervals for the pn, MOS1 and MOS2 cameras, respectively; (10) Total EPIC counts (pn+MOS1+MOS2); (11) Sample the source belongs to.

Table 3. Power law fit results.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample	Best fit
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ000027.7–250442	AGN1	0.336	$1.87^{+0.10}_{-0.08}$	<0.02	0.72	43.48	64/66	55.4	BSS	Y
XBSJ000031.7–245502	AGN1	0.284	$2.29^{+0.14}_{-0.13}$	<0.02	0.52	43.22	31/33	56.0	BSS	Y
XBSJ000100.2–250501	AGN1	0.850	$1.67^{+0.14}_{-0.18}$	$0.71^{+0.22}_{-0.24}$	1.38	44.66	48/41	22.4	BSS	Y
XBSJ000102.4–245850	AGN1	0.433	$2.12^{+0.13}_{-0.12}$	<0.02	0.67	43.74	40/40	47.7	BSS	Y
XBSJ000532.7+200716 ^p	AGN1 ^e	0.119	$3.42^{+0.15}_{-0.15}$	<0.01	0.21	42.03	64/37	0.4	BSS	N,5
XBSJ001831.6+162925	AGN1	0.553	$2.39^{+0.06}_{-0.06}$	<0.01	0.76	44.11	123/126	54.8	BSS	Y
XBSJ002618.5+105019	AGN1	0.473	$2.04^{+0.06}_{-0.06}$	<0.009	2.78	44.44	115/102	17.6	HBSS,BSS	Y
XBSJ002637.4+165953	AGN1	0.554	$2.15^{+0.07}_{-0.05}$	<0.02	1.23	44.27	123/138	81.9	BSS	Y
XBSJ002707.5+170748	AGN1	0.930	$1.86^{+0.11}_{-0.11}$	$0.08^{+0.08}_{-0.07}$	1.05	44.68	97/83	13.6	BSS	Y
XBSJ003255.9+394619	AGN1	1.139	$2.64^{+1.52}_{-0.87}$	<0.54	0.25	44.52	3/1	11.3	BSS	Y
XBSJ003315.5–120700	AGN1	1.206	$2.01^{+0.28}_{-0.16}$	<0.21	0.99	44.98	12/14	57.1	BSS	Y
XBSJ003316.0–120456	AGN1	0.660	$2.60^{+0.72}_{-0.29}$	<0.22	0.26	43.89	5/6	61.1	BSS	Y
XBSJ003418.9–115940	AGN1	0.850	$2.10^{+0.44}_{-0.26}$	$0.04^{+0.25}_{-0.04}$	0.55	44.37	6/9	71.3	BSS	Y
XBSJ005009.9–515934	AGN1	0.610	$2.28^{+0.15}_{-0.13}$	<0.03	0.53	44.04	69/65	34.4	BSS	Y
XBSJ005031.1–520012 ^p	AGN1	0.463	$2.27^{+0.10}_{-0.10}$	<0.01	0.73	43.87	97/80	9.9	BSS	N,5
XBSJ005032.3–521543	AGN1	1.216	$2.21^{+0.36}_{-0.25}$	$0.07^{+0.22}_{-0.07}$	0.52	44.78	118/118	48.7	BSS	Y
XBSJ010421.4–061418	AGN1	0.520	$1.87^{+0.25}_{-0.15}$	<0.12	0.66	43.88	11/18	91.0	BSS	Y
XBSJ010432.8–583712	AGN1	1.640	$1.95^{+0.08}_{-0.06}$	<0.03	1.24	45.38	97/86	19.6	BSS	Y
XBSJ010701.5–172748	AGN1	0.890	$2.02^{+0.35}_{-0.18}$	<0.19	1.51	44.84	22/19	28.9	BSS	Y
XBSJ010747.2–172044	AGN1	0.980	$2.46^{+0.25}_{-0.23}$	<0.08	0.69	44.74	23/20	29.7	BSS	Y
XBSJ012000.0–110429	AGN1	0.351	$3.01^{+3.62}_{-0.50}$	<0.54	0.12	42.91	9/6	17.0	BSS	Y
XBSJ012025.2–105441	AGN1	1.338	$2.40^{+0.34}_{-0.29}$	$0.03^{+0.46}_{-0.03}$	0.52	44.95	11/22	96.8	BSS	Y
XBSJ012057.4–110444	AGN2	0.072	$1.94^{+0.31}_{-0.26}$	<0.05	1.37	42.30	9/8	31.8	BSS	Y
XBSJ012119.9–110418	AGN1	0.204	$2.66^{+0.38}_{-0.23}$	$0.02^{+0.07}_{-0.02}$	1.33	43.32	28/27	41.2	BSS	Y
XBSJ012505.4+014624	AGN1	1.567	$1.65^{+0.05}_{-0.06}$	<0.06	1.62	45.33	84/104	92.1	BSS	Y
XBSJ012540.2+015752	AGN1 ^e	0.123	$1.83^{+0.13}_{-0.11}$	<0.02	0.89	42.60	54/43	12.0	BSS	Y
XBSJ012654.3+191246 ^p	AGN ^e	0.043	1.90^f	<0.06	0.65	41.51	27/13	1.2	BSS	N,9
XBSJ013204.9–400050	AGN1	0.450	$2.42^{+0.28}_{-0.23}$	$0.04^{+0.07}_{-0.04}$	0.43	43.64	38/40	57.7	BSS	Y
XBSJ013240.1–133307	AGN2	0.562	1.90^f	$2.55^{+0.70}_{-0.57}$	1.76	44.43	21/19	33.0	HBSS,BSS	Y
XBSJ013811.7–175416	BL Lac	0.530	$2.57^{+0.26}_{-0.23}$	$0.11^{+0.08}_{-0.07}$	0.33	43.74	40/30	11.4	BSS	Y
XBSJ013944.0–674909	AGN1 ^e	0.104	$1.95^{+0.13}_{-0.12}$	<0.02	1.13	42.56	32/35	64.5	HBSS,BSS	Y
XBSJ014109.9–675639	AGN1 ^e	0.226	$1.78^{+0.39}_{-0.34}$	$0.12^{+0.21}_{-0.12}$	2.02	43.53	12/16	77.9	BSS	Y
XBSJ014227.0+133453	AGN1 ^e	0.275	$1.94^{+0.27}_{-0.23}$	<0.05	0.66	43.25	13/11	28.3	BSS	Y
XBSJ014251.5+133352	AGN1	1.071	$1.86^{+0.24}_{-0.19}$	<0.26	0.81	44.72	6/8	69.3	BSS	Y
XBSJ015957.5+003309 ^p	AGN1	0.310	$2.22^{+0.09}_{-0.08}$	<0.007	2.59	44.00	115/82	1.0	HBSS,BSS	N,5
XBSJ020029.0+002846	AGN1	0.174	$2.42^{+0.17}_{-0.16}$	<0.02	0.94	42.99	45/35	12.1	BSS	Y
XBSJ020757.3+351828	AGN1	0.188	$1.99^{+0.21}_{-0.16}$	$0.02^{+0.05}_{-0.02}$	0.78	42.96	41/45	64.2	BSS	Y
XBSJ020845.1+351438	AGN1	0.415	$1.93^{+0.11}_{-0.09}$	$0.01^{+0.04}_{-0.01}$	1.22	43.93	86/94	69.6	BSS	Y

Table 3. continued.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample	Best fit
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ021640.7-044404	AGN1	0.873	$2.24^{+0.09}_{-0.07}$	<0.02	1.10	44.74	85/72	14.6	HBSS,BSS	Y
XBSJ021642.3-043553	AGN2	1.985	$1.91^{+0.18}_{-0.17}$	$4.20^{+1.20}_{-1.00}$	1.02	45.49	45/39	24.3	BSS	Y
XBSJ021808.3-045845 ^p	AGN1	0.712	$2.21^{+0.03}_{-0.02}$	<0.002	2.05	44.78	536/419	0.01	HBSS,BSS	N,5
XBSJ021817.4-045113	AGN1	1.080	$1.83^{+0.07}_{-0.05}$	$0.02^{+0.05}_{-0.02}$	2.57	45.22	207/192	20.7	HBSS,BSS	Y
XBSJ021820.6-050427	AGN1	0.646	$1.81^{+0.07}_{-0.06}$	<0.01	0.87	44.21	112/111	46.3	BSS	Y
XBSJ021822.2-050615 ^p	AGN2 ^e	0.044	$1.90^{+0.12}_{-0.07}$	$20.78^{+4.22}_{-0.01}$	2.94	42.60	151/84	0.001	HBSS	N,8
XBSJ021923.2-045148	AGN1	0.632	$2.41^{+0.12}_{-0.07}$	<0.03	0.42	44.00	121/109	20.8	BSS	Y
XBSJ022253.0-044515	AGN1	1.420	$1.84^{+0.16}_{-0.13}$	<0.15	0.59	44.86	28/23	21.7	BSS	Y
XBSJ022707.7-050819	AGN2	0.358	$1.68^{+0.48}_{-0.37}$	$1.31^{+0.80}_{-0.60}$	2.58	44.09	16/21	78.0	BSS	Y
XBSJ023459.7-294436	AGN1	0.446	$1.93^{+1.83}_{-0.57}$	<0.77	0.59	43.69	10/10	40.2	BSS	Y
XBSJ023530.2-523045 ^p	AGN1	0.428	$2.14^{+0.16}_{-0.16}$	<0.03	0.85	43.84	28/19	9.3	BSS	N,5
XBSJ023713.5-522734 ^p	AGN1	0.193	$2.43^{+0.08}_{-0.08}$	<0.003	1.70	43.35	120/89	1.6	HBSS,BSS	N,5
XBSJ023853.2-521911	AGN1	0.648	$2.12^{+0.15}_{-0.10}$	<0.06	0.71	44.20	40/44	66.6	BSS	Y
XBSJ024200.9+000020	AGN1	1.112	$2.03^{+0.08}_{-0.07}$	<0.03	1.06	44.93	78/80	52.4	BSS	Y
XBSJ024204.7+000814	AGN1	0.383	$2.26^{+0.65}_{-0.38}$	<0.16	0.24	43.20	8/6	23.0	BSS	Y
XBSJ024207.3+000037	AGN1	0.385	$2.52^{+0.20}_{-0.14}$	<0.04	0.38	43.43	25/24	41.6	BSS	Y
XBSJ024325.6-000413	AGN1	0.356	$1.74^{+0.28}_{-0.15}$	$0.09^{+0.17}_{-0.09}$	0.75	43.54	23/20	29.9	BSS	Y
XBSJ025606.1+001635	AGN1	0.629	$2.20^{+0.41}_{-0.21}$	<0.14	0.52	44.05	22/25	63.2	BSS	Y
XBSJ025645.4+000031	AGN1 ^e	0.358	$2.06^{+0.23}_{-0.20}$	<0.05	0.51	43.42	30/25	21.4	BSS	Y
XBSJ030206.8-000121	AGN1	0.641	$1.89^{+0.05}_{-0.05}$	<0.02	2.28	44.64	144/129	16.1	HBSS,BSS	Y
XBSJ030614.1-284019 ^p	AGN1	0.278	$1.66^{+0.16}_{-0.13}$	<0.04	2.85	43.87	42/21	0.4	HBSS,BSS	Y
XBSJ030641.0-283559 ^p	AGN1	0.367	$2.56^{+0.24}_{-0.12}$	<0.02	0.31	43.30	44/26	1.7	BSS	N,6
XBSJ031015.5-765131	AGN1	1.187	$1.91^{+0.04}_{-0.04}$	<0.03	3.44	45.47	180/194	73.4	HBSS,BSS	Y
XBSJ031146.1-550702	AGN2	0.162	$2.05^{+0.12}_{-0.10}$	<0.02	2.98	43.40	38/35	32.4	HBSS,BSS	Y
XBSJ031311.7-765428 ^p	AGN1	1.274	$1.89^{+0.10}_{-0.09}$	<0.05	0.89	44.95	67/52	7.6	BSS	N,7
XBSJ031401.3-545959	AGN1	0.841	$1.84^{+0.35}_{-0.36}$	<0.11	0.70	44.40	13/8	10.7	BSS	Y
XBSJ031549.4-551811	AGN1	0.808	$1.87^{+0.23}_{-0.21}$	<0.07	0.75	44.39	9/10	52.3	BSS	Y
XBSJ031851.9-441815 ^p	AGN1	1.360	$1.86^{+0.21}_{-0.20}$	<0.08	0.72	44.91	32/17	1.6	BSS	N,5
XBSJ031859.2-441627	AGN1 ^e	0.140	$1.48^{+0.31}_{-0.26}$	$0.28^{+0.27}_{-0.20}$	1.63	42.97	8/14	88.6	HBSS,BSS	Y
XBSJ033208.7-274735	AGN1	0.544	$1.99^{+0.15}_{-0.11}$	$0.02^{+0.04}_{-0.02}$	0.60	43.92	55/69	88.3	BSS	Y
XBSJ033226.9-274107	AGN1	0.736	$2.21^{+0.07}_{-0.07}$	<0.01	0.58	44.27	111/102	25.2	BSS	Y
XBSJ033435.5-254259	AGN1	1.190	$2.18^{+0.39}_{-0.29}$	<0.20	0.39	44.61	4/8	86.4	BSS	Y
XBSJ033453.9-254154	AGN1	1.160	$1.87^{+0.26}_{-0.25}$	$0.18^{+0.27}_{-0.18}$	0.93	44.87	12/15	69.2	BSS	Y
XBSJ033506.0-255619 ^p	AGN1	1.430	$2.10^{+0.30}_{-0.22}$	<0.16	0.80	45.11	37/12	0.02	BSS	Y
XBSJ033845.7-352253 ^p	AGN2	0.113	$1.90^{+0.12}_{-0.07}$	$23.5^{+7.7}_{-6.4}$	1.54	43.15	81/38	0.006	HBSS	N,8
XBSJ033851.4-352646	AGN1	1.070	$1.78^{+0.08}_{-0.08}$	$0.15^{+0.07}_{-0.07}$	0.64	44.59	158/153	38.1	BSS	Y
XBSJ033912.1-352813	AGN1	0.466	$1.46^{+0.10}_{-0.12}$	$0.07^{+0.06}_{-0.04}$	0.51	43.59	77/66	18.2	BSS	Y
XBSJ033942.8-352411	AGN1	1.043	$2.50^{+0.06}_{-0.07}$	<0.02	0.36	44.53	139/130	28.0	BSS	Y

Table 3. continued.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{ s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample	Best fit
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ040658.8–712457	AGN2	0.181	1.90^f	$20.60^{+15.10}_{-8.30}$	1.79	43.60	9/9	40.2	HBSS	Y
XBSJ040758.9–712833 ^p	AGN2	0.134	1.90^f	$10.70^{+8.90}_{-4.30}$	1.77	43.21	24/12	1.9	HBSS	N,8
XBSJ041108.1–711341	AGN1	0.923	$1.91^{+0.52}_{-0.32}$	<0.62	0.85	44.60	6/6	42.7	HBSS,BSS	Y
XBSJ043448.3–775329 ^p	AGN1 ^e	0.097	1.90^f	$0.48^{+0.27}_{-0.21}$	2.06	42.28	21/13	7.4	BSS	N,7
XBSJ045942.4+015843	AGN1	0.248	$1.55^{+0.42}_{-0.27}$	$0.04^{+0.37}_{-0.04}$	1.47	43.46	8/11	67.9	BSS	Y
XBSJ050011.7+013948	AGN1	0.360	$2.01^{+0.41}_{-0.33}$	$0.29^{+0.36}_{-0.26}$	0.99	43.71	18/18	46.4	BSS	Y
XBSJ050446.3–283821	AGN1	0.840	$1.97^{+0.18}_{-0.14}$	<0.05	0.56	44.33	30/28	38.2	BSS	Y
XBSJ050453.4–284532	AGN1 ^e	0.204	$1.46^{+0.11}_{-0.11}$	$0.04^{+0.05}_{-0.04}$	1.13	43.16	46/56	80.6	BSS	Y
XBSJ050501.8–284149	AGN1	0.257	$2.18^{+0.09}_{-0.09}$	<0.01	0.53	43.11	77/77	47.4	BSS	Y
XBSJ050536.6–290050	AGN2	0.577	$1.85^{+0.20}_{-0.18}$	$0.61^{+0.21}_{-0.17}$	1.33	44.30	39/44	70.5	HBSS,BSS	Y
XBSJ051651.9+794314	AGN1	0.557	$1.56^{+0.08}_{-0.09}$	$0.24^{+0.09}_{-0.08}$	1.95	44.37	66/92	97.9	BSS	Y
XBSJ051655.3–104104	AGN1	0.568	$1.88^{+0.62}_{-0.30}$	<0.57	1.20	44.23	10/9	38.5	BSS	Y
XBSJ051822.6+793208	AGN1 ^e	0.052	$1.83^{+0.12}_{-0.10}$	<0.02	0.94	41.85	41/44	60.4	BSS	Y
XBSJ051955.5–455727	AGN1	0.562	$2.09^{+0.06}_{-0.06}$	<0.007	1.07	44.21	124/119	36.4	BSS	Y
XBSJ052022.0–252309	AGN1	0.745	$2.05^{+0.38}_{-0.17}$	$0.22^{+0.28}_{-0.22}$	0.73	44.35	15/17	56.1	BSS	Y
XBSJ052108.5–251913 ^p	AGN1	1.196	$1.97^{+0.08}_{-0.07}$	<0.03	2.04	45.27	88/66	3.7	HBSS,BSS	N,7
XBSJ052116.2–252957	AGN1 ^e	0.332	$2.21^{+0.71}_{-0.44}$	$0.07^{+0.22}_{-0.07}$	0.36	43.21	5/7	71.2	BSS	Y
XBSJ052128.9–253032	AGN2 ^e	0.588	1.90^f	$13.80^{+7.50}_{-4.57}$	1.42	44.49	3/4	57.6	HBSS	Y
XBSJ052144.1–251518	AGN1	0.321	$2.10^{+0.41}_{-0.27}$	<0.08	0.60	43.39	8/8	43.0	BSS	Y
XBSJ052543.6–334856 ^p	AGN1	0.735	$1.88^{+0.17}_{-0.16}$	<0.03	0.51	44.13	31/19	4.5	BSS	N,6
XBSJ061342.7+710725	BL Lac	0.267	$2.70^{+0.19}_{-0.13}$	$0.16^{+0.05}_{-0.04}$	1.84	43.75	60/71	81.5	BSS	Y
XBSJ062134.8–643150	AGN1	1.277	$1.89^{+2.87}_{-0.65}$	$0.06^{+5.57}_{-0.06}$	0.66	44.82	0.43/1	51.4	BSS	Y
XBSJ065214.1+743230	AGN1	0.620	$2.01^{+0.16}_{-0.15}$	$0.11^{+0.08}_{-0.07}$	1.58	44.47	78/99	94.2	BSS	Y
XBSJ065400.0+742045	AGN1	0.362	$2.30^{+0.31}_{-0.19}$	$0.02^{+0.08}_{-0.02}$	0.65	43.57	66/65	42.4	BSS	Y
XBSJ065744.3–560817	AGN1	0.120	$2.35^{+0.14}_{-0.12}$	<0.02	0.56	42.41	51/46	29.3	BSS	Y
XBSJ065839.5–560813 ^p	AGN1	0.211	$3.41^{+0.24}_{-0.09}$	<0.007	0.14	42.43	151/70	6×10^{-6}	BSS	N,5
XBSJ074202.7+742625	AGN1	0.599	$2.01^{+0.16}_{-0.14}$	$0.07^{+0.06}_{-0.06}$	1.65	44.46	64/75	82.7	HBSS,BSS	Y
XBSJ074312.1+742937 ^p	AGN1	0.312	$2.27^{+0.06}_{-0.03}$	<0.001	7.56	44.48	281/194	0.004	HBSS,BSS	N,5
XBSJ074338.7+495431	AGN1	0.221	$2.06^{+0.17}_{-0.15}$	<0.03	1.27	43.33	29/21	12.2	BSS	Y
XBSJ074352.0+744258	AGN1	0.800	$2.03^{+0.11}_{-0.10}$	<0.04	1.03	44.56	61/59	39.1	BSS	Y
XBSJ075117.9+180856	AGN1 ^e	0.255	$1.61^{+0.13}_{-0.13}$	$0.13^{+0.07}_{-0.06}$	1.59	43.53	31/36	71.8	BSS	Y
XBSJ080504.6+245156	AGN1	0.980	$2.08^{+0.17}_{-0.16}$	<0.04	0.46	44.44	32/27	24.2	BSS	Y
XBSJ080608.1+244420	AGN1	0.357	$2.49^{+0.07}_{-0.05}$	<0.006	1.55	43.96	161/152	29.1	BSS	Y
XBSJ083049.8+524908	AGN1	1.200	$1.76^{+0.06}_{-0.05}$	<0.06	1.12	44.94	129/156	94.0	BSS	Y
XBSJ083737.0+255151	AGN1 ^e	0.105	$1.79^{+0.48}_{-0.41}$	$0.36^{+0.39}_{-0.30}$	3.15	43.02	4/9	90.9	HBSS,BSS	Y
XBSJ083737.1+254751	AGN1	0.080	$1.92^{+0.14}_{-0.12}$	$0.15^{+0.04}_{-0.02}$	6.62	43.09	58/69	82.0	HBSS,BSS	Y
XBSJ083838.6+253616	AGN1	0.601	$2.22^{+0.49}_{-0.27}$	<0.22	0.42	43.91	6/8	67.8	BSS	Y
XBSJ083905.9+255010	AGN1	0.250	$2.01^{+0.60}_{-0.35}$	<0.13	0.56	43.10	1/5	94.9	BSS	Y

Table 3. continued.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{ s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample	Best fit
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ084026.2+650638	AGN1	1.144	$1.65^{+0.18}_{-0.19}$	$0.16^{+0.30}_{-0.16}$	1.22	44.90	22/25	63.1	BSS	Y
XBSJ084651.7+344634	AGN1	1.115	$2.15^{+0.12}_{-0.12}$	<0.05	0.37	44.51	54/42	10.2	BSS	Y
XBSJ085530.7+585129	AGN1	0.905	$1.67^{+0.29}_{-0.28}$	<0.12	0.99	44.58	20/25	75.1	BSS	Y
XBSJ090729.1+620824	AGN2 ^e	0.388	1.90^f	$0.90^{+1.60}_{-0.62}$	0.70	43.63	5/3	19.3	BSS	Y
XBSJ091828.4+513931 ^p	AGN1	0.185	1.90^f	$4.58^{+1.02}_{-0.80}$	2.41	43.53	40/17	0.1	HBSS,BSS	N,8
XBSJ094526.2-085006	AGN1 ^e	0.314	$2.25^{+1.23}_{-0.78}$	$0.36^{+0.81}_{-0.36}$	0.77	43.49	8/5	14.2	BSS	Y
XBSJ094548.3-084824	AGN1	1.748	$1.75^{+0.83}_{-0.08}$	<1.18	0.61	45.15	0.56/4	96.9	BSS	Y
XBSJ095054.5+393924	AGN1	1.299	$2.01^{+0.38}_{-0.26}$	<0.32	0.58	44.82	10/9	36.3	BSS	Y
XBSJ095218.9-013643 ^p	AGN1	0.020	$3.65^{+0.50}_{-0.43}$	$0.05^{+0.06}_{-0.05}$	0.65	40.85	214/58	1×10^{-17}	HBSS,BSS	N,8
XBSJ095309.7+013558	AGN1	0.477	$1.89^{+0.40}_{-0.28}$	<0.16	0.74	43.85	8/11	70.1	BSS	Y
XBSJ095509.6+174124	AGN1	1.290	$1.90^{+0.14}_{-0.09}$	<0.11	0.79	44.91	84/85	51.0	BSS	Y
XBSJ095606.4+411814	AGN1	1.150	$2.27^{+0.40}_{-0.32}$	$0.22^{+0.28}_{-0.20}$	0.45	44.68	13/14	48.9	BSS	Y
XBSJ100032.5+553626	AGN2 ^e	0.216	$1.85^{+0.21}_{-0.18}$	<0.03	1.01	43.19	22/18	24.1	BSS	Y
XBSJ100100.0+252103	AGN1	0.794	$2.20^{+0.12}_{-0.07}$	<0.05	0.56	44.33	92/84	24.5	BSS	Y
XBSJ100309.4+554135	AGN1	0.673	$2.27^{+0.12}_{-0.10}$	<0.03	0.51	44.13	61/63	55.0	BSS	Y
XBSJ100828.8+535408	AGN1	0.384	$2.04^{+0.20}_{-0.15}$	<0.04	0.44	43.43	30/23	15.1	BSS	Y
XBSJ100921.7+534926	AGN1	0.387	$2.35^{+0.14}_{-0.08}$	<0.03	0.62	43.63	53/60	71.4	BSS	Y
XBSJ100926.5+533426 ^p	AGN1	1.718	$2.01^{+0.13}_{-0.12}$	<0.06	0.66	45.19	51/39	9.8	BSS	N,10
XBSJ101506.0+520157 ^p	AGN1	0.610	$2.00^{+1.57}_{-1.03}$	<0.19	0.53	43.98	20/9	1.6	BSS	Y
XBSJ101511.8+520708	AGN1	0.888	$2.34^{+0.66}_{-0.46}$	$0.15^{+0.30}_{-0.15}$	0.54	44.48	19/20	52.3	BSS	Y
XBSJ101706.5+520245	BL Lac	0.377	$2.68^{+0.43}_{-0.21}$	<0.07	0.34	43.38	37/30	16.8	BSS	Y
XBSJ101838.0+411635	AGN1	0.577	$2.36^{+0.11}_{-0.10}$	<0.01	0.49	43.96	53/56	59.5	BSS	Y
XBSJ101843.0+413515	AGN1 ^e	0.084	$1.86^{+0.20}_{-0.11}$	<0.04	0.99	42.30	39/37	39.2	BSS	Y
XBSJ101850.5+411506	AGN1	0.577	$2.30^{+0.08}_{-0.05}$	<0.02	1.50	44.43	133/122	23.8	HBSS,BSS	Y
XBSJ101922.6+412049 ^p	AGN1	0.239	$1.88^{+0.06}_{-0.06}$	<0.004	2.12	43.62	136/112	6.4	HBSS,BSS	N,7
XBSJ102252.0+194837	AGN1	0.910	$2.27^{+0.11}_{-0.11}$	<0.02	0.52	44.47	58/57	44.0	BSS	Y
XBSJ102412.3+042023 ^p	AGN1	1.458	$2.17^{+0.14}_{-0.12}$	<0.09	0.39	44.84	55/35	1.7	BSS	N,10
XBSJ102417.5+041656	AGN1	1.712	$1.78^{+0.12}_{-0.12}$	$0.48^{+0.24}_{-0.23}$	1.08	45.29	55/60	67.7	BSS	Y
XBSJ103120.0+311404	AGN1	1.190	$1.85^{+0.19}_{-0.14}$	$0.04^{+0.18}_{-0.04}$	1.22	45.00	29/32	60.0	BSS	Y
XBSJ103154.1+310732	AGN1	0.299	$1.88^{+0.22}_{-0.19}$	<0.03	0.65	43.32	21/19	36.6	BSS	Y
XBSJ103745.7+532353	AGN1	2.347	$2.25^{+0.71}_{-0.49}$	$0.55^{+1.17}_{-0.55}$	0.37	45.39	27/25	37.5	BSS	Y
XBSJ103909.4+205222	AGN1	0.980	$1.96^{+0.26}_{-0.16}$	<0.15	1.18	44.82	23/17	13.4	BSS	Y
XBSJ103932.7+205426	AGN1	0.237	$1.87^{+0.18}_{-0.15}$	<0.03	1.16	43.34	14/20	81.1	BSS	Y
XBSJ103935.8+533036	AGN1	0.229	$2.08^{+0.25}_{-0.16}$	$0.01^{+0.05}_{-0.01}$	1.38	43.40	60/73	87.0	BSS	Y
XBSJ104026.9+204542	AGN1	0.465	$1.99^{+0.05}_{-0.05}$	<0.009	6.21	44.76	135/138	54.3	HBSS,BSS	Y
XBSJ104034.3+205110	AGN1	0.670	$2.26^{+0.39}_{-0.18}$	<0.07	0.40	44.02	10/11	48.6	BSS	Y
XBSJ104425.0-013521	AGN1	1.571	$1.85^{+0.16}_{-0.14}$	$0.12^{+0.26}_{-0.12}$	0.84	45.13	41/41	45.6	BSS	Y
XBSJ104509.3-012442	AGN1	0.472	$2.14^{+0.18}_{-0.10}$	<0.03	0.49	43.70	41/46	68.5	BSS	Y

Table 3. continued.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{ s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample	Best fit
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ104522.1–012843 ^p	AGN1	0.782	$2.00^{+0.13}_{-0.03}$	<0.02	2.26	44.87	101/83	8.3	HBSS,BSS	Y
XBSJ104912.8+330459	AGN1	0.226	$1.67^{+0.19}_{-0.15}$	$0.02^{+0.05}_{-0.02}$	1.84	43.48	97/112	82.9	HBSS,BSS	Y
XBSJ105014.9+331013	AGN1	1.012	$2.33^{+0.61}_{-0.33}$	$0.05^{+0.30}_{-0.05}$	0.36	44.44	8/12	81.0	BSS	Y
XBSJ105239.7+572431	AGN1	1.113	$2.10^{+0.04}_{-0.03}$	<0.006	0.83	44.85	290/261	10.8	BSS	Y
XBSJ105316.9+573551	AGN1	1.204	$1.80^{+0.04}_{-0.03}$	<0.009	1.77	45.16	251/249	45.4	BSS	Y
XBSJ105335.0+572540	AGN1	0.784	$1.72^{+0.08}_{-0.06}$	<0.03	0.50	44.15	80/86	65.5	BSS	Y
XBSJ105339.7+573104	AGN1	0.586	$2.16^{+0.07}_{-0.07}$	<0.007	0.35	43.79	109/119	73.4	BSS	Y
XBSJ105624.2–033522	AGN1	0.635	$2.16^{+0.15}_{-0.10}$	$0.008^{+0.056}_{-0.008}$	0.58	44.10	81/76	32.1	BSS	Y
XBSJ110652.0–182738	AGN1	1.435	$1.55^{+0.36}_{-0.29}$	<0.52	1.03	45.00	17/21	73.9	BSS	Y
XBSJ111928.5+130250	AGN1	2.394	$1.93^{+0.09}_{-0.06}$	<0.15	1.31	45.80	81/85	60.9	BSS	Y
XBSJ111933.0+212756	AGN1	0.282	$1.92^{+0.32}_{-0.40}$	<0.10	0.79	43.35	6/4	20.8	BSS	Y
XBSJ111942.1+211516	AGN1	1.288	$1.92^{+0.40}_{-0.28}$	<0.41	0.44	44.66	11/7	16.1	BSS	Y
XBSJ112022.3+125252	AGN1	0.406	$2.22^{+0.15}_{-0.13}$	<0.03	1.11	43.91	26/26	46.4	BSS	Y
XBSJ112026.7+431520 ^p	AGN2 ^e	0.146	1.90^f	$6.67^{+1.34}_{-1.06}$	1.83	43.23	31/17	2.0	HBSS	N,8
XBSJ112046.7+125429	AGN1	0.382	$2.29^{+0.25}_{-0.15}$	<0.06	1.04	43.83	20/20	47.4	BSS	Y
XBSJ113106.9+312518 ^p	AGN1	1.482	$1.72^{+0.27}_{-0.24}$	<0.30	1.11	45.13	31/21	7.4	HBSS,BSS	Y
XBSJ113121.8+310252	AGN2	0.190	$1.54^{+0.24}_{-0.22}$	$0.67^{+0.27}_{-0.24}$	2.83	43.51	35/30	23.4	HBSS,BSS	Y
XBSJ113128.6–195903	AGN1	0.363	$2.35^{+0.31}_{-0.21}$	<0.06	0.53	43.49	9/10	55.8	BSS	Y
XBSJ113148.7+311358	AGN2	0.500	1.90^f	$2.90^{+0.66}_{-0.60}$	2.85	44.52	32/33	52.3	HBSS,BSS	Y
XBSJ113837.9–373402	AGN1	0.120	$2.49^{+0.05}_{-0.05}$	<0.003	1.33	42.79	158/140	13.4	BSS	Y
XBSJ115317.9+364712	AGN1	0.725	$2.00^{+1.48}_{-0.73}$	$0.22^{+0.92}_{-0.22}$	0.38	44.02	6/8	69.6	BSS	Y
XBSJ115846.9+551625	AGN1	0.518	$2.11^{+0.30}_{-0.22}$	<0.07	0.63	43.90	31/26	22.4	BSS	Y
XBSJ120359.1+443715 ^p	AGN1	0.641	$2.43^{+0.19}_{-0.19}$	$0.04^{+0.06}_{-0.04}$	0.43	44.03	81/59	3.2	BSS	Y
XBSJ120413.7+443149	AGN1	0.492	$2.23^{+0.22}_{-0.14}$	<0.06	0.44	43.71	45/42	35.6	BSS	Y
XBSJ121501.7+140113	AGN1	0.596	$2.12^{+0.11}_{-0.09}$	<0.03	0.75	44.13	69/57	13.8	BSS	Y
XBSJ122350.4+752231	AGN1	0.565	$1.92^{+0.12}_{-0.11}$	<0.03	0.92	44.12	39/36	33.2	BSS	Y
XBSJ122656.5+013126	AGN2	0.733	$1.61^{+0.24}_{-0.21}$	$2.40^{+0.65}_{-0.56}$	2.23	44.73	69/67	41.1	HBSS,BSS	Y
XBSJ123036.2+642531	AGN1	0.744	$2.25^{+0.34}_{-0.21}$	<0.12	0.32	44.02	24/21	28.4	BSS	Y
XBSJ123116.5+641115	AGN1	0.454	$1.92^{+0.09}_{-0.08}$	<0.02	0.58	43.70	111/126	83.0	BSS	Y
XBSJ123218.5+640311	AGN1	1.013	$1.88^{+0.25}_{-0.22}$	$0.15^{+0.16}_{-0.14}$	0.79	44.66	68/67	43.3	BSS	Y
XBSJ123538.6+621644	AGN1	0.717	$1.97^{+0.07}_{-0.05}$	<0.03	1.03	44.43	104/117	80.3	BSS	Y
XBSJ123759.6+621102	AGN1	0.910	$2.05^{+0.06}_{-0.06}$	<0.02	0.98	44.69	139/131	30.7	BSS	Y
XBSJ123800.9+621338	AGN1	0.440	$2.54^{+0.07}_{-0.09}$	<0.01	0.23	43.36	113/101	18.8	BSS	Y
XBSJ124214.1–112512	AGN1	0.820	$1.81^{+0.09}_{-0.08}$	<0.03	1.05	44.54	66/57	18.6	BSS	Y
XBSJ124557.6+022659	AGN1	0.708	$2.72^{+0.60}_{-0.47}$	$0.14^{+0.22}_{-0.14}$	0.26	44.00	9/18	96.0	BSS	Y
XBSJ124607.6+022153	AGN1	0.491	$2.46^{+0.19}_{-0.14}$	<0.04	0.63	43.91	33/32	43.8	BSS	Y
XBSJ124641.8+022412	AGN1	0.934	$2.21^{+0.11}_{-0.09}$	<0.04	1.36	44.90	76/69	26.7	HBSS,BSS	Y
XBSJ124647.9+020955	AGN1	1.074	$2.08^{+0.66}_{-0.45}$	<0.27	0.39	44.48	0.84/1	54.0	BSS	Y

Table 3. continued.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample	Best fit
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ124903.6–061049	AGN1	0.646	$2.24^{+0.09}_{-0.08}$	<0.02	0.81	44.27	89/75	13.1	BSS	Y
XBSJ124914.6–060910	AGN1	1.627	$2.14^{+0.13}_{-0.12}$	<0.07	0.41	44.97	56/45	13.2	BSS	Y
XBSJ124949.4–060722	AGN1	1.053	$2.16^{+0.11}_{-0.10}$	<0.03	0.57	44.64	57/53	32.6	BSS	Y
XBSJ125457.2+564940	AGN1	1.261	$1.74^{+0.16}_{-0.14}$	$0.18^{+0.17}_{-0.15}$	1.43	45.09	34/35	54.3	BSS	Y
XBSJ125648.4+570349	AGN1	0.860	$1.59^{+0.19}_{-0.13}$	<0.16	0.98	44.51	31/23	12.3	BSS	Y
XBSJ130619.7–233857	AGN1	0.351	$2.49^{+0.46}_{-0.34}$	$0.07^{+0.14}_{-0.07}$	0.97	43.74	22/17	18.3	BSS	Y
XBSJ130658.1–234849	AGN1	0.375	$1.96^{+0.40}_{-0.30}$	<0.21	0.80	43.65	6/7	55.6	BSS	Y
XBSJ132038.0+341124	AGN1	0.065	$1.74^{+0.10}_{-0.10}$	$0.14^{+0.03}_{-0.03}$	2.51	42.47	102/94	28.5	HBSS,BSS	Y
XBSJ132052.5+341742	AGN1	0.844	$1.73^{+0.20}_{-0.19}$	$0.18^{+0.18}_{-0.15}$	1.18	44.60	20/23	63.2	BSS	Y
XBSJ132101.6+340656	AGN1	0.335	$2.44^{+0.06}_{-0.06}$	<0.005	0.94	43.66	155/134	10.1	BSS	Y
XBSJ132105.5+341459	AGN1	0.452	$1.87^{+0.14}_{-0.12}$	$0.02^{+0.05}_{-0.02}$	1.72	44.16	60/55	30.2	BSS	Y
XBSJ133023.8+241707	AGN1	1.438	$2.23^{+0.12}_{-0.08}$	<0.09	0.48	44.94	62/65	58.3	BSS	Y
XBSJ133026.6+241520	BL Lac	0.460	$2.72^{+0.10}_{-0.06}$	<0.02	0.44	43.72	114/137	92.9	BSS	Y
XBSJ133807.5+242411	AGN1	0.631	$2.08^{+0.16}_{-0.14}$	<0.01	0.61	44.09	40/30	10.7	BSS	Y
XBSJ133942.6–315004	AGN1 ^e	0.114	$1.66^{+0.22}_{-0.20}$	$0.25^{+0.12}_{-0.09}$	1.71	42.81	38/31	18.6	HBSS,BSS	Y
XBSJ134656.7+580315	AGN2 ^e	0.373	$1.90^{+0.12}_{-0.10}$	$9.27^{+4.23}_{-2.91}$	1.49	44.04	48/41	21.1	HBSS	Y
XBSJ134749.9+582111	AGN1	0.646	$2.20^{+0.03}_{-0.03}$	<0.001	5.13	45.07	562/562	46.7	HBSS,BSS	Y
XBSJ140100.0–110942 ^p	AGN1 ^e	0.164	$2.37^{+0.12}_{-0.08}$	<0.008	0.39	42.55	81/63	6.5	BSS	N,6
XBSJ140102.0–111224	AGN1	0.037	$1.91^{+0.03}_{-0.03}$	<0.004	5.16	41.80	292/307	70.7	HBSS,BSS	Y
XBSJ140113.4+024016	AGN1	0.631	$1.99^{+0.45}_{-0.21}$	<0.17	0.36	43.85	30/25	24.2	HBSS	Y
XBSJ140127.7+025605 ^p	AGN1	0.265	$1.58^{+0.05}_{-0.05}$	$0.15^{+0.02}_{-0.02}$	7.34	44.23	419/361	1.9	HBSS,BSS	N,6
XBSJ140921.1+261336	AGN1	1.100	$1.48^{+0.08}_{-0.04}$	$0.17^{+0.10}_{-0.09}$	1.99	45.02	103/102	45.3	BSS	Y
XBSJ141235.8–030909	AGN2	0.601	$1.77^{+0.26}_{-0.23}$	$0.23^{+0.18}_{-0.14}$	1.39	44.34	16/20	69.9	BSS	Y
XBSJ141531.5+113156	AGN1	0.257	$1.82^{+0.06}_{-0.05}$	<0.01	2.05	43.67	129/126	43.1	HBSS,BSS	Y
XBSJ141643.8+521434	AGN1	0.531	$1.82^{+0.07}_{-0.08}$	<0.02	0.89	44.02	80/80	46.7	BSS	Y
XBSJ141658.8+521202	AGN1	0.600	$2.15^{+0.12}_{-0.11}$	<0.02	0.35	43.82	27/40	94.8	BSS	Y
XBSJ141722.6+251335	AGN1	0.560	$2.26^{+0.37}_{-0.17}$	<0.07	0.53	43.94	12/14	62.7	BSS	Y
XBSJ141736.3+523028 ^p	AGN1	0.985	$2.00^{+0.07}_{-0.07}$	<0.02	0.67	44.59	122/99	6.0	BSS	Y
XBSJ141809.1+250040	AGN1	0.727	$1.93^{+0.19}_{-0.16}$	<0.06	0.78	44.31	15/16	51.0	BSS	N,5
XBSJ142741.8+423335	AGN2 ^e	0.142	$1.90^{+0.12}_{-0.10}$	$4.48^{+0.93}_{-0.76}$	2.12	43.23	19/21	60.9	HBSS	Y
XBSJ143835.1+642928	AGN2 ^e	0.118	$1.75^{+0.41}_{-0.34}$	$1.74^{+0.68}_{-0.53}$	2.42	43.04	12/16	74.0	HBSS,BSS	Y
XBSJ143911.2+640526	AGN2 ^e	0.113	$1.90^{+0.12}_{-0.10}$	$17.50^{+8.10}_{-5.40}$	0.97	42.89	9/10	55.1	HBSS	Y
XBSJ144021.0+642144 ^p	AGN1	0.720	$1.88^{+0.22}_{-0.19}$	$0.67^{+0.34}_{-0.25}$	1.18	44.48	51/31	1.3	BSS	Y
XBSJ144937.5+090826	AGN1	1.260	$1.81^{+0.11}_{-0.07}$	$0.009^{+0.106}_{-0.009}$	1.42	45.12	81/79	43.7	BSS	Y
XBSJ144945.8+085921	AGN1^e	0.265	$1.97^{+0.09}_{-0.09}$	<0.02	0.58	43.14	64/57	24.0	BSS	Y
XBSJ145857.1–313535	AGN1	1.045	$1.97^{+0.26}_{-0.15}$	<0.22	0.53	44.55	54/53	44.9	BSS	Y
XBSJ150428.3+101856	AGN1	1.000	$2.31^{+0.25}_{-0.16}$	<0.14	0.77	44.76	25/20	19.3	BSS	Y
XBSJ151815.0+060851	AGN1	1.294	$1.90^{+0.12}_{-0.10}$	<0.62	0.76	44.90	2/2	31.1	BSS	Y

Table 3. continued.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{ s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample	Best fit
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ153205.7–082952	AGN1	1.239	$1.99^{+0.13}_{-0.12}$	<0.13	0.62	44.80	41/42	51.5	BSS	Y
XBSJ153419.0+011808	AGN1	1.283	$2.52^{+0.64}_{-0.36}$	$0.11^{+0.47}_{-0.11}$	0.40	44.83	18/20	57.9	BSS	Y
XBSJ153452.3+013104	AGN1	1.435	$1.76^{+0.07}_{-0.06}$	<0.07	7.50	45.95	109/114	61.0	HBSS,BSS	Y
XBSJ153456.1+013033 ^p	AGN1	0.310	$2.59^{+0.14}_{-0.13}$	<0.01	0.85	43.55	86/64	3.4	BSS	N,5
XBSJ160645.9+081525	AGN2	0.618	$1.68^{+0.77}_{-0.65}$	$13.80^{+6.70}_{-5.15}$	3.91	44.92	12/10	28.7	HBSS,BSS	Y
XBSJ160706.6+075709	AGN1	0.233	$2.42^{+0.15}_{-0.14}$	<0.02	0.61	43.10	22/29	81.0	BSS	Y
XBSJ160731.5+081202	AGN1	0.226	$2.67^{+0.36}_{-0.22}$	$0.02^{+0.06}_{-0.02}$	0.40	42.91	31/27	26.8	BSS	Y
XBSJ161544.2+121708	AGN1^e	0.207	$2.22^{+0.39}_{-0.19}$	<0.07	0.56	42.93	18/21	63.9	BSS	Y
XBSJ161615.1+121353	AGN1	0.843	$2.01^{+0.28}_{-0.17}$	<0.17	0.57	44.36	31/23	12.7	BSS	Y
XBSJ161820.7+124116	AGN2 ^e	0.361	$1.90^{+0.17}_{-0.17}$	$3.65^{+6.75}_{-2.43}$	0.83	43.68	0.22/1	64.3	HBSS	Y
XBSJ161825.4+124145	AGN1	0.396	$2.29^{+0.66}_{-0.45}$	<0.16	0.43	43.48	6/7	51.3	BSS	Y
XBSJ162813.9+780342	AGN1	0.640	$2.30^{+0.34}_{-0.29}$	<0.10	1.25	44.45	5/8	72.5	BSS	Y
XBSJ163141.1+781239	AGN1	0.380	$2.52^{+0.33}_{-0.25}$	$0.03^{+0.07}_{-0.03}$	0.43	43.48	26/23	31.2	BSS	Y
XBSJ163223.6+052547	AGN1	0.146	$2.33^{+0.14}_{-0.12}$	<0.01	0.69	42.68	54/46	18.8	BSS	Y
XBSJ163309.8+571039	AGN1	0.288	$2.23^{+0.32}_{-0.18}$	<0.06	0.85	43.44	49/48	43.7	BSS	Y
XBSJ163332.3+570520	AGN1 ^e	0.386	$2.31^{+0.72}_{-0.50}$	$0.08^{+0.18}_{-0.08}$	0.40	43.43	13/19	82.4	BSS	Y
XBSJ163427.5+781002	AGN1	0.376	$2.20^{+0.54}_{-0.41}$	$0.33^{+0.34}_{-0.26}$	0.70	43.64	11/12	50.3	BSS	Y
XBSJ164237.9+030014	AGN1	1.338	$1.72^{+0.31}_{-0.21}$	<0.41	0.99	44.99	9/13	74.0	BSS	Y
XBSJ165314.4+141943	AGN1	0.465	$2.05^{+0.29}_{-0.19}$	<0.10	0.83	43.91	17/21	68.7	BSS	Y
XBSJ165406.6+142123	AGN1	0.641	$1.88^{+0.20}_{-0.14}$	$0.003^{+0.096}_{-0.003}$	0.79	44.18	52/46	26.0	BSS	Y
XBSJ165425.3+142159	AGN1	0.178	$2.11^{+0.06}_{-0.03}$	<0.01	6.33	43.82	324/306	21.3	HBSS,BSS	Y
XBSJ165448.5+141311	AGN1	0.320	$1.81^{+0.12}_{-0.07}$	<0.03	4.76	44.25	89/97	68.9	HBSS,BSS	Y
XBSJ165800.7+352333	AGN1^e	0.127	$1.86^{+0.75}_{-0.39}$	$0.05^{+0.17}_{-0.05}$	1.01	42.69	9/6	15.9	BSS	Y
XBSJ172230.6+341344	AGN1	0.425	$1.90^{+0.87}_{-0.31}$	<0.22	1.52	44.04	12/9	23.5	BSS	Y
XBSJ185518.7–462504	AGN1	0.788	$1.42^{+0.53}_{-0.41}$	<0.48	1.34	44.51	12/17	79.4	BSS	Y
XBSJ185613.7–462239	AGN1	0.768	$2.17^{+0.27}_{-0.23}$	<0.06	1.43	44.70	31/29	35.9	BSS	Y
XBSJ193138.9–725115	AGN1	0.701	$1.89^{+0.27}_{-0.16}$	$0.01^{+0.15}_{-0.01}$	1.14	44.43	46/38	17.6	BSS	Y
XBSJ193248.8–723355	AGN2 ^e	0.287	$1.45^{+0.23}_{-0.22}$	$0.67^{+0.32}_{-0.27}$	2.33	43.80	30/34	69.1	HBSS,BSS	Y
XBSJ204043.4–004548	AGN2	0.615	$1.90^{+0.17}_{-0.17}$	$3.28^{+1.35}_{-0.97}$	1.75	44.52	17/12	16.5	HBSS,BSS	Y
XBSJ204159.2–321439 ^p	AGN1	0.738	$2.14^{+0.20}_{-0.10}$	<0.04	0.88	44.43	40/28	6.8	BSS	N,10
XBSJ204204.1–321601	AGN1	0.384	$2.02^{+0.38}_{-0.21}$	<0.11	0.38	43.35	13/13	43.1	BSS	Y
XBSJ204208.2–323523	AGN1	1.184	$2.01^{+0.37}_{-0.23}$	<0.37	0.37	44.53	19/15	21.7	BSS	Y
XBSJ204548.4–025234 ^p	AGN1	2.188	$1.98^{+0.24}_{-0.25}$	<0.23	0.65	45.42	33/23	8.2	BSS	Y
XBSJ205411.9–160804	AGN1	1.466	$1.86^{+0.28}_{-0.25}$	<0.30	1.06	45.16	7/10	73.0	BSS	Y
XBSJ205429.9–154937	AGN1	1.297	$1.92^{+0.60}_{-0.38}$	$0.20^{+0.83}_{-0.20}$	0.58	44.80	7/9	58.6	BSS	Y
XBSJ205635.7–044717	AGN1	0.217	$2.40^{+0.17}_{-0.14}$	<0.03	1.14	43.30	29/23	17.4	HBSS,BSS	Y
XBSJ205829.9–423634	AGN1	0.232	$1.90^{+0.09}_{-0.08}$	$0.09^{+0.03}_{-0.03}$	3.17	43.76	92/123	98.1	HBSS,BSS	Y
XBSJ210325.4–112011	AGN1	0.720	$1.85^{+0.34}_{-0.20}$	$0.21^{+0.24}_{-0.14}$	0.85	44.33	30/36	75.7	BSS	Y

Table 3. continued.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability %	Sample	Best fit
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ210355.3–121858	AGN1	0.792	$2.20^{+0.41}_{-0.15}$	<0.13	0.58	44.35	18/15	24.5	BSS	Y
XBSJ212759.5–443924	AGN1	0.860	$1.99^{+0.24}_{-0.18}$	<0.10	0.54	44.35	32/38	74.1	BSS	Y
XBSJ213002.3–153414	AGN1	0.562	$2.06^{+0.21}_{-0.19}$	$0.04^{+0.10}_{-0.04}$	1.89	44.46	23/27	67.6	HBSS,BSS	Y
XBSJ213719.6–433347	AGN1	0.793	$1.70^{+0.32}_{-0.30}$	$0.55^{+0.52}_{-0.38}$	2.34	44.83	12/14	60.2	BSS	Y
XBSJ213729.7–423601	AGN1	0.664	$2.02^{+0.43}_{-0.30}$	<0.18	0.81	44.26	12/15	66.1	BSS	Y
XBSJ213733.2–434800	AGN1	0.427	$2.38^{+0.73}_{-0.50}$	$0.10^{+0.18}_{-0.10}$	0.37	43.51	5/5	47.9	BSS	Y
XBSJ213757.6–422334	AGN1	0.364	$2.59^{+2.19}_{-0.72}$	$0.04^{+0.47}_{-0.04}$	0.27	43.22	16/11	13.3	BSS	Y
XBSJ213820.2–142536	AGN1	0.369	$1.60^{+0.14}_{-0.12}$	$0.51^{+0.14}_{-0.12}$	2.27	44.04	97/92	35.8	HBSS,BSS	Y
XBSJ213824.0–423019	AGN1	0.257	$2.16^{+0.11}_{-0.16}$	<0.02	1.49	43.56	72/63	21.6	BSS	Y
XBSJ213829.8–423958	AGN1	1.469	$2.61^{+0.40}_{-0.33}$	<0.24	0.39	44.99	17/12	13.6	BSS	Y
XBSJ213852.2–434714	AGN1	0.461	$3.02^{+0.59}_{-0.38}$	<0.08	0.16	43.32	19/13	11.3	BSS	Y
XBSJ214041.4–234720	AGN1	0.490	$2.17^{+0.09}_{-0.08}$	<0.01	1.73	44.29	103/98	33.5	HBSS,BSS	Y
XBSJ215244.2–302407	AGN1	0.539	$2.36^{+0.10}_{-0.09}$	<0.02	1.00	44.19	59/71	83.8	BSS	Y
XBSJ220320.8+184930	AGN1	0.309	$2.12^{+0.21}_{-0.18}$	$0.30^{+0.13}_{-0.10}$	1.29	43.69	39/36	32.1	BSS	Y
XBSJ220446.8–014535	AGN1	0.540	$1.75^{+0.25}_{-0.20}$	$0.17^{+0.21}_{-0.14}$	0.92	44.04	22/24	56.1	BSS	Y
XBSJ220601.5–015346	AGN1	0.211	$1.65^{+0.16}_{-0.14}$	<0.03	1.50	43.33	29/22	14.6	HBSS,BSS	Y
XBSJ221623.3–174317	AGN1	0.754	$1.82^{+0.25}_{-0.16}$	<0.14	0.68	44.27	17/18	53.8	BSS	Y
XBSJ221722.4–082018	AGN1	1.160	$1.51^{+0.24}_{-0.19}$	$0.05^{+0.27}_{-0.05}$	1.88	45.05	30/34	65.7	BSS	Y
XBSJ221729.3–081154	AGN1	1.008	$2.19^{+0.38}_{-0.19}$	$0.005^{+0.240}_{-0.005}$	0.85	44.78	27/32	74.5	BSS	Y
XBSJ221821.9–081332 ^p	AGN1	0.803	$2.35^{+0.22}_{-0.13}$	<0.04	0.74	44.51	55/40	5.4	BSS	Y
XBSJ221951.6+120123	AGN2	0.532	$1.90^{+0.13}_{-0.12}$	$0.35^{+0.33}_{-0.26}$	0.82	44.01	20/15	17.1	BSS	Y
XBSJ223547.9–255836 ^p	AGN1	0.304	$2.00^{+0.13}_{-0.12}$	<0.02	0.61	43.33	47/34	7.3	BSS	N,6
XBSJ223555.0–255833	AGN1	1.800	$2.17^{+0.15}_{-0.13}$	<0.15	0.43	45.13	26/30	68.5	BSS	Y
XBSJ223949.8+080926	AGN1	1.406	$2.35^{+1.66}_{-0.90}$	<1.20	0.35	44.82	6/3	12.6	BSS	Y
XBSJ224756.6–642721	AGN1	0.598	$2.00^{+0.21}_{-0.17}$	<0.07	0.66	44.06	14/12	30.3	BSS	Y
XBSJ225025.1–643225	AGN1	1.206	$2.09^{+0.19}_{-0.14}$	<0.14	0.44	44.65	34/32	38.3	BSS	Y
XBSJ225050.2–642900	AGN1	1.251	$2.04^{+0.07}_{-0.07}$	<0.03	1.38	45.18	113/102	22.0	BSS	Y
XBSJ225118.0–175951 ^p	AGN1	0.172	$3.18^{+0.11}_{-0.10}$	<0.005	0.36	42.62	212/90	8×10^{-10}	BSS	N,5
XBSJ230400.4–083755	AGN1	0.411	$2.72^{+0.90}_{-0.56}$	$0.12^{+0.29}_{-0.12}$	0.15	43.13	8/8	45.1	BSS	Y
XBSJ230401.0+031519	AGN1 ^e	0.036	$2.05^{+0.60}_{-0.42}$	$0.06^{+0.16}_{-0.06}$	0.56	41.30	8/10	58.5	BSS	Y
XBSJ230434.1+122728	AGN1 ^e	0.232	$1.60^{+0.37}_{-0.30}$	<0.13	1.16	43.29	7/10	72.0	BSS	Y
XBSJ230443.8+121636	AGN1	1.405	$1.95^{+0.29}_{-0.27}$	$0.35^{+0.52}_{-0.35}$	0.82	45.04	19/17	35.0	BSS	Y
XBSJ230459.6+121205	AGN1	0.560	$1.58^{+0.32}_{-0.27}$	<0.28	1.31	44.20	8/10	63.7	BSS	Y
XBSJ230522.1+122121	AGN2	0.326	$1.87^{+0.53}_{-0.45}$	$0.22^{+0.32}_{-0.22}$	0.95	43.58	8/10	66.3	BSS	Y
XBSJ231342.5–423210	AGN1	0.973	$2.14^{+0.13}_{-0.07}$	<0.02	1.10	44.83	166/151	18.2	BSS	Y
XBSJ231546.5–590313 ^p	AGN2 ^e	0.045	$1.90^{+0.07}_{-0.06}$	$0.15^{+0.07}_{-0.06}$	0.99	41.73	38/16	0.15	BSS	N,9
XBSJ231601.7–424038 ^p	AGN1	0.383	$2.70^{+0.20}_{-0.17}$	<0.02	0.28	43.32	68/53	7.4	BSS	N,5
XBSJ233421.9–151219	AGN1	0.992	$2.00^{+0.17}_{-0.14}$	<0.06	0.71	44.62	19/24	72.3	BSS	Y

Table 3. continued.

Source	Type	z	Γ	N_H 10^{22} cm^{-2}	$f_{2-10\text{keV}}$ 10^{-13} $\text{erg cm}^{-2} \text{ s}^{-1}$	$\text{Log } L_{2-10\text{keV}}$ erg s^{-1}	$\chi^2/\text{d.o.f}$	Probability	Sample	Best fit
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)	(11)
XBSJ235036.9+362204	BL Lac	0.317	$1.90^{+0.34}_{-0.22}$	$0.03^{+0.19}_{-0.03}$	2.03	43.88	14/13	39.2	BSS	Y

Columns: (1) Source name; (2) Class; (3) Redshift; (4) Photon index; (5) Intrinsic column density; (6) Observed flux in the 2-10 keV band, de-absorbed by our Galaxy; (7) Intrinsic luminosity in the 2-10 keV band; (8) χ^2 to number of degrees of freedom; (9) Null hypothesis probability; (10) Sample the source belongs to; (11) Whether the simple power law is considered as our best fit (Y) or not (N). In the case it is not, a number indicates the corresponding table's number where our considered best fit is reported.

^p: Null hypothesis probability < 10%. ^e: Elusive AGN. ^f: Fixed parameter. AGN class and redshift in bold face mark new optical identifications. Note: Errors and upper limits are at 90% confidence level. Fluxes and luminosities refer to the MOS2 calibration.